

NONRECIPROcity WITHOUT MAGNETISM: ANGULAR-MOMENTUM BIASED METAMATERIALS

Andrea Alù

Contributors: Dimitrios Sounas, Romain Fleury, Michael Haberman, Caleb F. Sieck

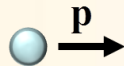
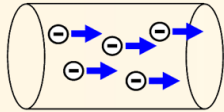
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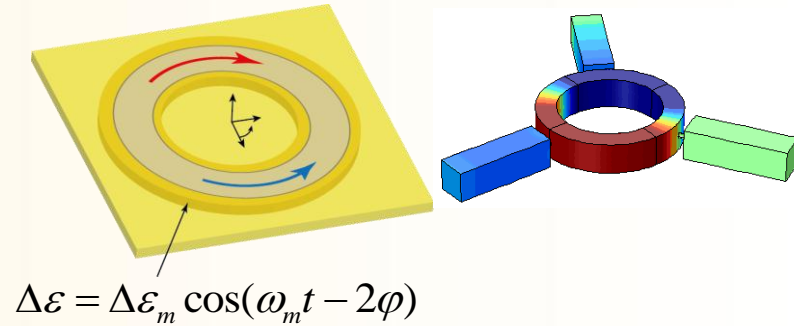
OVERVIEW



Direct current Linear Momentum

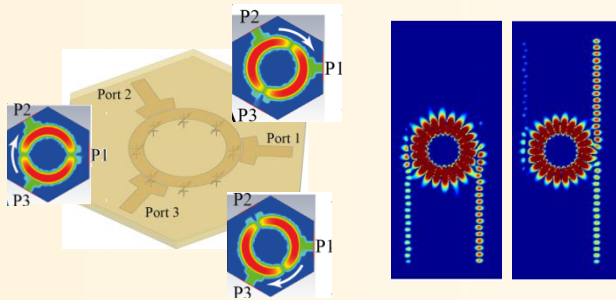


Available approaches to break reciprocity

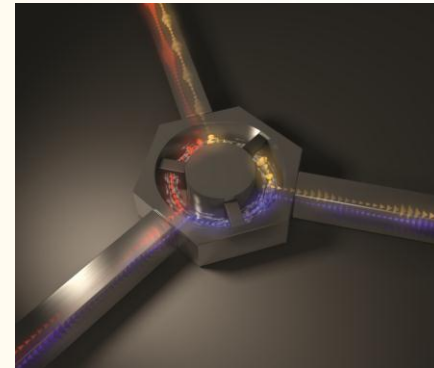


$$\Delta\varepsilon = \Delta\varepsilon_m \cos(\omega_m t - 2\varphi)$$

Angular-momentum biased 'Zeeman' meta-molecules



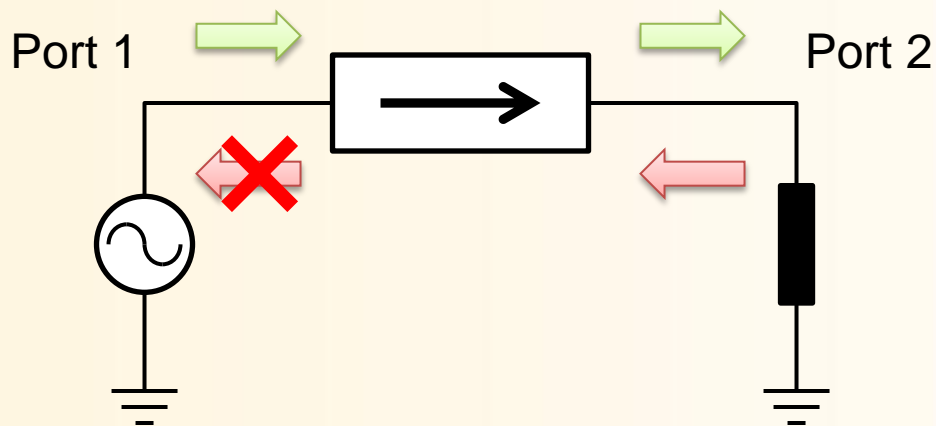
Proposed electromagnetic devices: from RF to optics



A first-off-its-kind acoustic circulator

CONVENTIONAL NONRECIPROCAL COMPONENTS

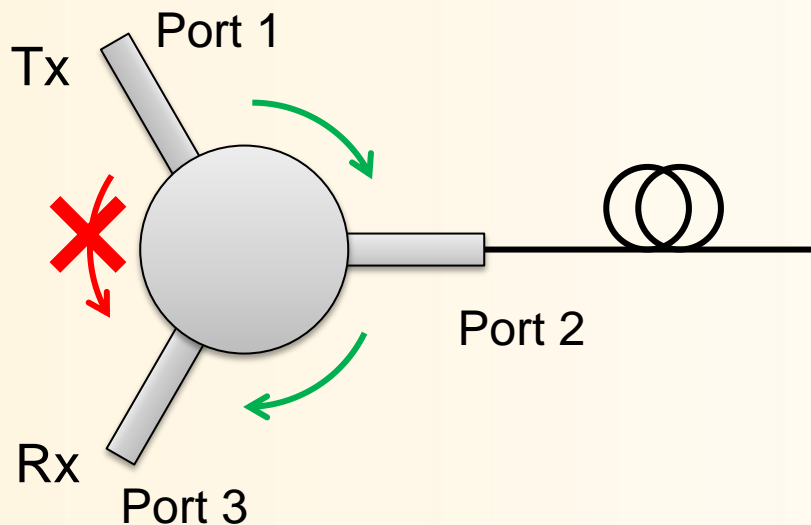
Isolator



$$S_{21} \neq S_{12}$$

$$\bar{S} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}$$

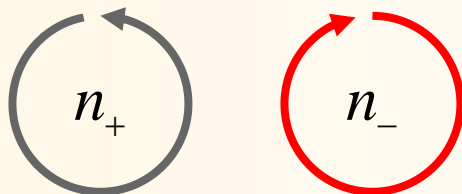
Circulator



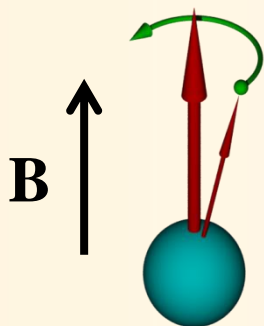
$$S_{ij} \neq S_{ji}$$

$$\bar{S} = \begin{bmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix}$$

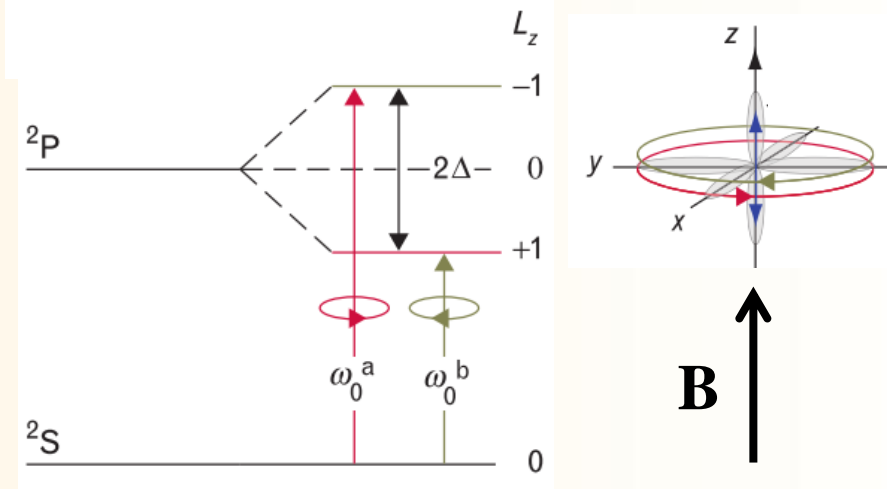
Circular Birefringence



Larmor precession

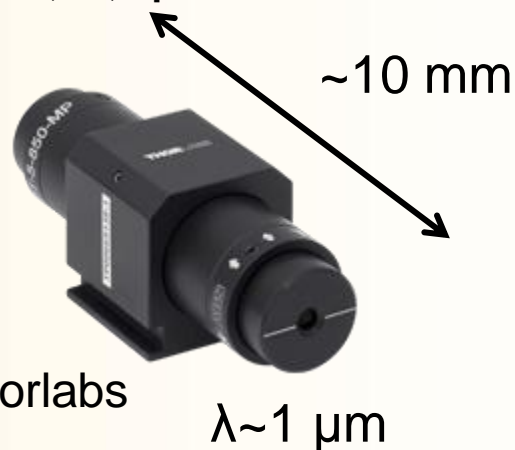


Zeeman splitting

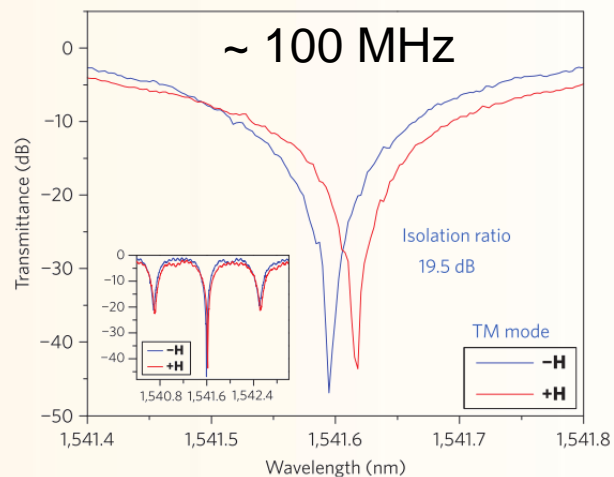
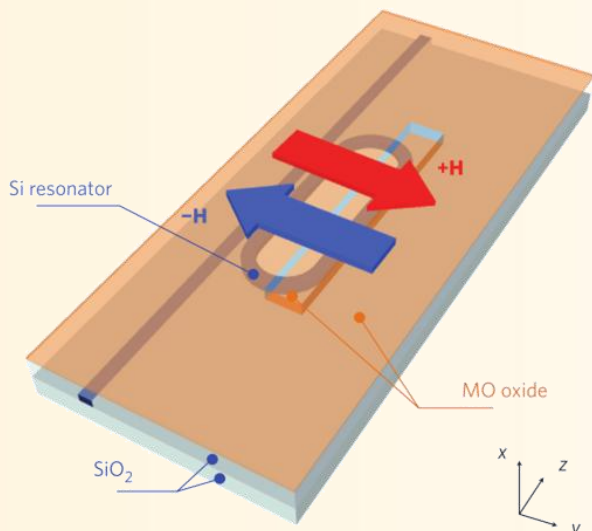


MAGNETICALLY BIASED NONRECIPROCAL DEVICES

Massive devices at microwaves, IR, optics



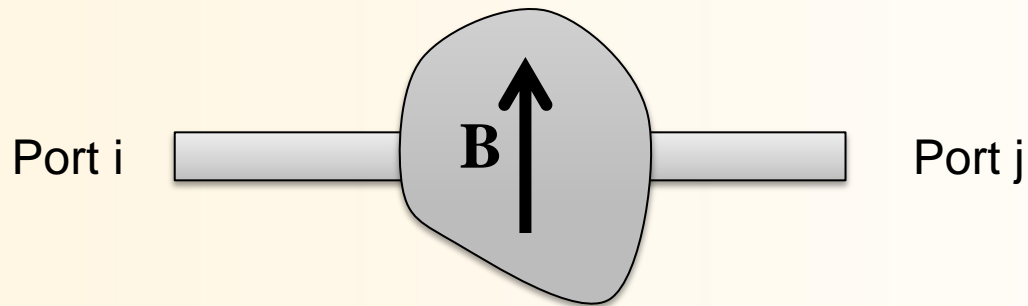
Inherently narrow bandwidths



Bi, L. et al. *Nature Photonics* 5, 758–762 (2011)

DO WE NEED MAGNETIC BIAS?

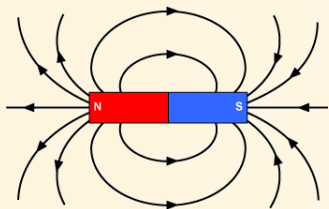
Onsager-Casimir Principle



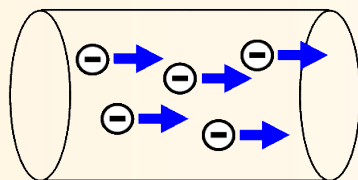
$$S_{ij}(\mathbf{B}) = S_{ji}(-\mathbf{B})$$

↑
Odd vector under time reversal

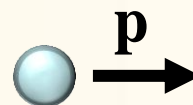
Magnetic Field



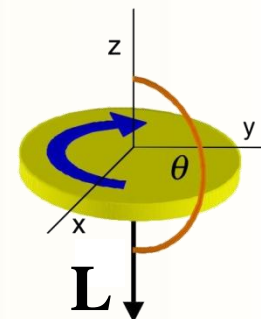
Direct current



Linear Momentum

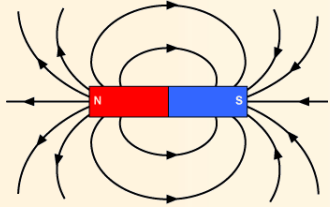


Angular Momentum

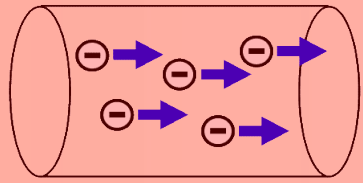


OTHER AVAILABLE APPROACHES

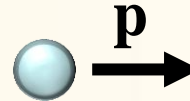
Magnetic Field



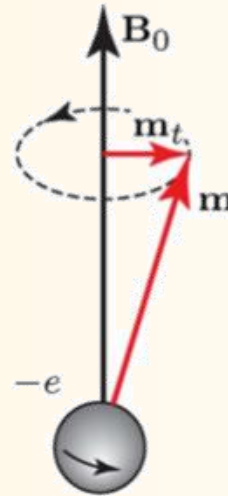
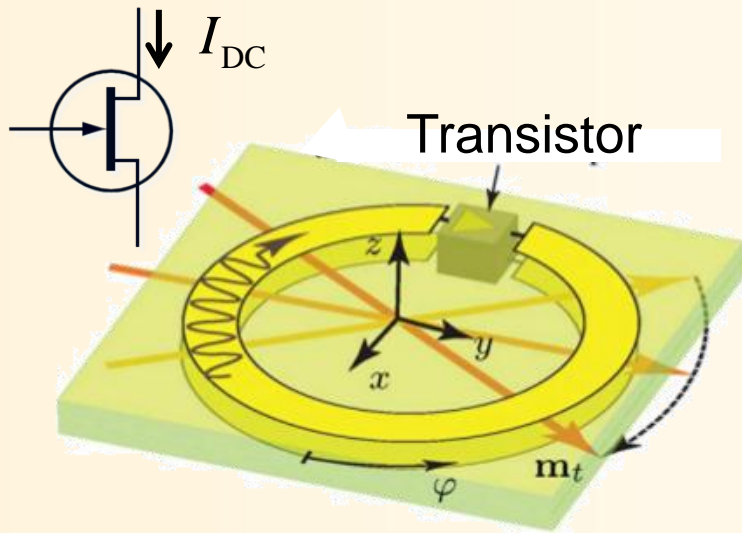
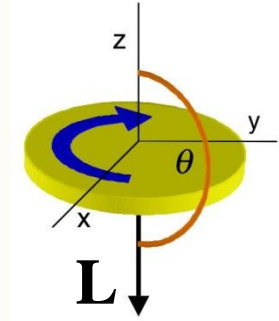
Direct current



Linear Momentum



Angular Momentum

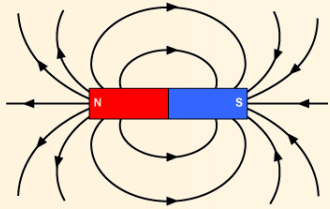


- Only at low microwaves
- Significant power consumption
- Limitations on technology

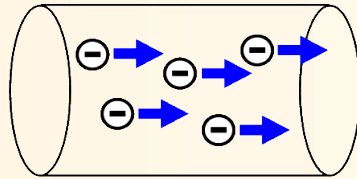
Kodera, T., Sounas, D. L. & Caloz, C. *Appl. Phys. Lett.* 99, 03114 (2011)

OTHER AVAILABLE APPROACHES

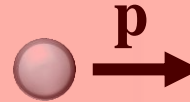
Magnetic Field



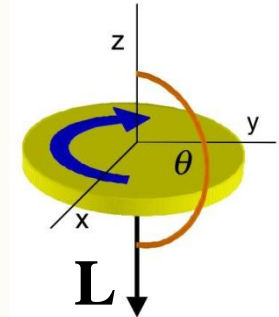
Direct current



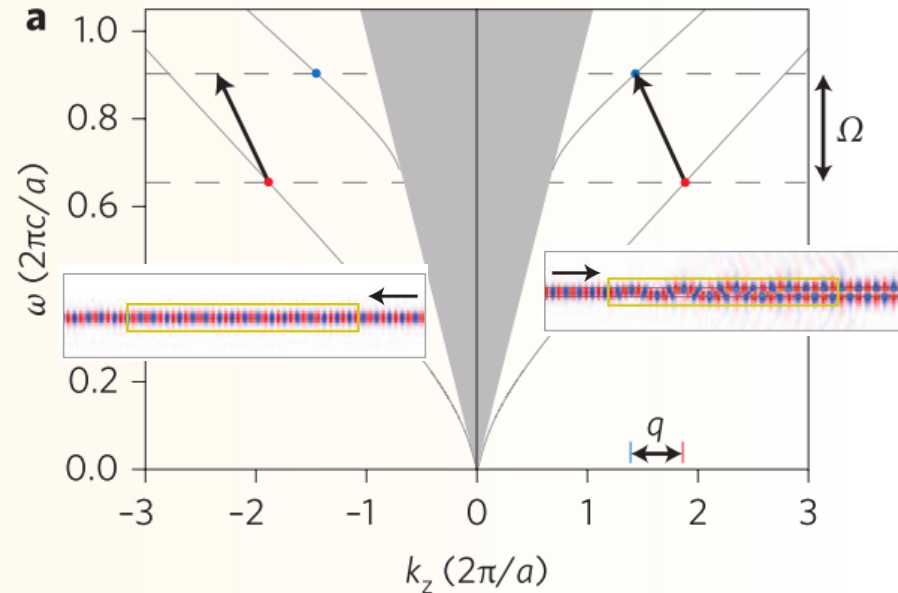
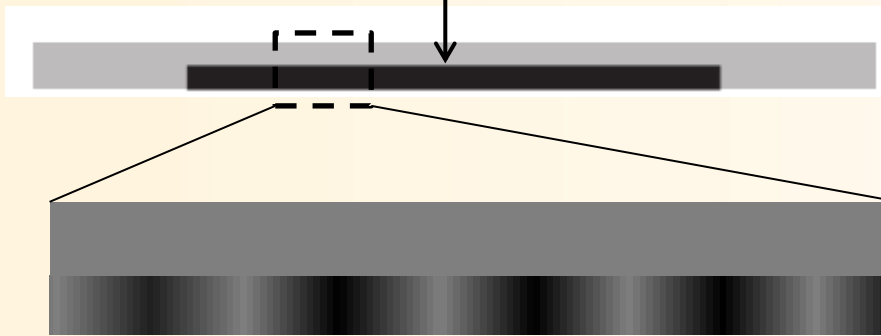
Linear Momentum



Angular Momentum



$$\Delta\epsilon = \Delta\epsilon_m \cos(\Omega t + qz)$$



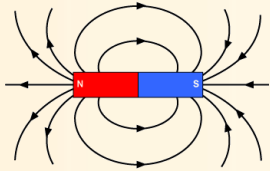
Very large volumes

Non-uniform modulation across the waveguide cross-section

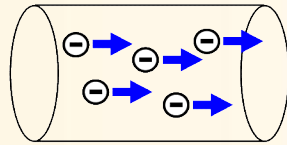
Yu, Z. & Fan, S. *Nature Photonics* 3, 91-94 (2009)

ANGULAR-MOMENTUM BIASING?

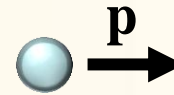
Magnetic Field



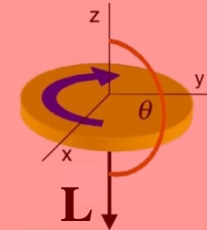
Direct current



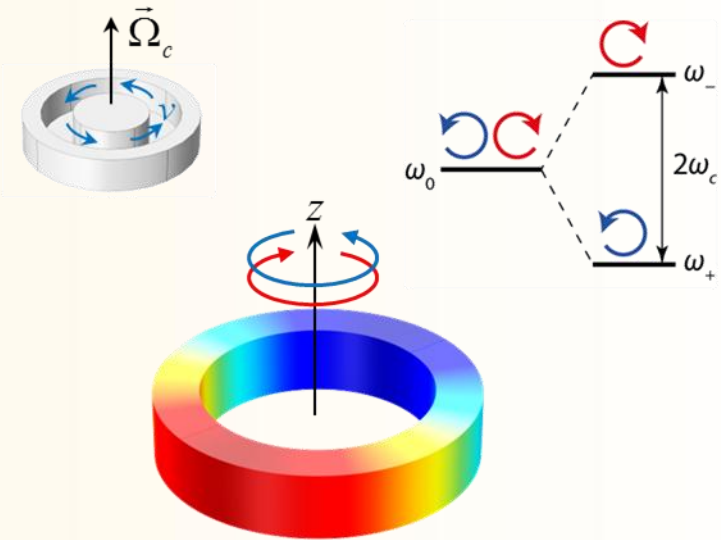
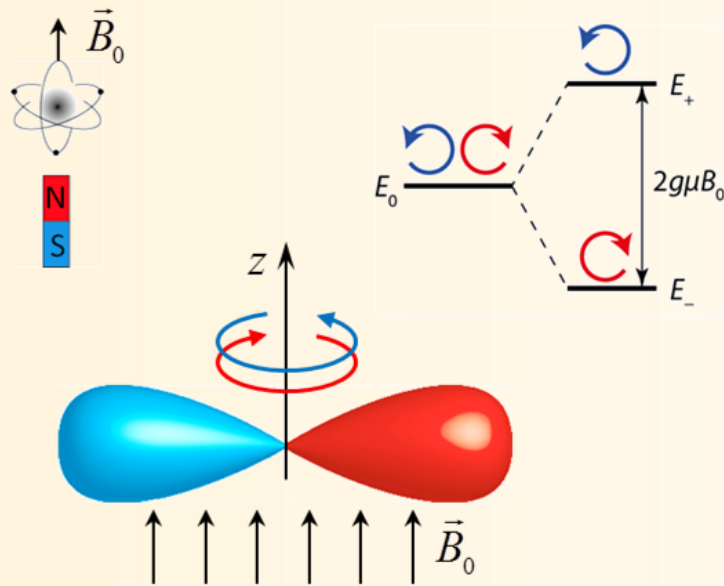
Linear Momentum



Angular Momentum



ANGULAR-MOMENTUM BIASING: A ZEEMAN META-MOLECULE



AN ACOUSTIC RING CAVITY: MODE SPLITTING

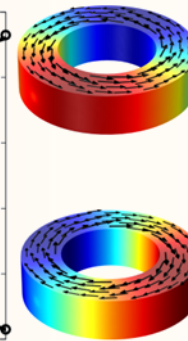
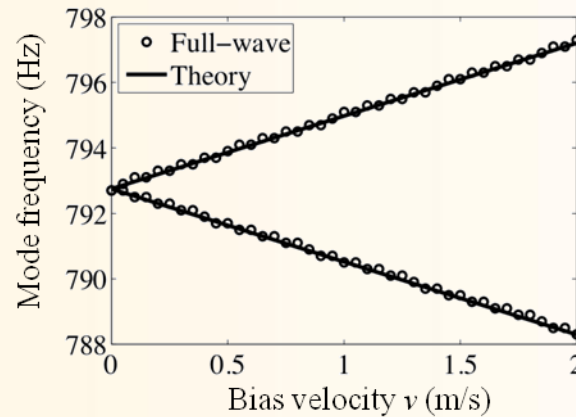
$$|\psi\rangle = \begin{pmatrix} \hat{p} \\ -i\hat{\phi} \end{pmatrix}$$

$$(H_0 + P)|\psi\rangle = \omega|\psi\rangle$$

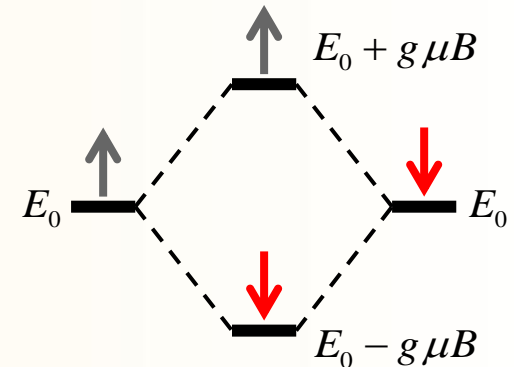
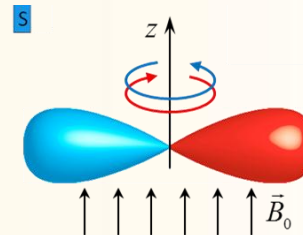
$$H_0 = \begin{pmatrix} 0 & c_0^2 \vec{\nabla} \cdot (\rho_0 \vec{\nabla}) \\ -1/\rho_0 & 0 \end{pmatrix}$$

$$P = \begin{pmatrix} -i(\vec{v} \cdot \vec{\nabla}) - \vec{\nabla} \cdot (i\vec{v}) & 0 \\ 0 & -i(\vec{v} \cdot \vec{\nabla}) \end{pmatrix}$$

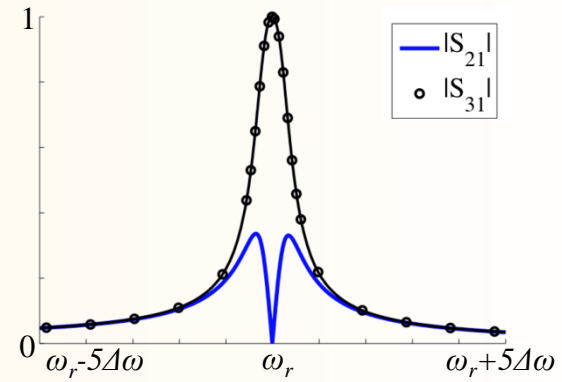
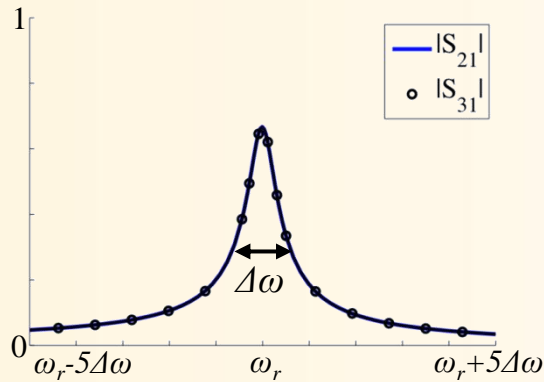
$$\omega^\pm = \omega_0 \pm \langle + | P | + \rangle$$



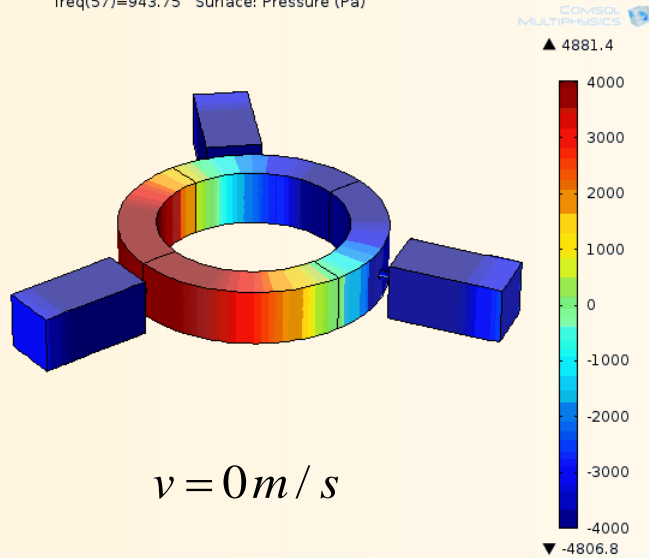
$$\omega_m^\pm = \omega_m \pm \frac{mv}{R_{av}}$$



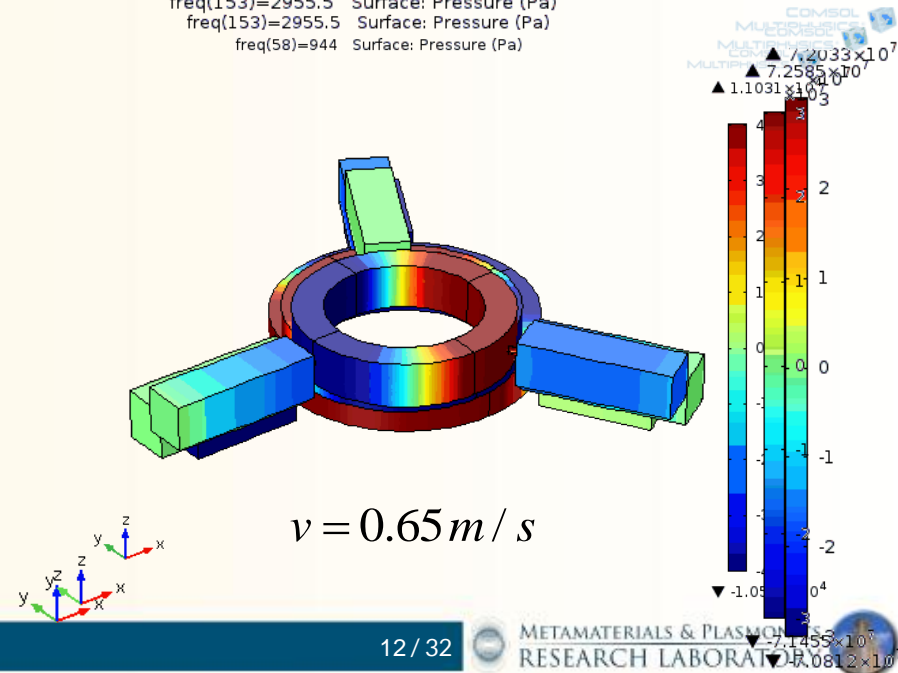
AN ACOUSTIC CIRCULATOR!



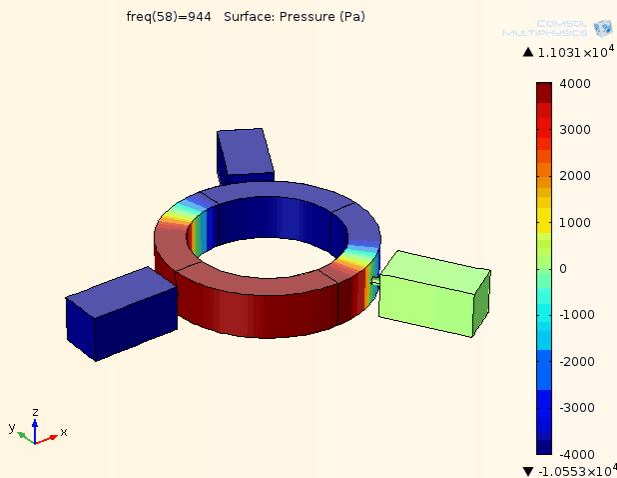
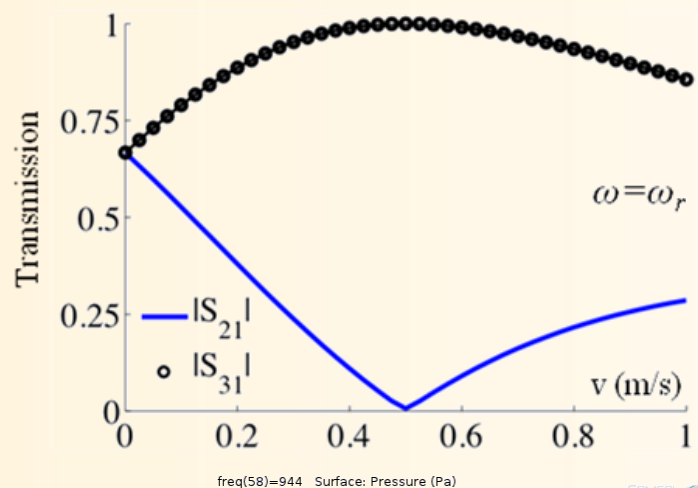
freq(57)=943.75 Surface: Pressure (Pa)



freq(153)=2955.5 Surface: Pressure (Pa)
freq(153)=2955.5 Surface: Pressure (Pa)
freq(58)=944 Surface: Pressure (Pa)



MODE EVOLUTION



$$T_{1 \rightarrow 2} = \left| \frac{2}{3} \left(\frac{e^{-i4\pi/3}}{1 - i(\omega - \omega^-) / \gamma^-} + \frac{e^{-i2\pi/3}}{1 - i(\omega - \omega^+) / \gamma^+} \right) \right|^2$$

$$T_{1 \rightarrow 3} = \left| \frac{2}{3} \left(\frac{e^{-i2\pi/3}}{1 - i(\omega - \omega^-) / \gamma^-} + \frac{e^{-i4\pi/3}}{1 - i(\omega - \omega^+) / \gamma^+} \right) \right|^2$$

$$\downarrow$$

$$\omega^\pm = \omega_r \pm \gamma / \sqrt{3}$$

$$v_{opt} = \gamma R_{av} / \sqrt{3}$$

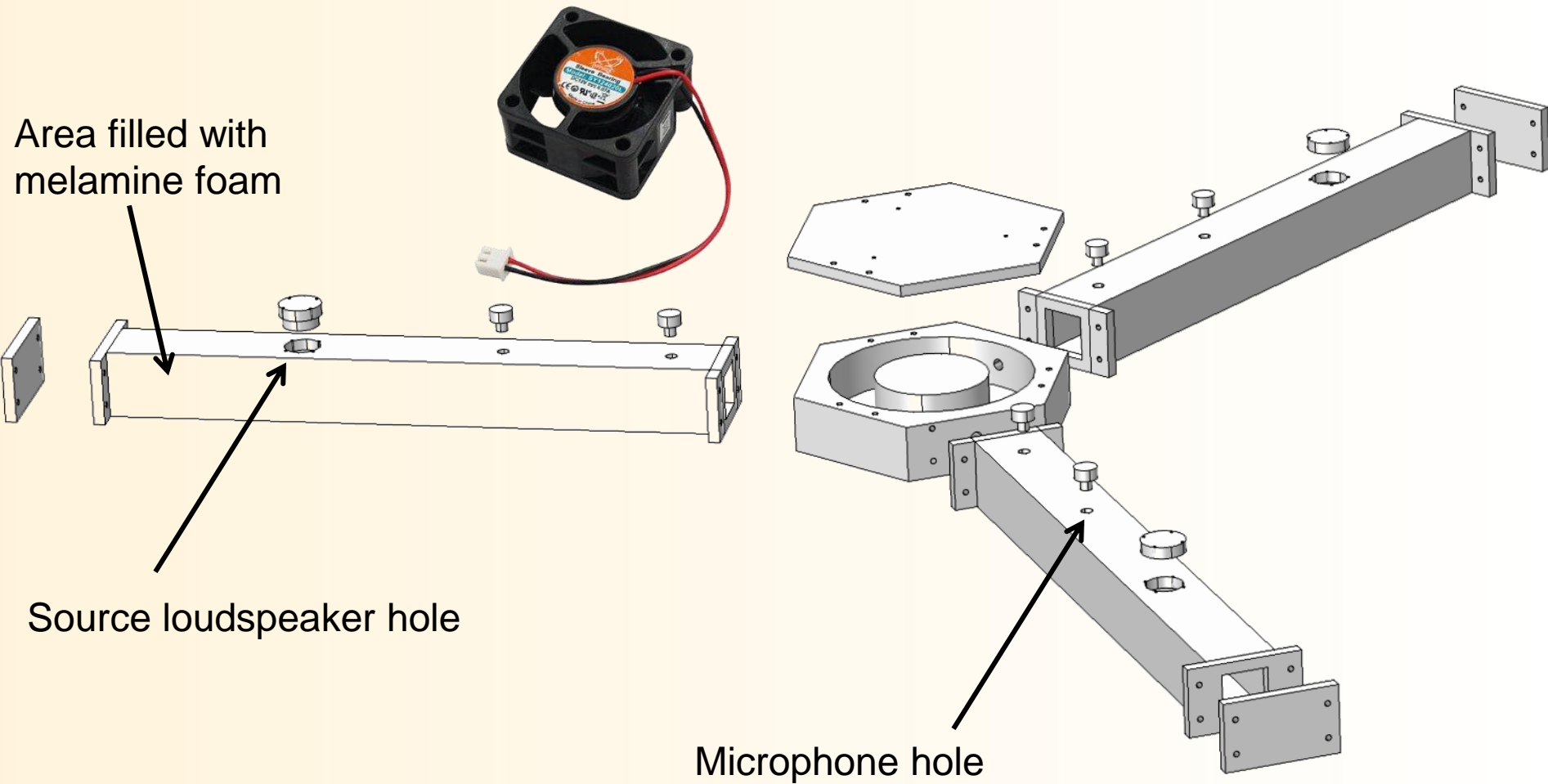
$$\omega_0 = c_0 / R_{av}$$

$$Q = \omega_0 / (2\gamma)$$

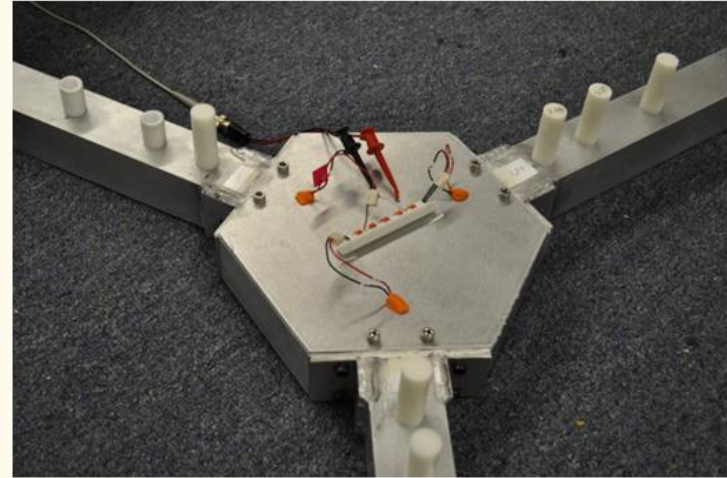
$$v_{opt} / c_0 = 1 / (2Q\sqrt{3})$$

R. Fleury, D. L. Sounas, C. Sieck, M. Haberman, A. Alù, *Science* in press (2014)

AN ACOUSTIC CIRCULATOR



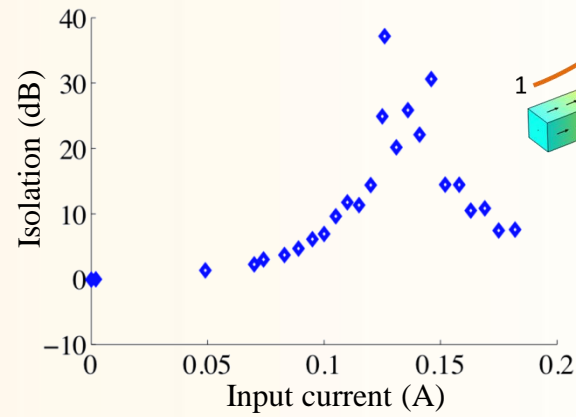
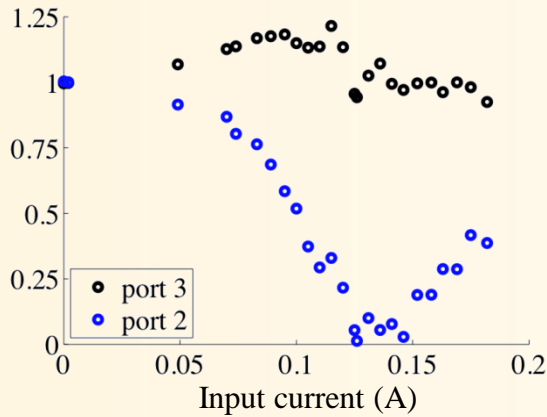
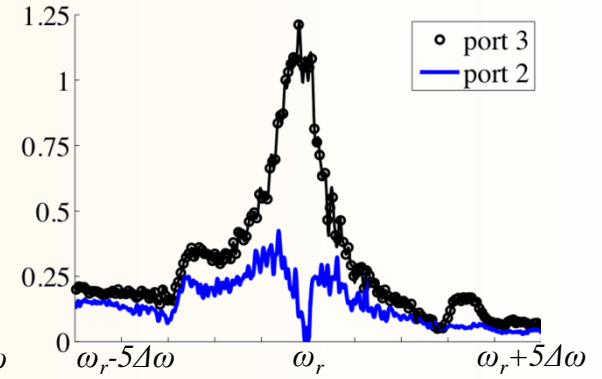
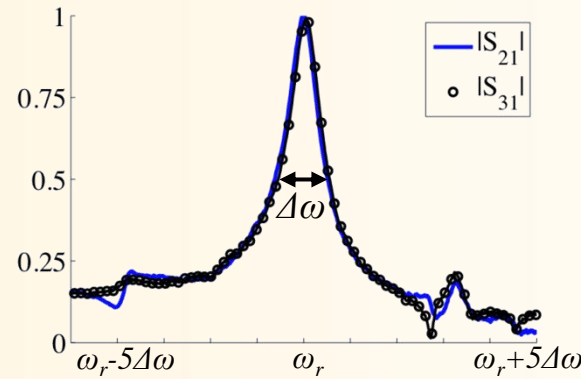
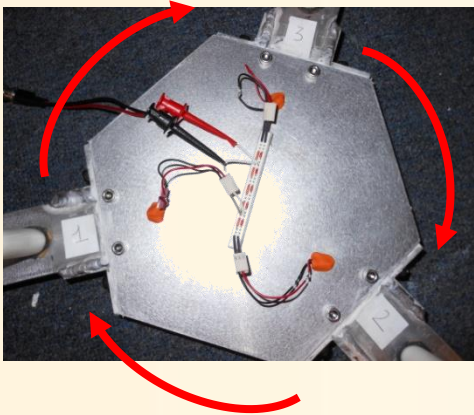
AN ACOUSTIC CIRCULATOR



R. Fleury, D. L. Sounas, C. Sieck, M. Haberman, A. Alù, *Science* in press (2014)



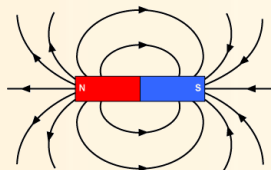
AN ACOUSTIC CIRCULATOR



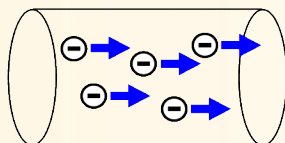
R. Fleury, D. L. Sounas, C. Sieck, M. Haberman, A. Alù, *Science* in press (2014)

HOW ABOUT EM WAVES? AZIMUTHAL SPATIO-TEMPORAL MODULATION

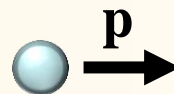
Magnetic Field



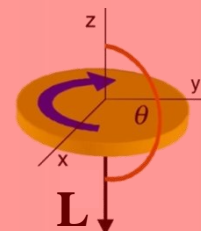
Direct current



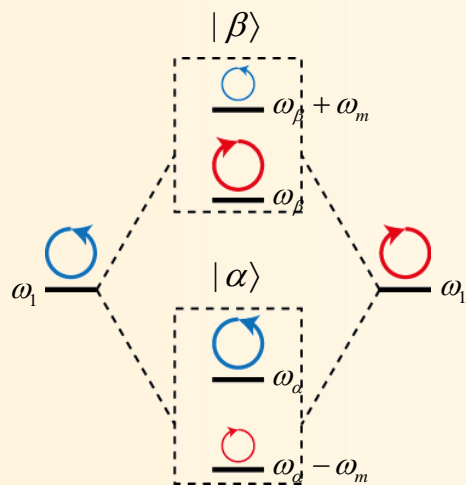
Linear Momentum



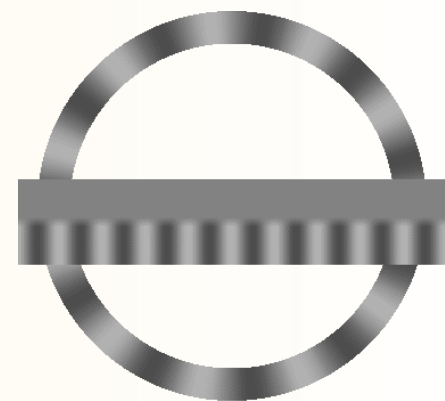
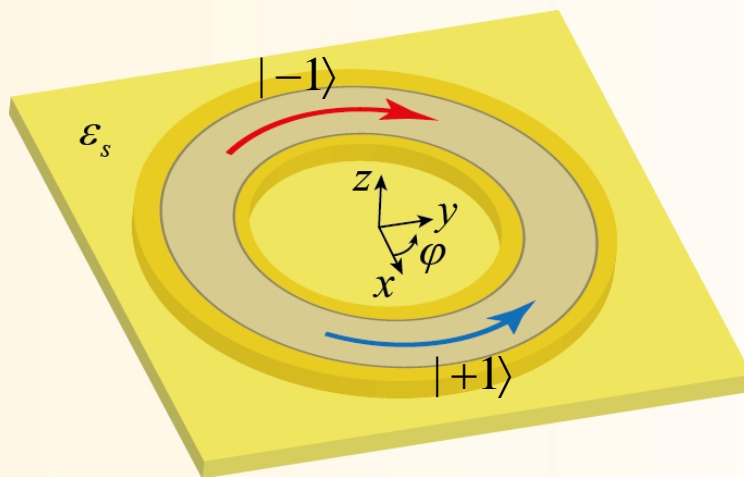
Angular Momentum



Effective Zeeman Effect



$$\Delta\varepsilon(\varphi, t) = \Delta\varepsilon_m \cos(\omega_m t - L_m \varphi)$$

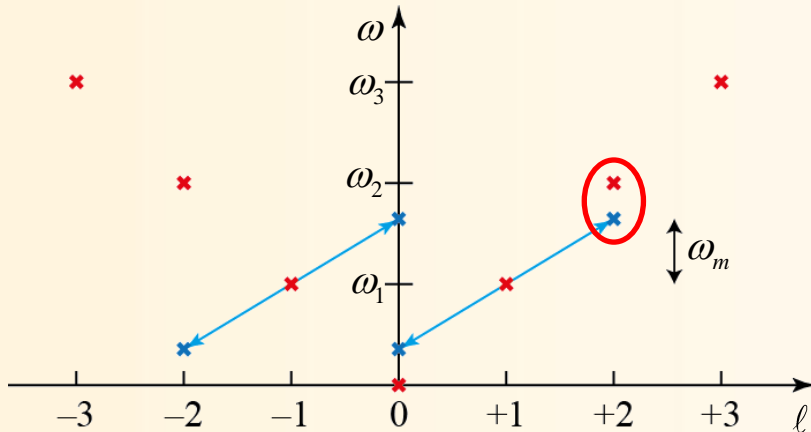


- Strong response & small footprint
- Microwaves to visible
- Weak modulation
- Moderate bandwidth

Sounas, D. L., Caloz, C. & Alù, A.
Nature Communications 4, No. 2407 (2013)

OPTIMAL CHOICE OF AZIMUTHAL SPATIO-TEMPORAL MODULATION

$$\Delta\epsilon(\varphi, t) = \Delta\epsilon_m \cos(\omega_m t - 1\varphi)$$

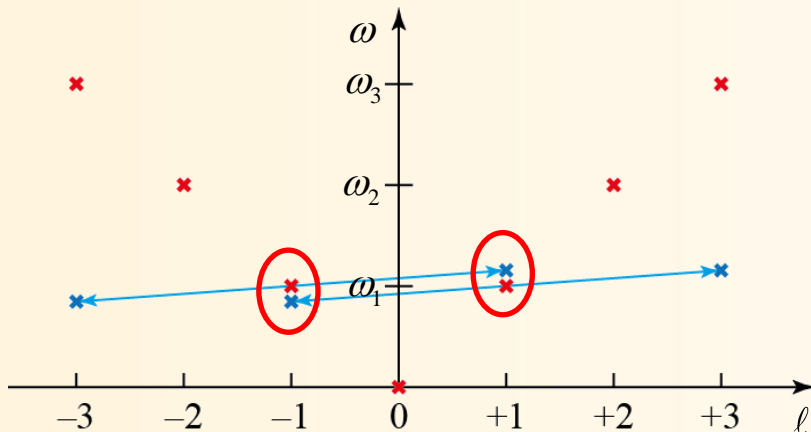


$$\omega_m = \omega_1$$



- Doable in microwaves
- Impossible in optics!**

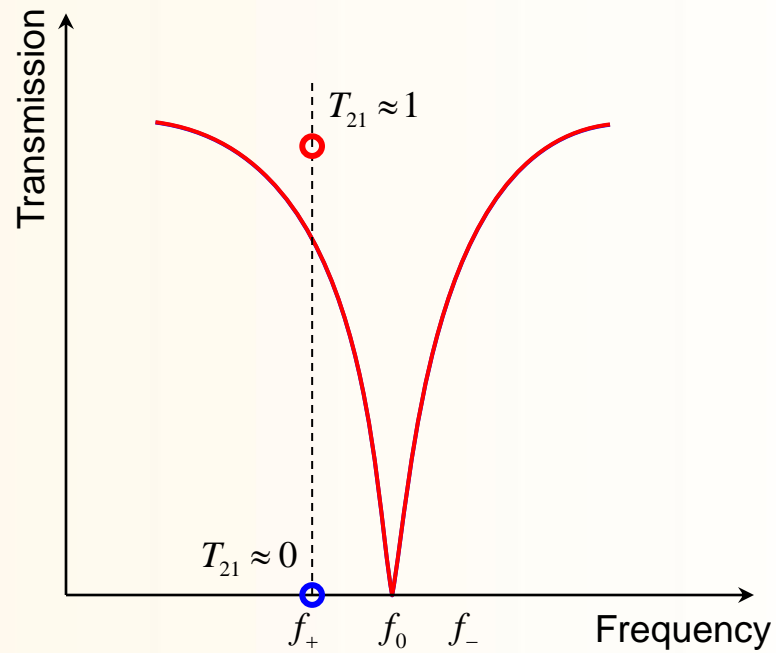
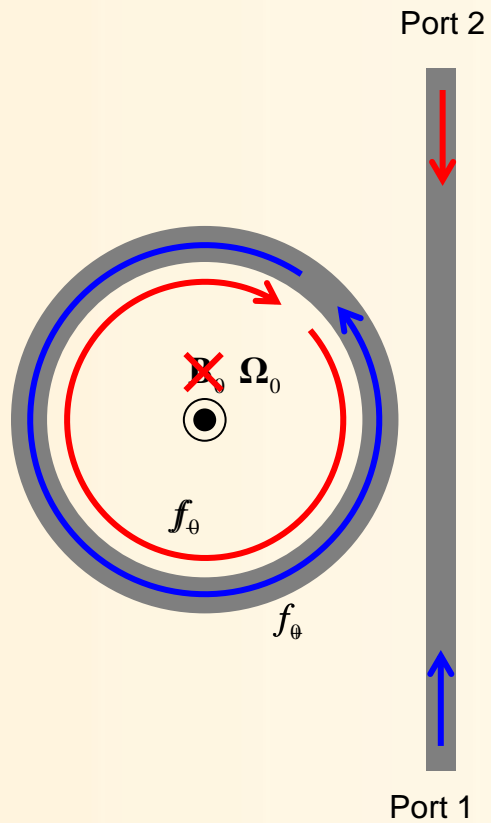
$$\Delta\epsilon(\varphi, t) = \Delta\epsilon_m \cos(\omega_m t - 2\varphi)$$



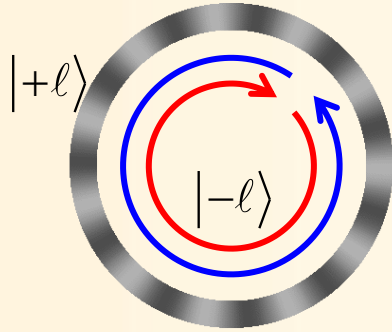
$$\omega_m \ll 0$$

- Almost zero modulation can produce strong non-reciprocity!
- Applicable over the entire EM spectrum**

OPTICAL ISOLATOR PRINCIPLE



ST-MODULATED RING – COUPLED-MODE EQUATIONS



$$\dot{a}_k = -i\omega_k a_k + i\omega_k \sum_{\ell} \kappa_{k\ell} a_{\ell} \quad \kappa_{k\ell} = \int d\mathbf{r}^3 \Delta\epsilon(\mathbf{r}, t) \mathbf{E}_k^* \cdot \mathbf{E}_{\ell}$$

$$\int d\mathbf{r}^3 \left(\epsilon |\mathbf{E}_k|^2 + \mu_0 |\mathbf{H}_k|^2 \right) = 1 \quad \Rightarrow \quad \int d\mathbf{r}^3 \epsilon |\mathbf{E}_k|^2 = \frac{1}{2}$$

$$\Delta\epsilon(\rho, \varphi, z) = \Delta\epsilon_m(\rho, z) \cos(\omega_m t - L_m \varphi)$$

$$L_m = 2\ell$$

$$\mathbf{E}_{\pm\ell}(\rho, \phi, z) = \mathbf{e}_{\ell}(\rho, z) e^{\pm i\ell\varphi} \quad \Rightarrow \quad 2\pi \int \epsilon |\mathbf{e}_{\ell}|^2 \rho d\rho dz = \frac{1}{2}$$

$$\kappa_{+l, -l} = \kappa_{-l, +l}^* = \pi e^{-i\omega_m t} \int \Delta\epsilon_m |\mathbf{e}_{\ell}|^2 \rho d\rho dz = \frac{1}{4} \frac{\int \Delta\epsilon_m |\mathbf{e}_{\ell}|^2 \rho d\rho dz}{\int \epsilon |\mathbf{e}_{\ell}|^2 \rho d\rho dz} e^{-i\omega_m t} \approx \frac{\Delta\epsilon_m}{4\epsilon} e^{-i\omega_m t}$$

$$\dot{a}_{+l} = -i\omega_{\alpha} a_{+l} + i \frac{1}{2} \omega_{\ell} \kappa_{\ell} a_{-l} e^{-i\omega_m t}$$

$$\dot{a}_{-l} = -i\omega_{\beta} a_{-l} + i \frac{1}{2} \omega_{\ell} \kappa_{\ell} a_{+l} e^{i\omega_m t}$$

$$\kappa_{\ell} = 2\kappa_{+l, -l} e^{i\omega_m t}$$

$$|\alpha\rangle = | +l \rangle e^{-i\omega_{\alpha} t} + \frac{\Delta\omega}{\omega_{\ell} \kappa_{\ell}} | -l \rangle e^{-i(\omega_{\alpha} - \omega_m) t}$$

$$\omega_{\alpha} = \omega_{\ell} - \Delta\omega/2$$

$$\omega_{\beta} = \omega_{\ell} + \Delta\omega/2$$

$$|\beta\rangle = | -l \rangle e^{-i\omega_{\beta} t} - \frac{\Delta\omega}{\omega_{\ell} \kappa_{\ell}} | +l \rangle e^{-i(\omega_{\beta} + \omega_m) t}$$

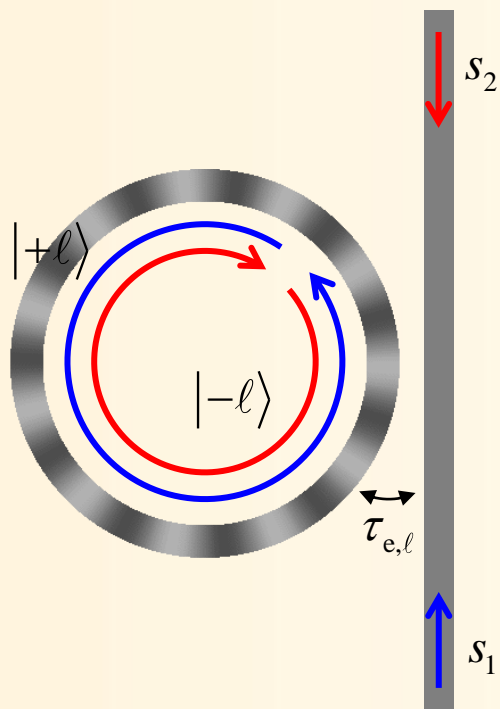
$$\Delta\omega = \sqrt{\omega_m^2 + \omega_{\ell}^2 \kappa_{\ell}^2} - \omega_m$$

ST-MODULATED RING – LOSS + EXCITATION

Loss $\rightarrow -i\omega_\ell \rightarrow -i\omega_\ell - \tau_\ell^{-1} \rightarrow$

$$\dot{a}_{+l} = (-i\omega_\ell - \tau_\ell^{-1})a_{+l} + i\frac{1}{2}\omega_\ell\kappa_\ell a_{-l}e^{-i\omega_m t}$$

$$\dot{a}_{-l} = (-i\omega_\ell - \tau_\ell^{-1})a_{-l} + i\frac{1}{2}\omega_\ell\kappa_\ell a_{+l}e^{i\omega_m t}$$

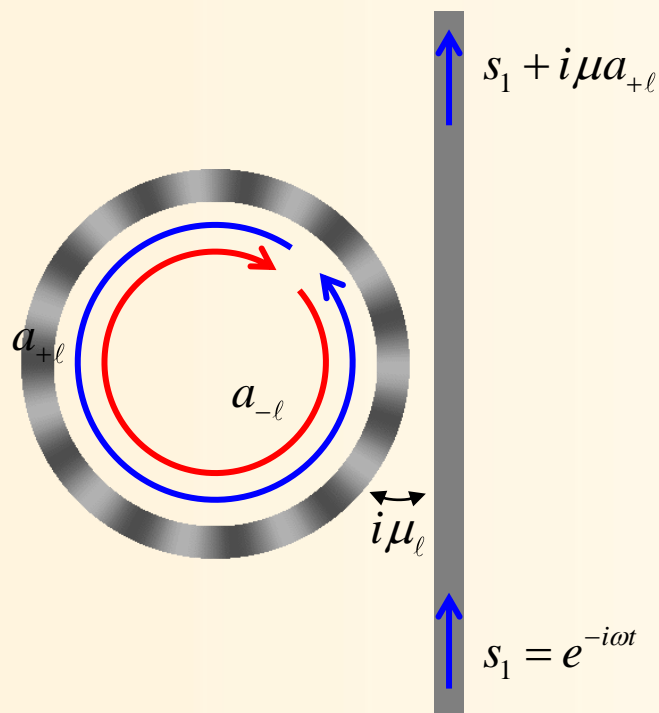


$$\dot{a}_{+l} = (-i\omega_\ell - \tau_\ell^{-1})a_{+l} + i\frac{1}{2}\omega_\ell\kappa_\ell a_{-l}e^{-i\omega_m t} + i\mu_\ell s_1$$

$$\dot{a}_{-l} = (-i\omega_\ell - \tau_\ell^{-1})a_{-l} + i\frac{1}{2}\omega_\ell\kappa_\ell a_{+l}e^{i\omega_m t} + i\mu_\ell s_2$$

$$\mu_\ell^2 = \frac{2}{\tau_{e,l}}$$

TRANSMISSION COEFFICIENTS



$$\dot{a}_{+l} = (-i\omega_l - \tau_l^{-1})a_{+l} + i\frac{1}{2}\omega_l\kappa_l a_{-l} e^{-i\omega_m t} + i\mu_l s_1$$

$$\dot{a}_{-l} = (-i\omega_l - \tau_l^{-1})a_{-l} + i\frac{1}{2}\omega_l\kappa_l a_{+l} e^{i\omega_m t}$$



$$a_{+l} = -\frac{\mu_l(\omega - \omega_m - \omega_l + i\tau_l^{-1})}{(\omega - \omega_l + i\tau_l^{-1})(\omega - \omega_m - \omega_l + i\tau_l^{-1}) - \omega_l^2\kappa_l^2/4} e^{i\omega t}$$



$$T_{21} = \frac{(\omega - \omega_l + i\tau_l^{-1} - i2\tau_{e,l}^{-1})(\omega - \omega_m - \omega_l + i\tau_l^{-1}) - \omega_l^2\kappa_l^2/4}{(\omega - \omega_l + i\tau_l^{-1})(\omega - \omega_m - \omega_l + i\tau_l^{-1}) - \omega_l^2\kappa_l^2/4}$$

ZERO TRANSMISSION CONDITION

$$(\omega - \omega_\ell)(\omega - \omega_m - \omega_\ell) - \tau_\ell^{-1}(\tau_\ell^{-1} - 2\tau_{e,\ell}^{-1}) - \omega_\ell^2 \kappa_\ell^2 / 4 = 0$$



$$(\omega - \omega_\ell)(\omega - \omega_m - \omega_\ell) - \omega_\ell^2 \kappa_\ell^2 / 4 \cong 0$$



$$T_{21} = 0$$



$$\omega = \omega_\alpha = \omega_\ell - \Delta\omega/2$$

$$\omega = \omega_\beta + \omega_m = \omega_\ell + \Delta\omega/2 + \omega_m$$

$$\tau_\ell^{-1}(\omega - \omega_\ell) + (\tau_\ell^{-1} - 2\tau_{e,\ell}^{-1})(\omega - \omega_m - \omega_\ell) = 0$$



$$\omega = \omega_\alpha$$



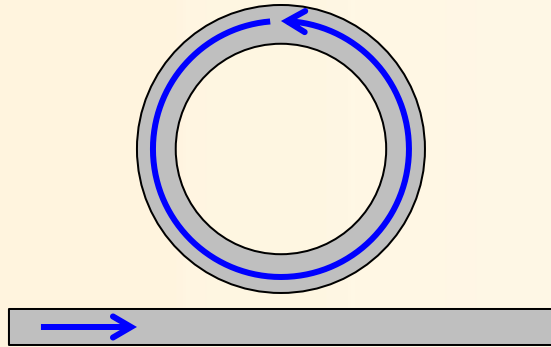
$$\omega = \omega_\beta$$

$$\tau_\ell = \tau_{e,\ell} \frac{\Delta\omega + \omega_m}{\Delta\omega + 2\omega_m}$$

No solution



BANDWIDTH REQUIREMENTS

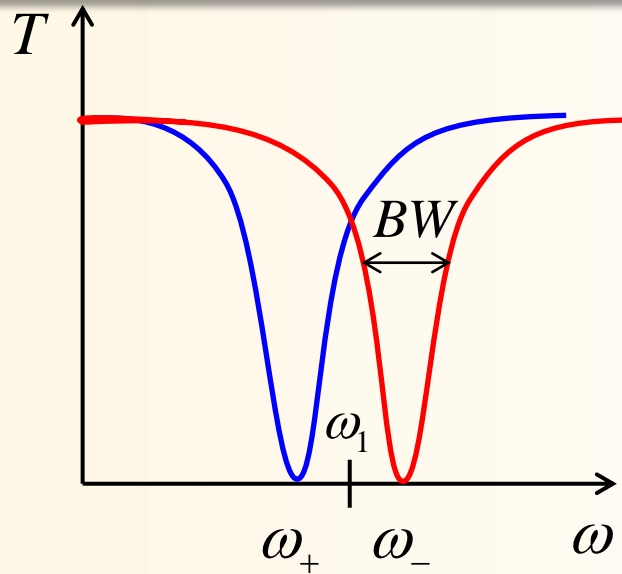


$$\Delta\omega = \sqrt{\omega_m^2 + \omega_1^2 C_m^2} - \omega_m$$

$$\Delta\omega_{\max} = \frac{\omega_1 C_m}{\sqrt{3}}$$

$$C_m \propto \Delta\epsilon_m$$

$$Q = \frac{\omega_1}{2BW}$$

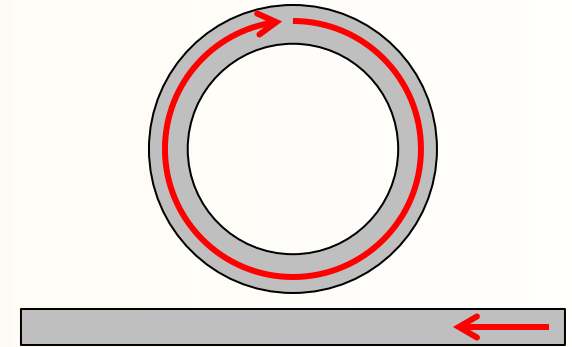


$$\Delta\omega = \omega_- - \omega_+ = BW$$

$$\frac{\omega_1 C_m}{\sqrt{3}} = BW$$

$$\Delta\epsilon_m \propto \frac{BW\sqrt{3}}{\omega_1} \propto \frac{1}{Q}$$

$$Q\Delta\epsilon_m \propto 1$$



High Q allows
weak modulation



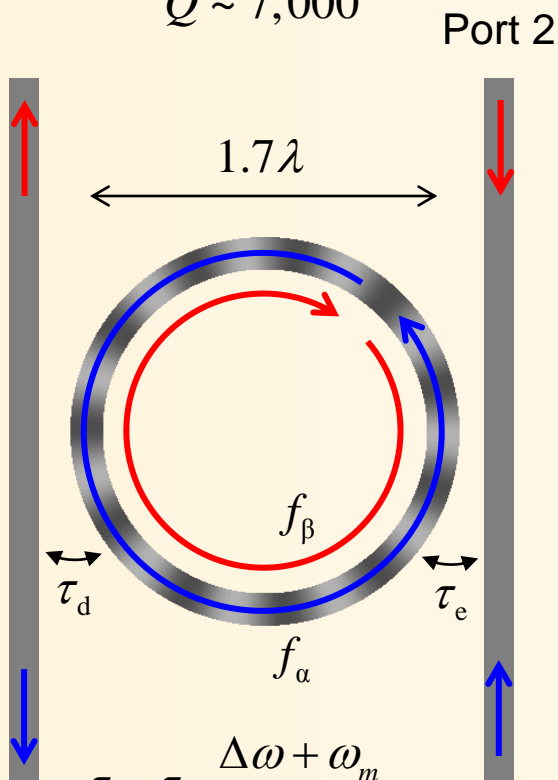
**Ideal for RF as well as
optics!**

ANGULAR-MOMENTUM-BIASED OPTICAL ISOLATOR

$$\Delta\varepsilon_m = 5 \times 10^{-4} \varepsilon !$$

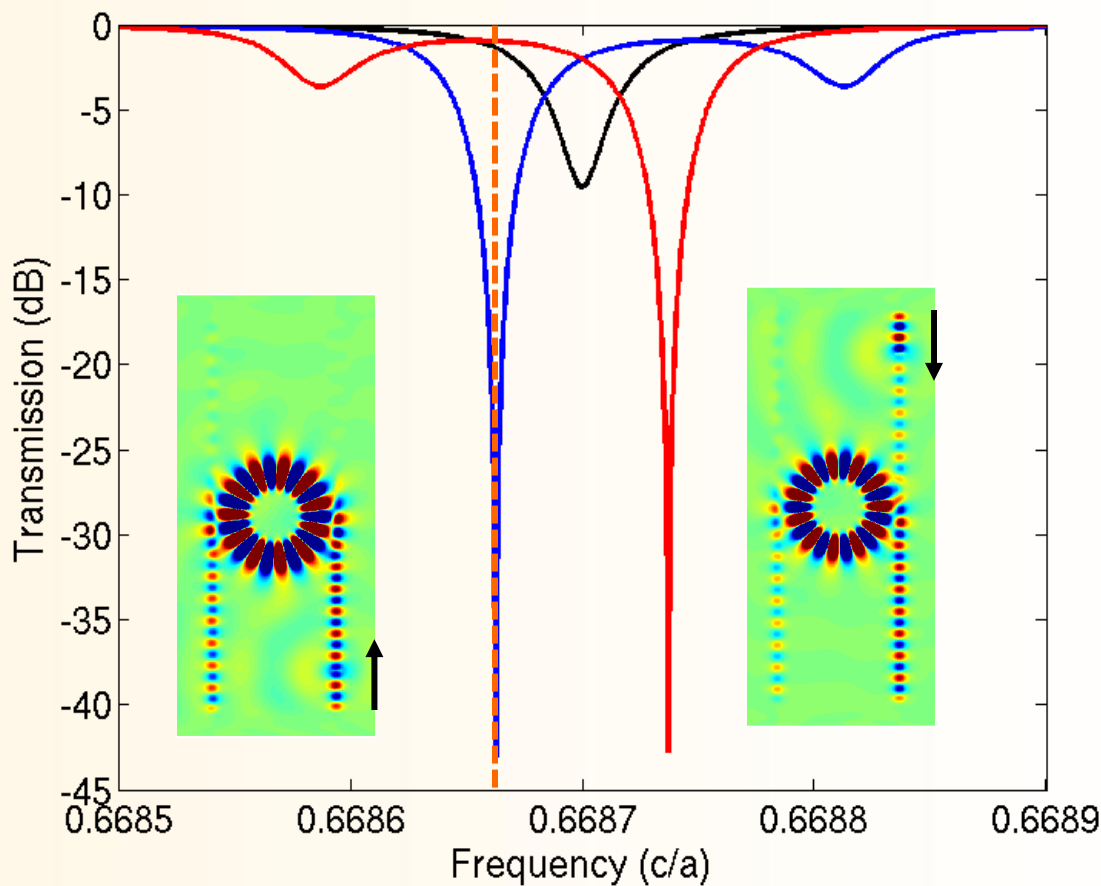
$$Q \sim 7,000$$

$$Q\Delta\varepsilon_m \square 1$$

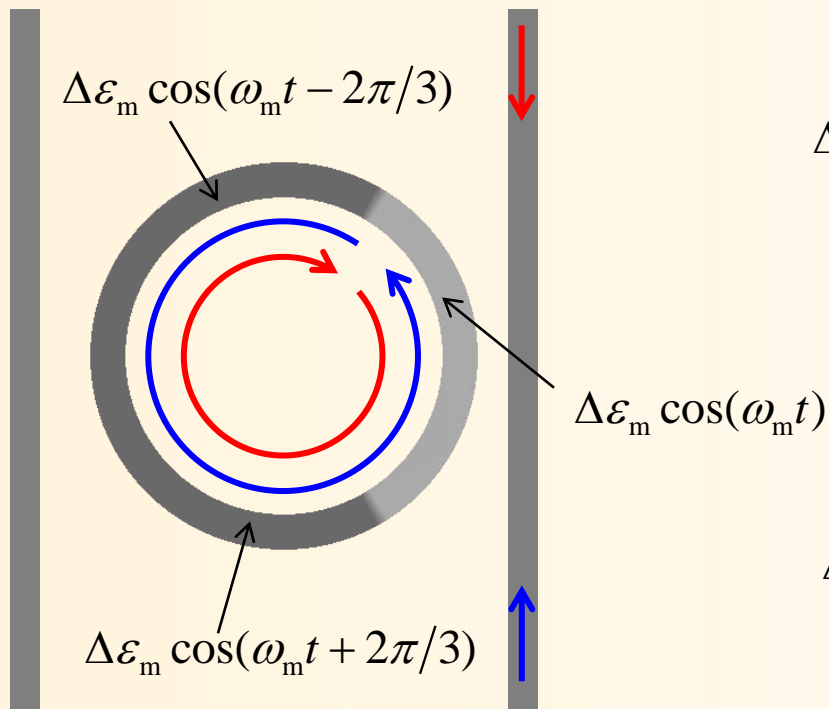


$$\tau = \tau_e \frac{\Delta\omega + \omega_m}{\Delta\omega + 2\omega_m}$$

$$\tau_d = \tau_e \left(1 + \frac{\Delta\omega}{\omega_m} \right)$$



ST MODULATION DISCRETIZATION: 3 STEPS



$$\begin{aligned}\Delta \varepsilon &= \Delta \varepsilon_m \sum_{n=0}^2 \cos[\omega_m t - 2(n-1)\pi/3] \text{rect}\left(\frac{\varphi - \varphi_n}{\Delta \varphi}\right) \\ &= \Delta \varepsilon_m \sum_{k=-\infty}^{\infty} \text{sinc}\left(\frac{3N+1}{3}\right) \cos[\omega_m t - (3k+1)\varphi]\end{aligned}$$



Poisson Summation Formula

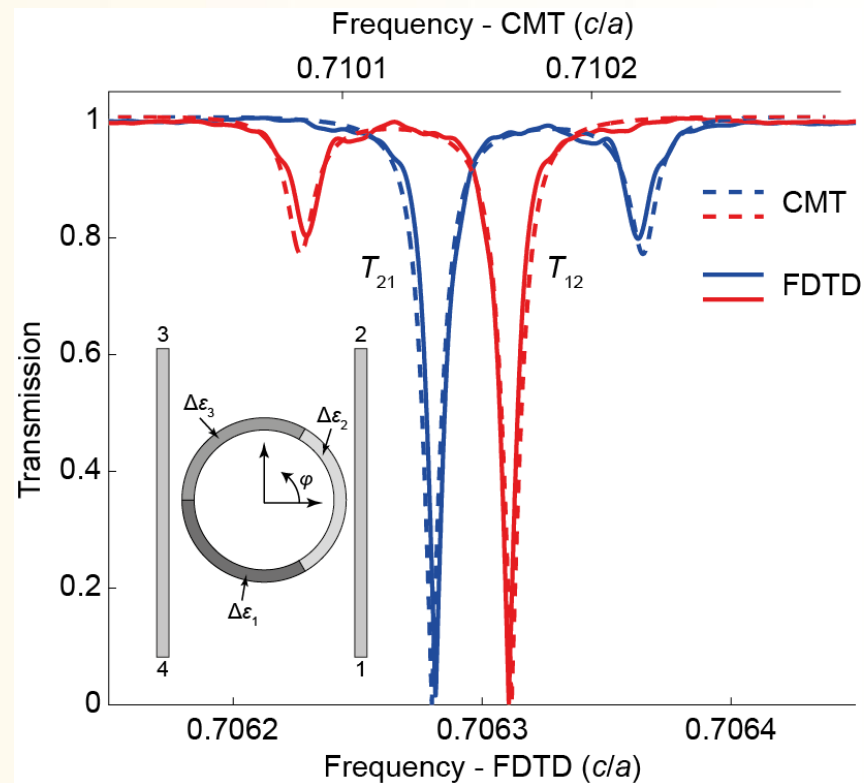
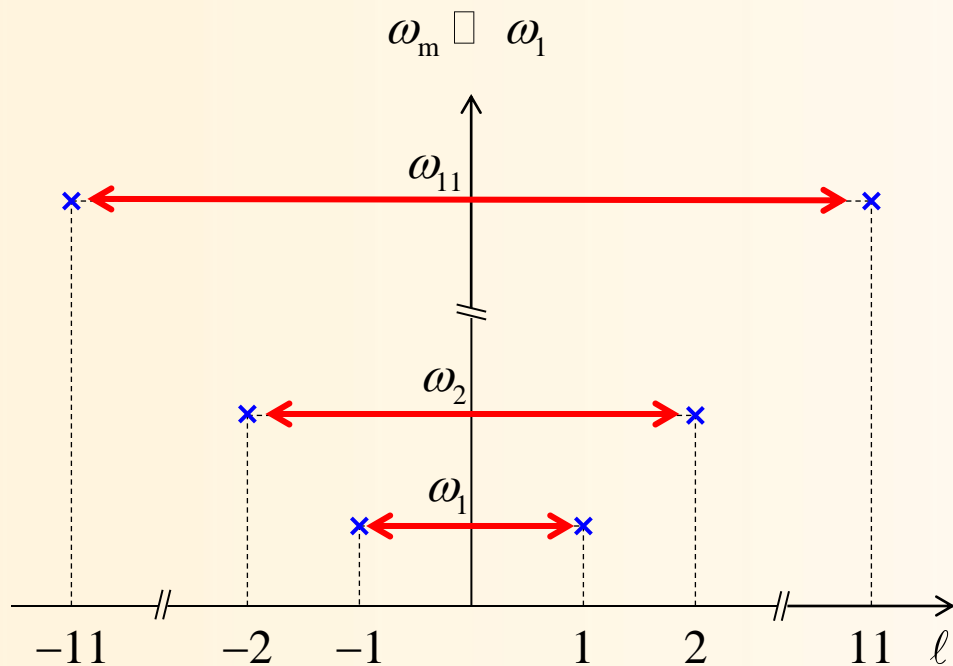
$$\Delta \varepsilon = \Delta \varepsilon_m \sum_{k=-\infty}^{\infty} \text{sinc}\left(\frac{3N+1}{3}\right) \cos[\omega_m t - (3k+1)\varphi]$$

$$\begin{aligned}\Delta \varepsilon &= \Delta \varepsilon_m \text{sinc}\left(\frac{1}{3}\right) \cos(\omega_m t - \varphi) + \Delta \varepsilon_m \text{sinc}\left(\frac{4}{3}\right) \cos(\omega_m t - 4\varphi) + \dots + \Delta \varepsilon_m \text{sinc}\left(\frac{22}{3}\right) \cos(\omega_m t - 22\varphi) + \dots \\ &+ \Delta \varepsilon_m \text{sinc}\left(\frac{2}{3}\right) \cos(\omega_m t + 2\varphi) + \Delta \varepsilon_m \text{sinc}\left(\frac{5}{3}\right) \cos(\omega_m t + 5\varphi) + \dots\end{aligned}$$

ST MODULATION DISCRETIZATION – 3 STEPS

$$\Delta\varepsilon = \Delta\varepsilon_m \operatorname{sinc}\left(\frac{1}{3}\right) \cos(\omega_m t - \varphi) + \Delta\varepsilon_m \operatorname{sinc}\left(\frac{4}{3}\right) \cos(\omega_m t - 4\varphi) + \dots + \Delta\varepsilon_m \operatorname{sinc}\left(\frac{22}{3}\right) \cos(\omega_m t - 22\varphi) + \dots$$

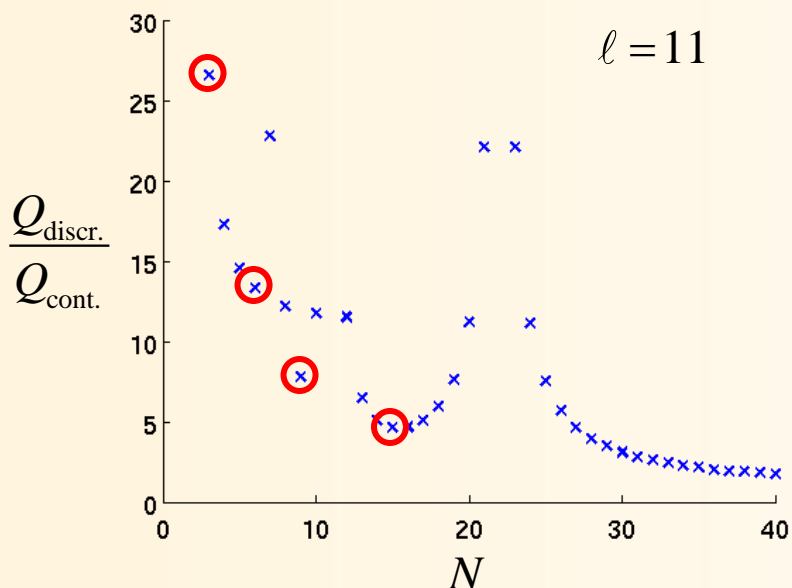
$$+ \Delta\varepsilon_m \operatorname{sinc}\left(\frac{2}{3}\right) \cos(\omega_m t + 2\varphi) + \Delta\varepsilon_m \operatorname{sinc}\left(\frac{5}{3}\right) \cos(\omega_m t + 5\varphi) + \dots$$



ST MODULATION DISCRETIZATION – N STEPS

$$\begin{aligned} \Delta \varepsilon &= \Delta \varepsilon_m \sum_{n=0}^{N-1} \cos[\omega_m t - L_m \varphi_n] \operatorname{rect}\left(\frac{\varphi - \varphi_n}{\Delta \varphi}\right) \\ &= \Delta \varepsilon_m \sum_{k=-\infty}^{\infty} (-1)^{k(N-1)} \operatorname{sinc}\left(\frac{L_m + kN}{N}\right) \cos[\omega_m t - \underbrace{(L_m + kN)\varphi}_{L_{m,\text{eff}}}] \end{aligned}$$

ℓ -th pair $\Rightarrow |L_{m,\text{eff}}| = |L_m + kN| = 2\ell \Rightarrow \Delta \varepsilon_{\text{eff}} = \Delta \varepsilon_m \operatorname{sinc}\left(\frac{2\ell}{N}\right) \Rightarrow Q_{\text{discr.}} = \frac{Q_{\text{cont.}}}{\operatorname{sinc}\left(\frac{2\ell}{N}\right)}$



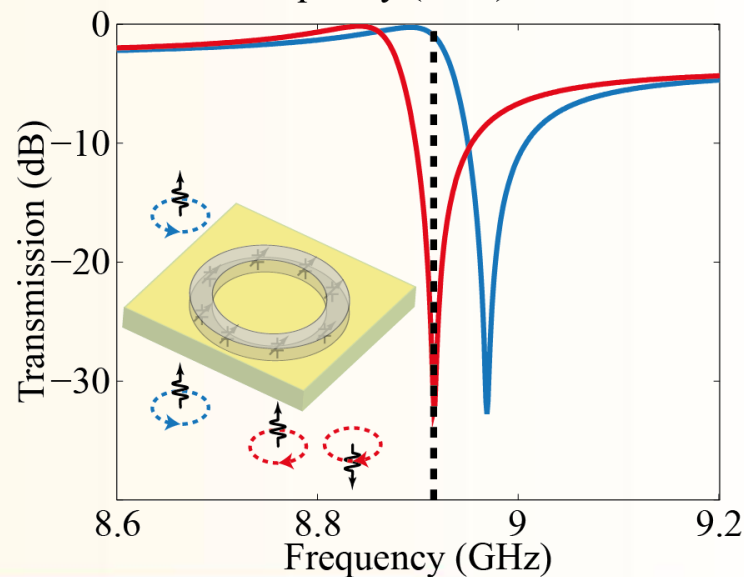
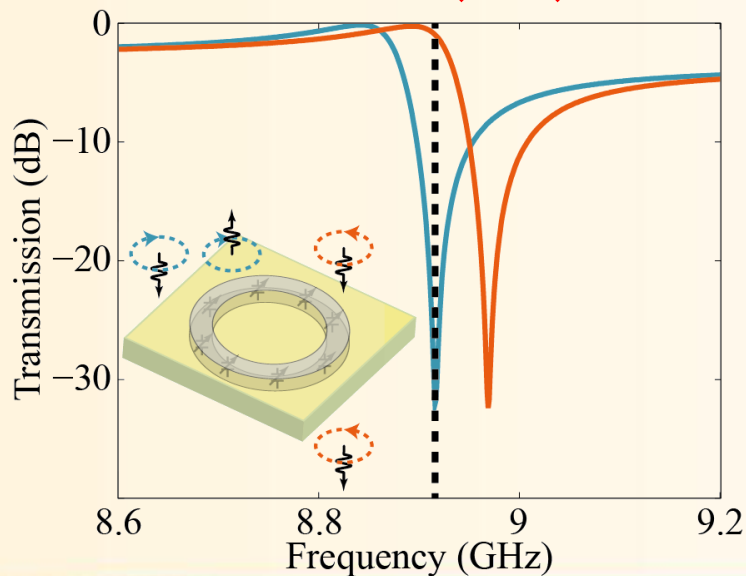
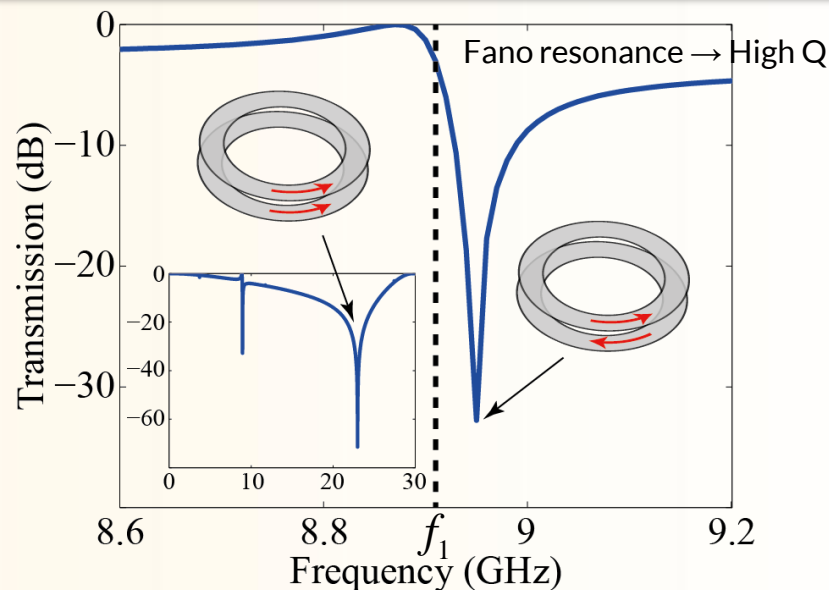
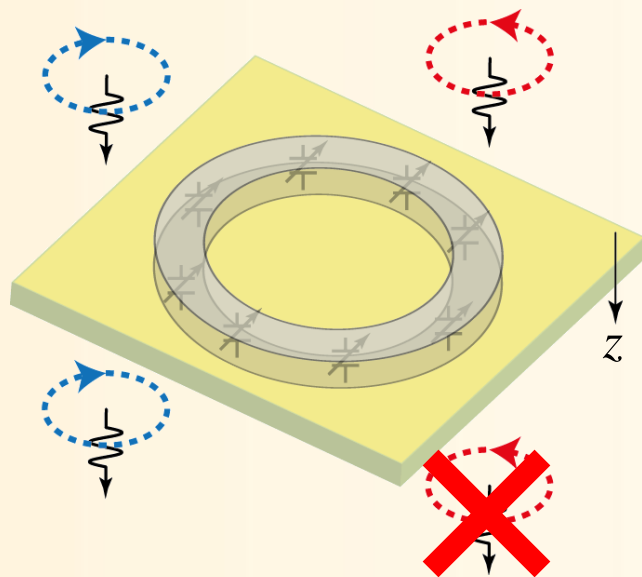
$$N = 3 \quad Q_{\text{discr.}} \approx 26 Q_{\text{cont.}}$$

$$N = 6 \quad Q_{\text{discr.}} \approx 14 Q_{\text{cont.}}$$

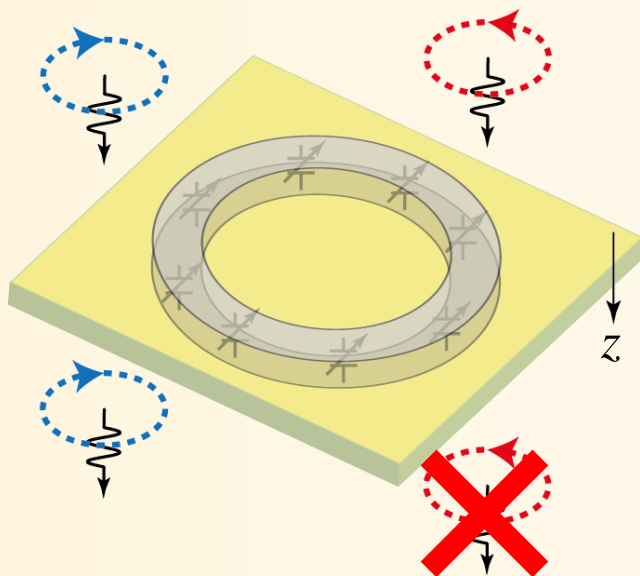
$$N = 9 \quad Q_{\text{discr.}} \approx 7 Q_{\text{cont.}}$$

$$N = 15 \quad Q_{\text{discr.}} \approx 5 Q_{\text{cont.}}$$

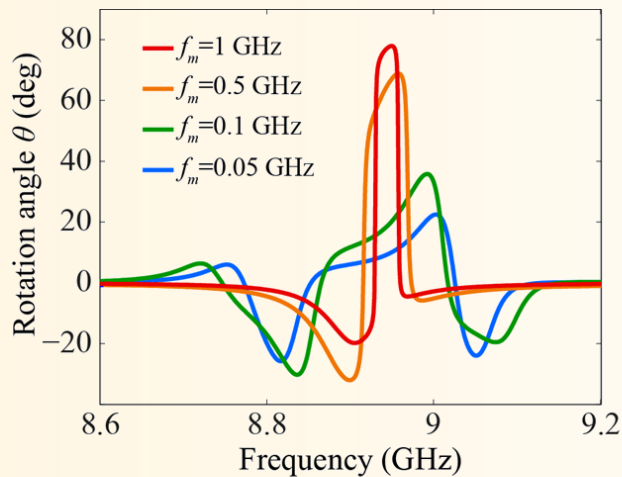
RF METASURFACE ISOLATOR FOR CIRCULARLY POLARIZED WAVES



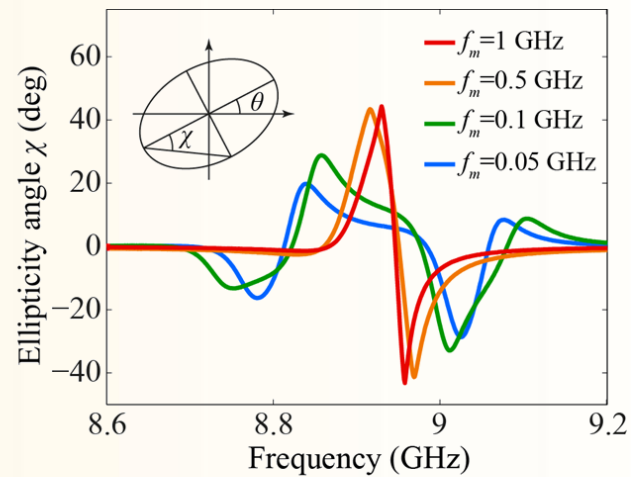
RF ISOLATOR FOR CIRCULARLY POLARIZED WAVES



6000°/wavelength !

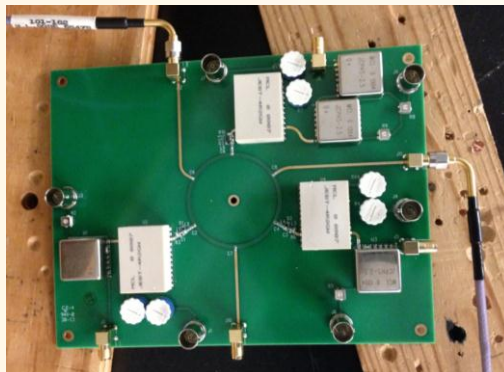
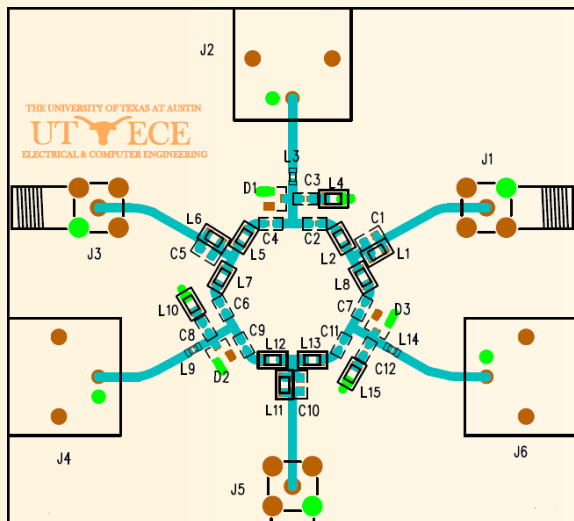


(a)

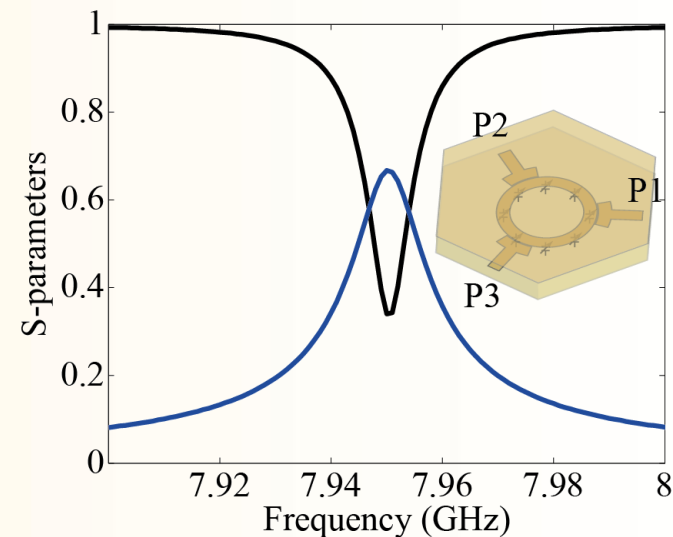


(b)

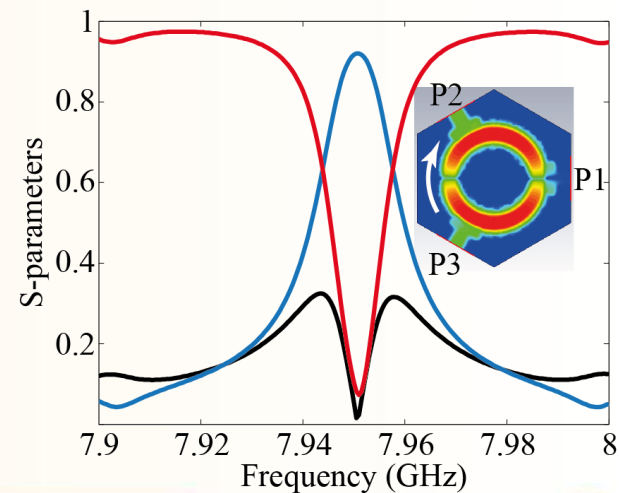
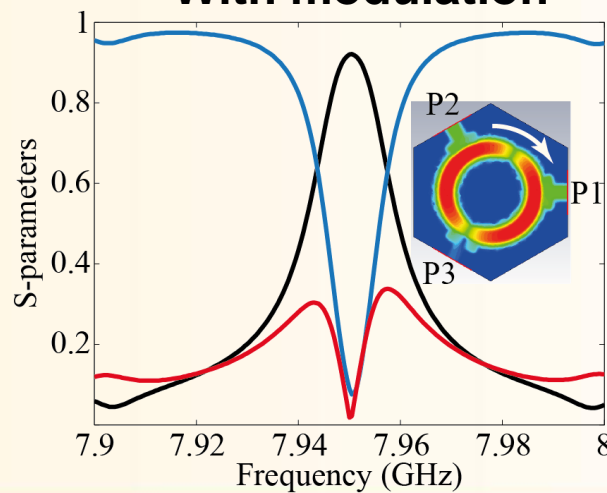
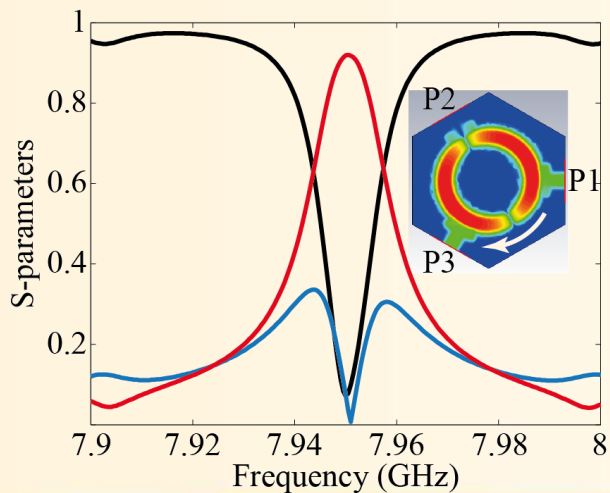
RF MAGNET-LESS INTEGRATED CIRCULATOR



Without modulation

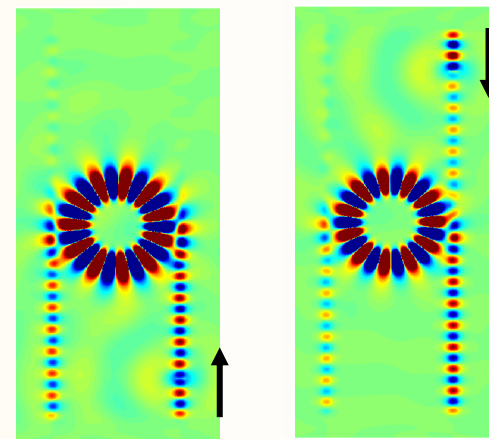
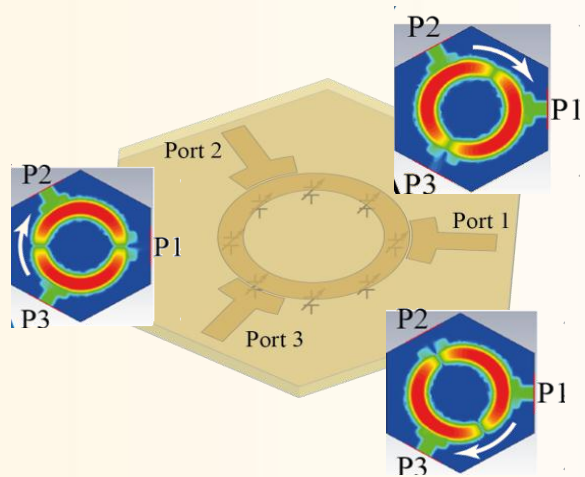
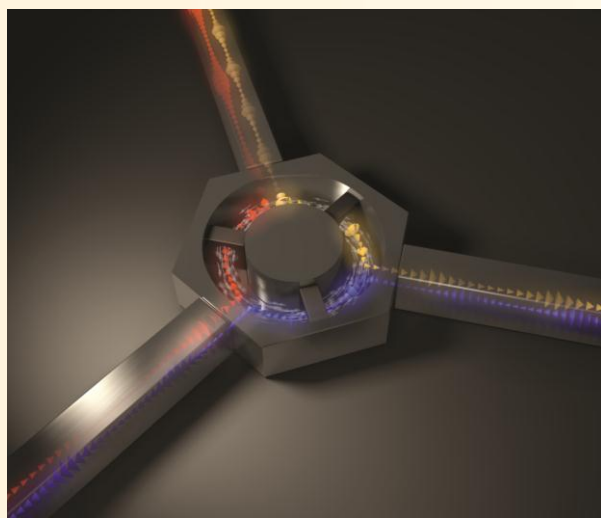


With modulation



NON-RECIPROCAL ANGULAR-MOMENTUM BIASED METAMATERIALS

Magnetic-free, linear, giant nonreciprocity at the subwavelength scale: angular-momentum biased meta-atoms, coupled to a circularly-symmetric high-Q resonance



D. L. Sounas, C. Caloz, and A. Alù, *Nature Communications* 4, No. 2407 (2013)
R. Fleury, D. L. Sounas, C. Sieck, M. Haberman, A. Alù, *Science* in press (2014)
D. L. Sounas, A. Alù, *ACS Photonics* under review (2014)