

TIME METAMATERIALS

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Physics Program, CUNY Graduate Center

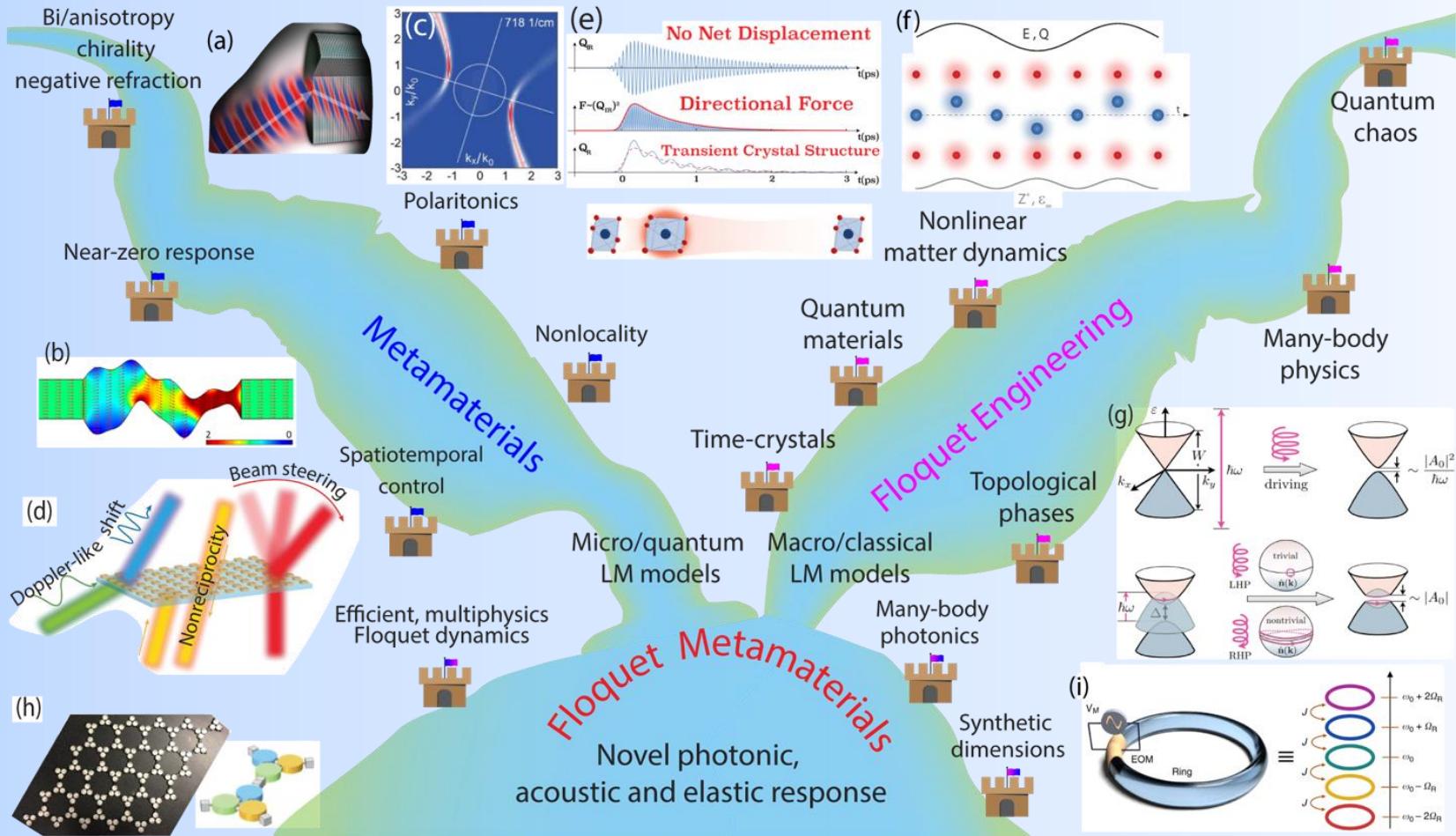
Department of Electrical Engineering, City College of New York, CUNY

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Supported by the **Air Force Office of Scientific Research (Dr. Arje Nachman)** through an **SBIR program** with **Silicon Audio, Inc.**



FLOQUET METAMATERIALS

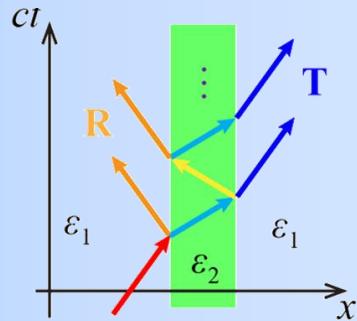


S. Yin, E. Galiffi, A. Alù, *E-Light* 2, 8 (2022)

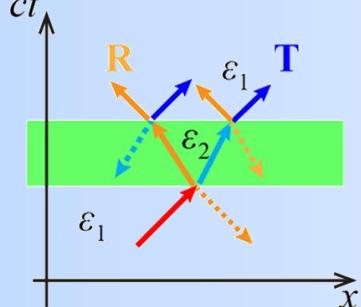
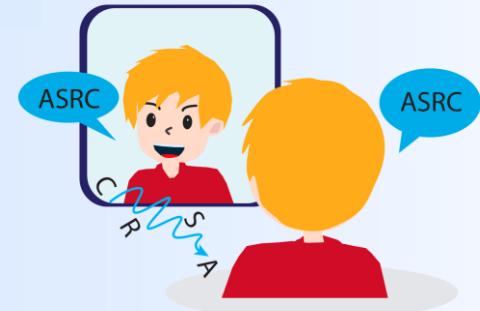


TIME AS A NEW DIMENSION FOR WAVE MANIPULATION

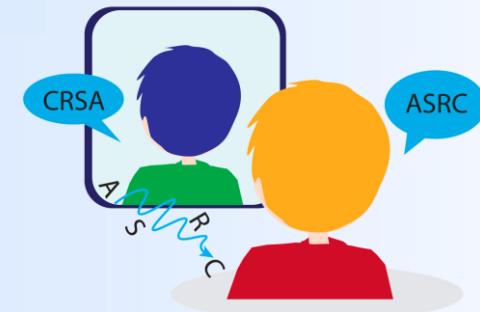
Spatial reflections



ω (frequency) is conserved
 k (momentum) is reversed



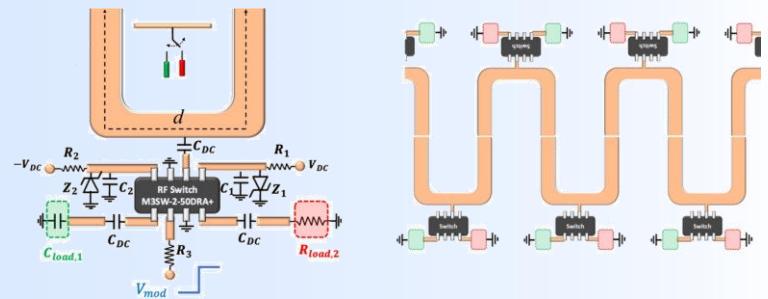
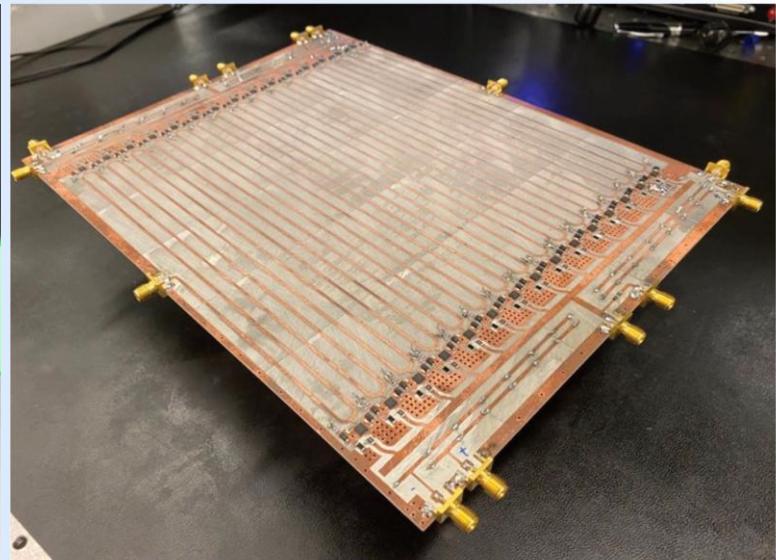
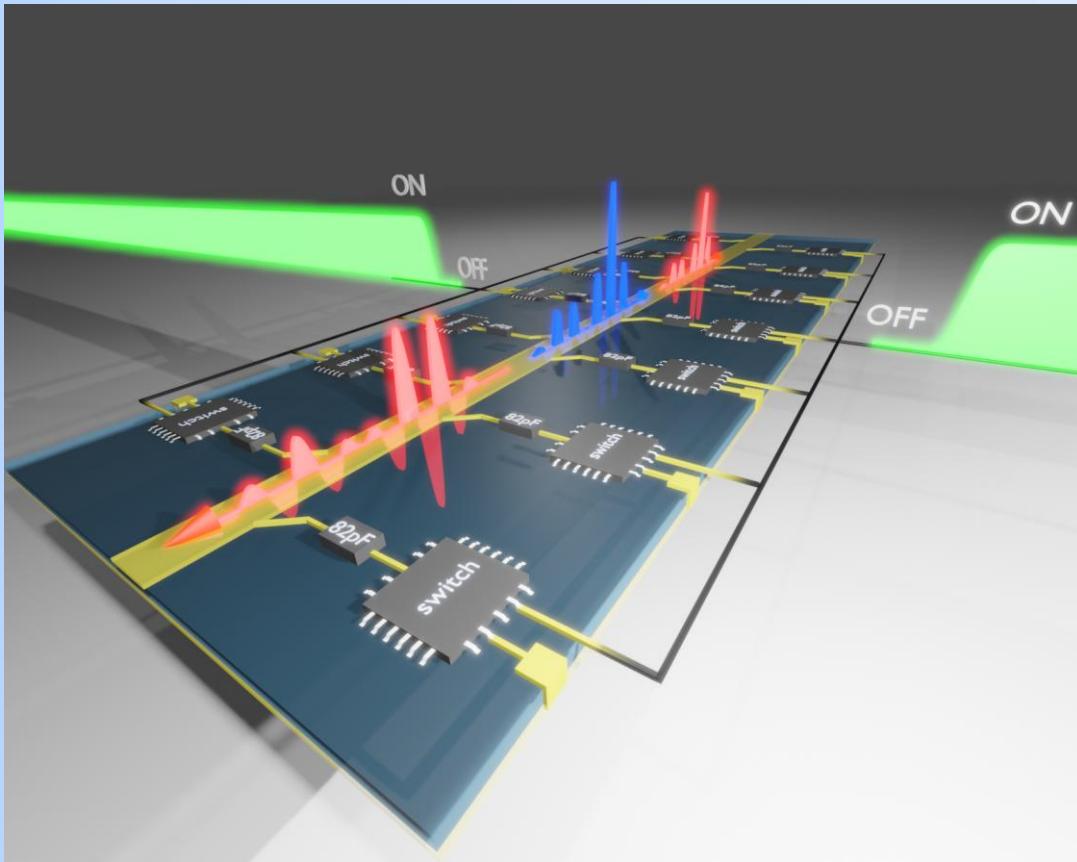
k (momentum) is conserved
 ω (frequency) changes



H. Moussa, et al., *Nature Physics* **19**, 863 (2023)



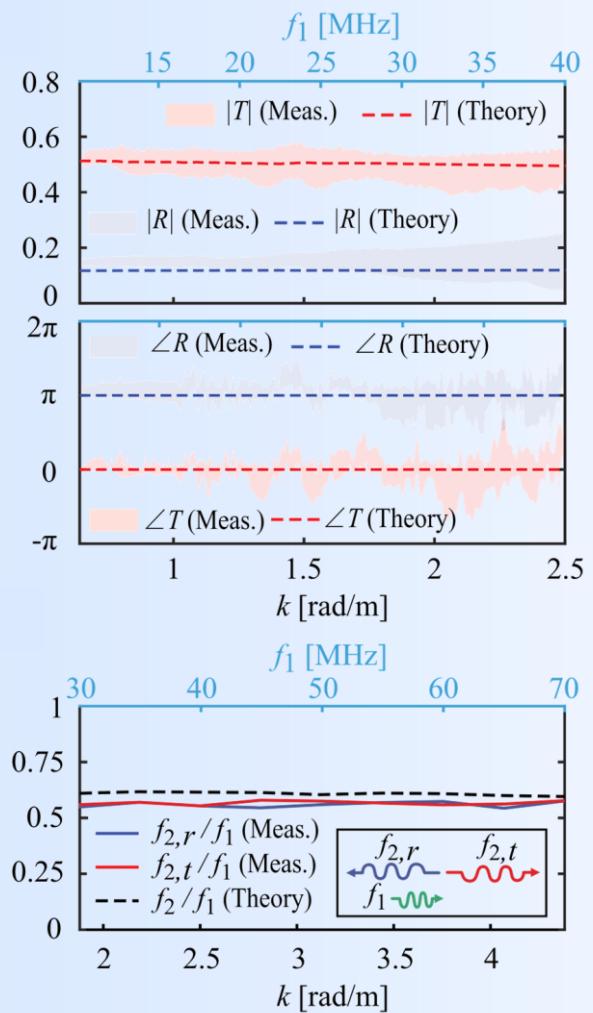
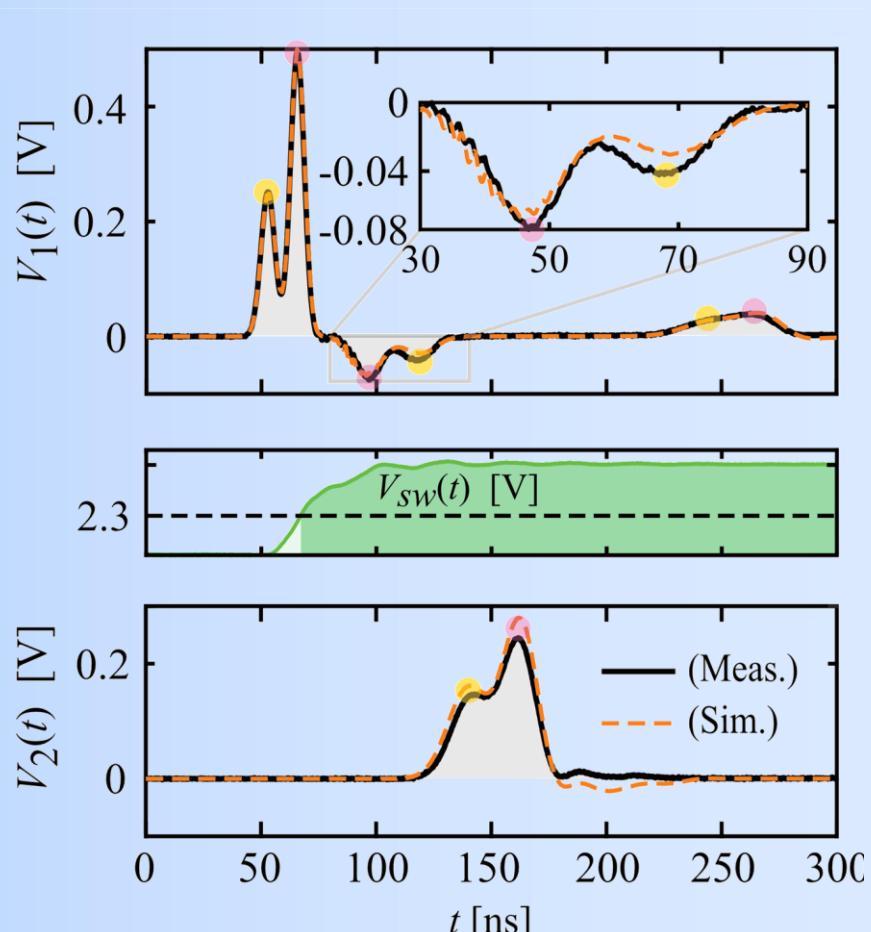
A TIME METAMATERIAL



H. Moussa, G. Xu, S. Yin, et al., *Nature Physics* **19**, 863 (2023)



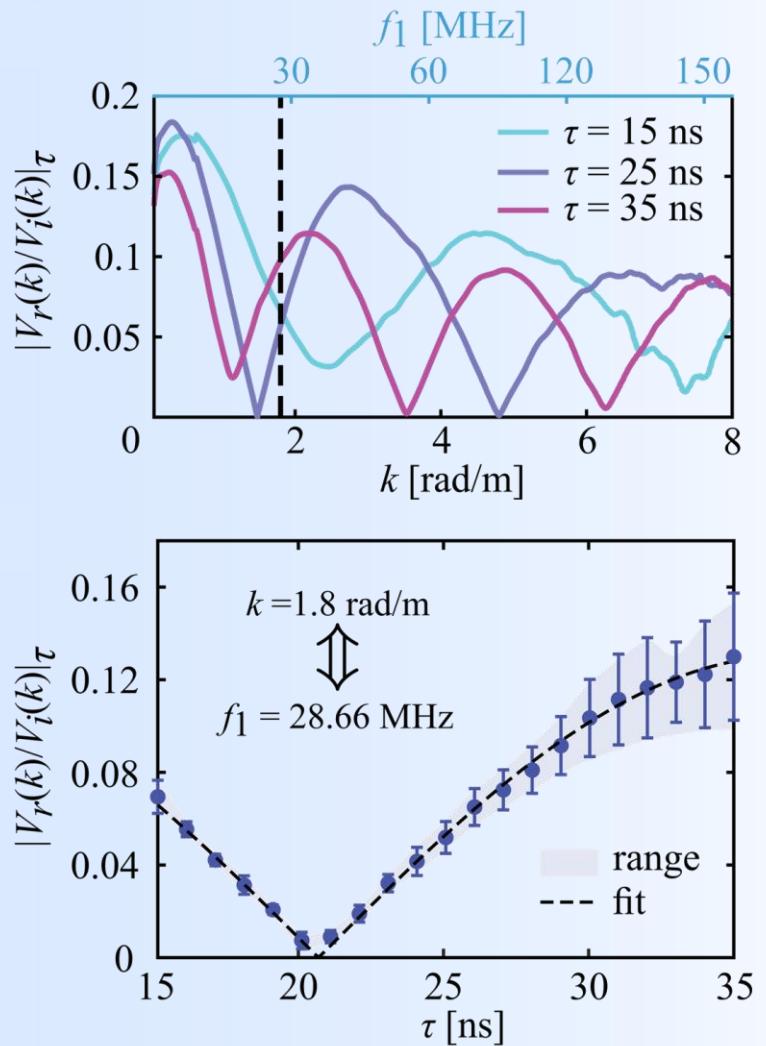
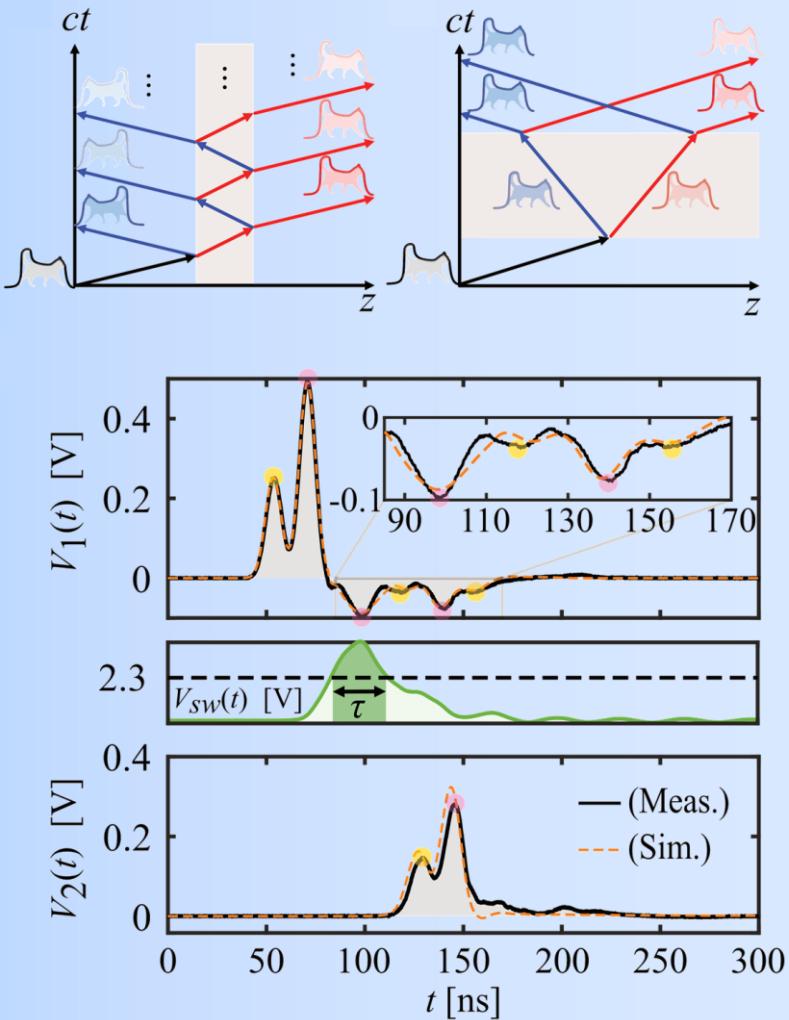
OBSERVATION OF BROADBAND TIME REVERSAL



H. Moussa, G. Xu, S. Yin, et al., *Nature Physics* **19**, 863 (2023)



SCATTERING AND INTERFERENCE WITH A TEMPORAL SLAB

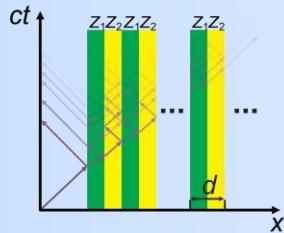


H. Moussa, et al., *Nature Physics* **19**, 863 (2023)



TIME CRYSTALS AND K-GAPS

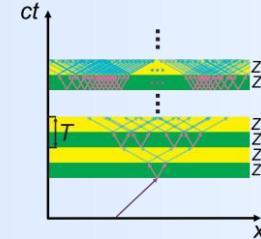
Photonic crystal



$$\psi(x + d) = \psi(x)e^{-jkd}$$

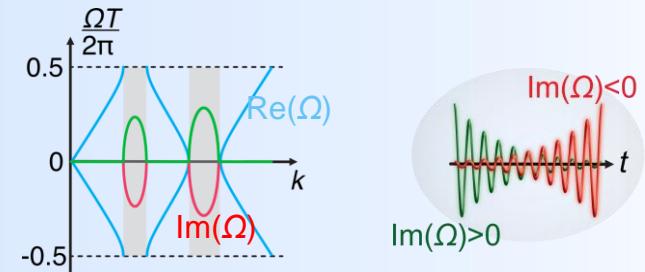
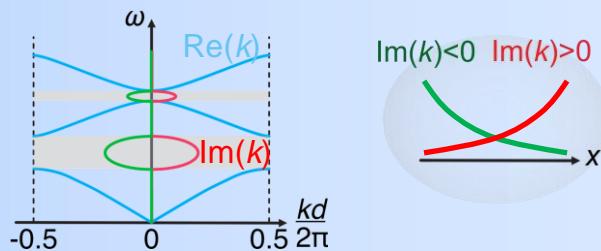
k : Bloch wavevector

Photonic time crystal

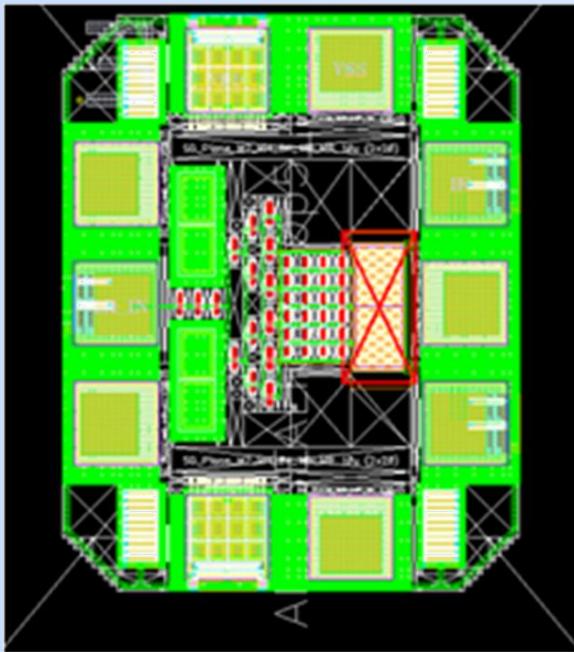


$$\psi(t + T) = \psi(t)e^{j\Omega T}$$

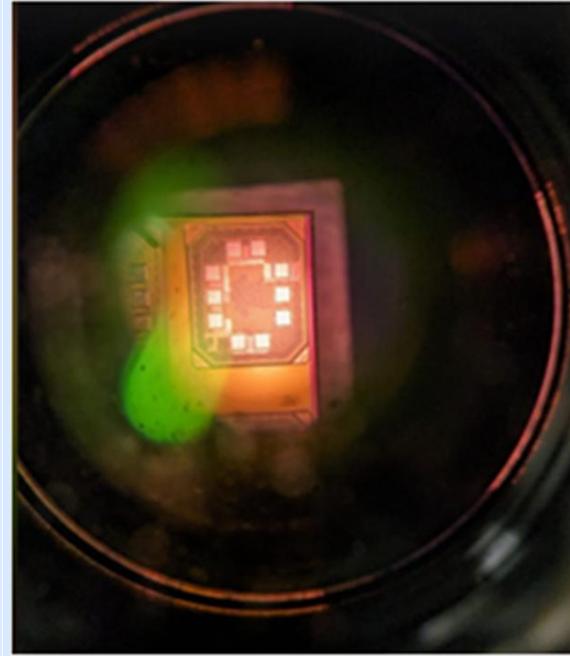
Ω : Floquet frequency



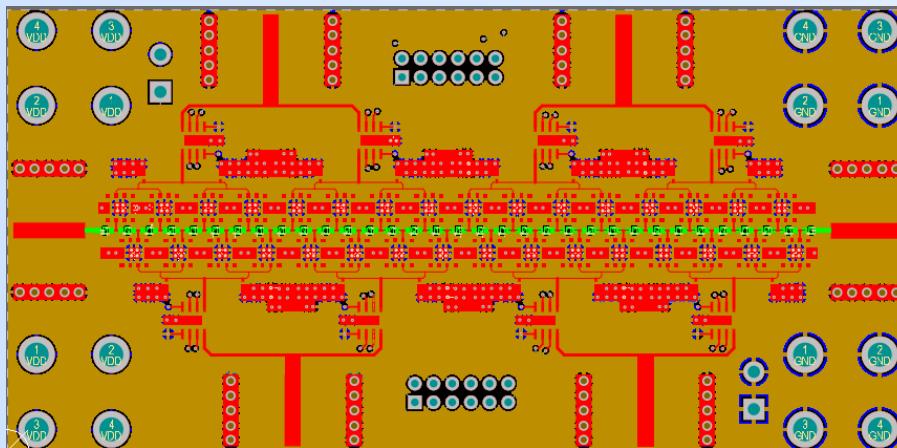
EXTENSION TO HIGHER FREQUENCIES



Taped-out chip

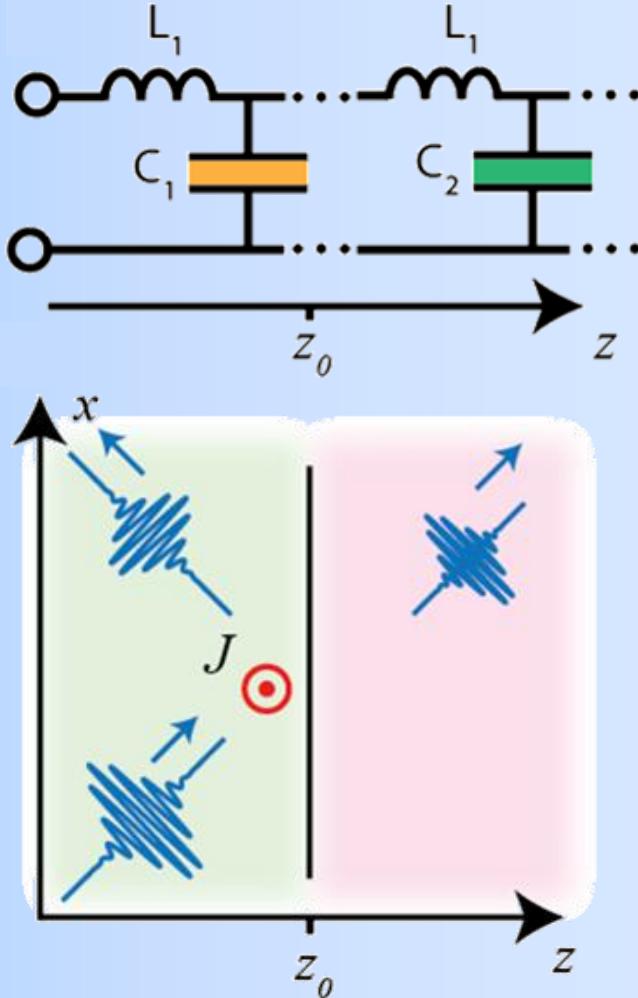


Microscopic view of the CMOS chip



with Silicon Audio

BOUNDARY CONDITIONS AND ENERGY REQUIREMENTS



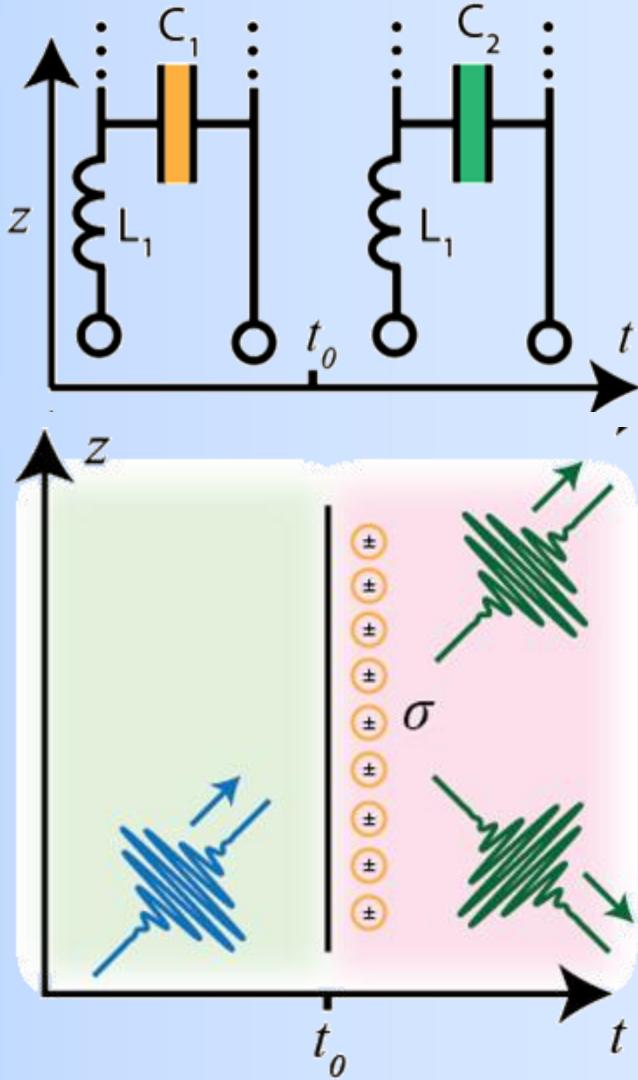
At a uniform spatial interface:

- ω is conserved
- energy is conserved
- spatial symmetry is broken (typically leading to a change in momentum)

$$E_{\tan}(z_o^-) - E_{\tan}(z_o^+) = -J_m$$

$$H_{\tan}(z_o^-) - H_{\tan}(z_o^+) = J_e$$

BOUNDARY CONDITIONS AND ENERGY REQUIREMENTS



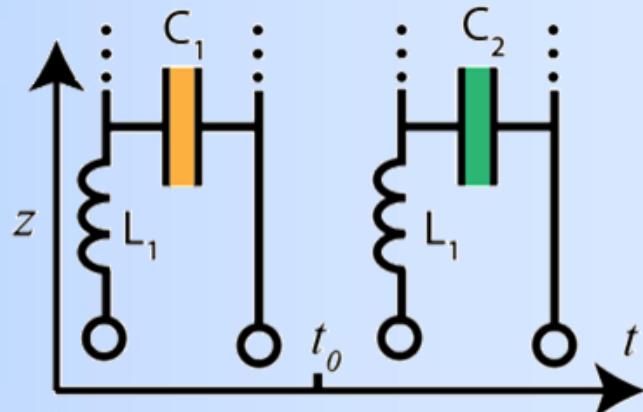
At a uniform time interface:

- ω is **not necessarily** conserved
- energy is **not necessarily** conserved
- spatial symmetry is preserved (conserving momentum)

$$D(t_o^-) - D(t_o^+) = \sigma$$

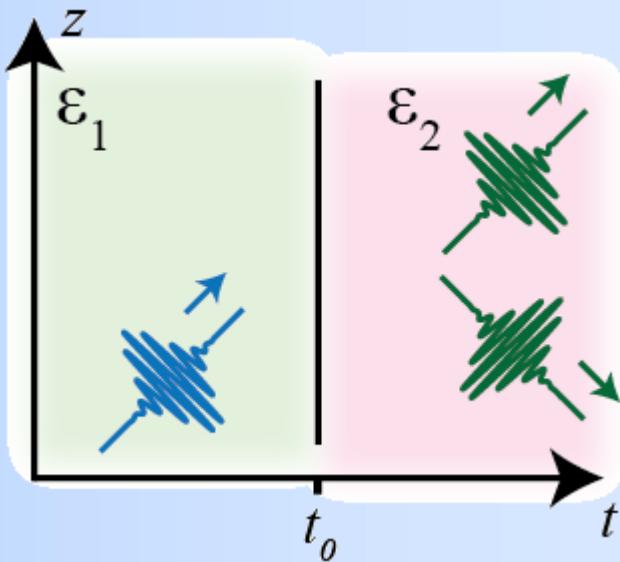
$$B(t_o^-) - B(t_o^+) = \sigma_m$$

BOUNDARY CONDITIONS AND ENERGY REQUIREMENTS



If the charge is conserved, the energy requirements may become very large:

$$P_{req} = \frac{D^2}{2\epsilon_2} - \frac{D^2}{2\epsilon_1}$$



$$\Delta U \sim N E_f \frac{\Delta \epsilon}{\epsilon_\infty}$$

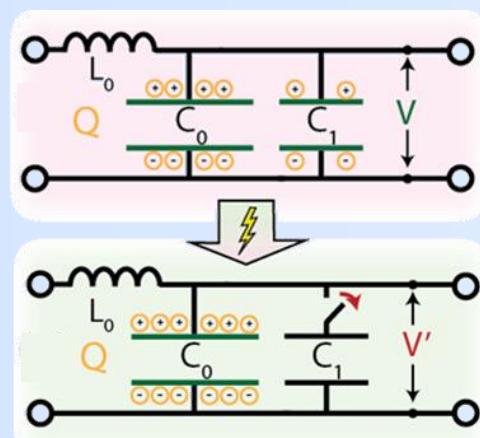
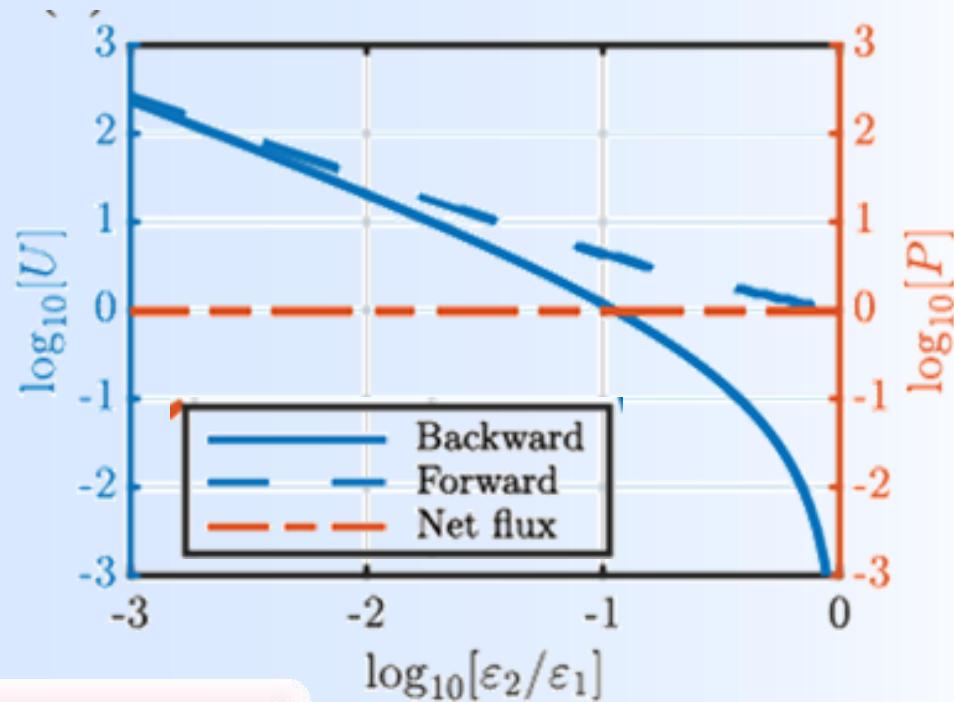
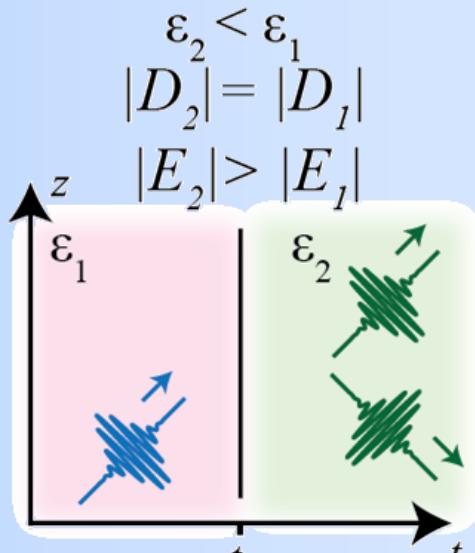
$\Delta \epsilon \sim 1$ requires $\Delta U \sim 10 \text{ J/cm}^3$!!!

Z. Hayran, J. B. Khurgin, F. Monticone. *Opt. Mater. Express, OME*. **12**, 3904–3917 (2022)

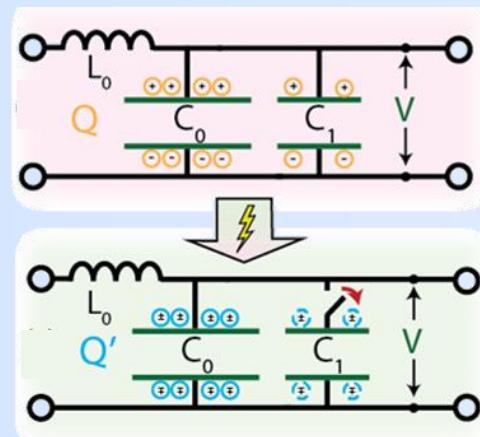
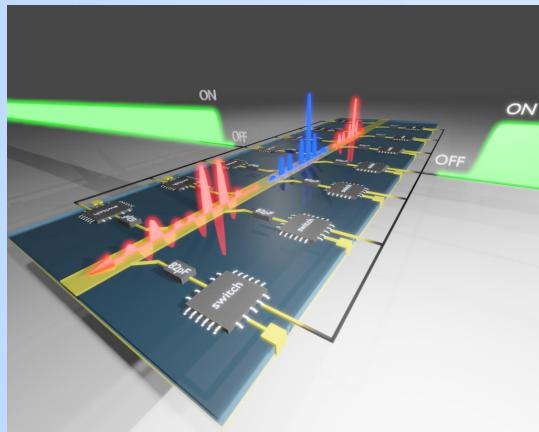
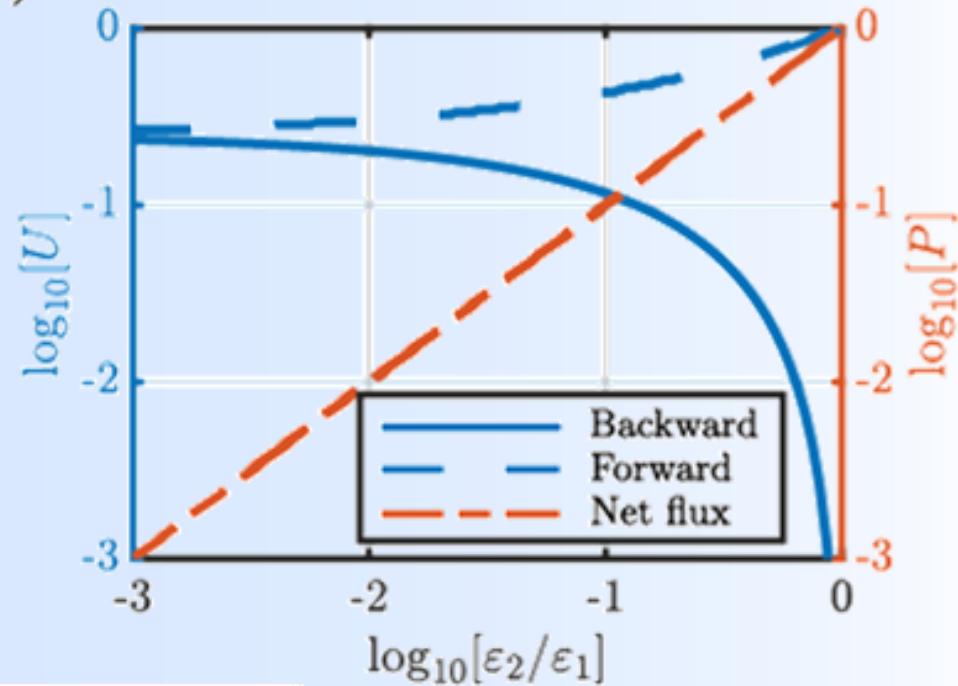
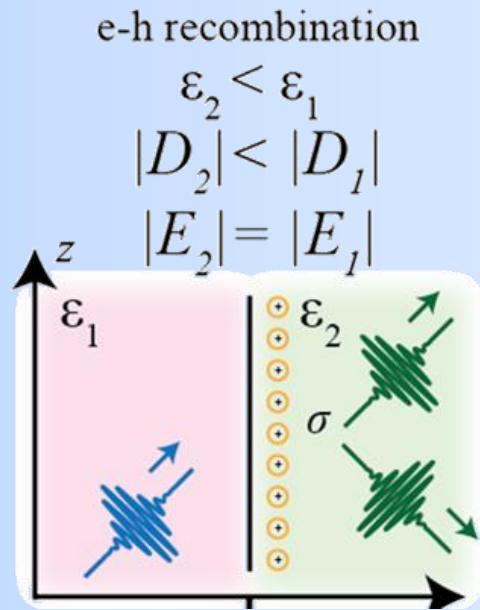


ENERGY REQUIREMENTS

Pockels/Kerr mod. (on)

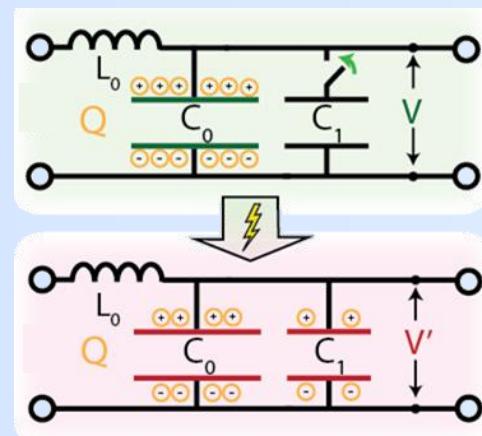
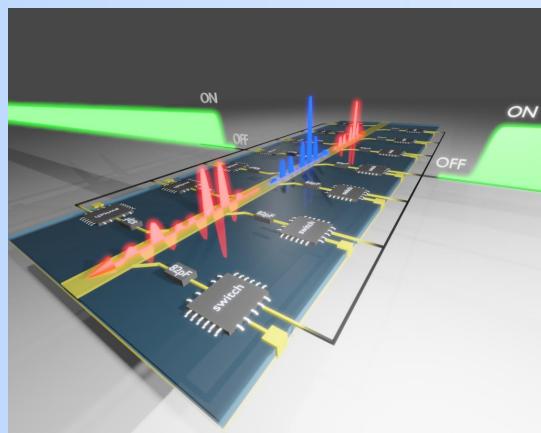
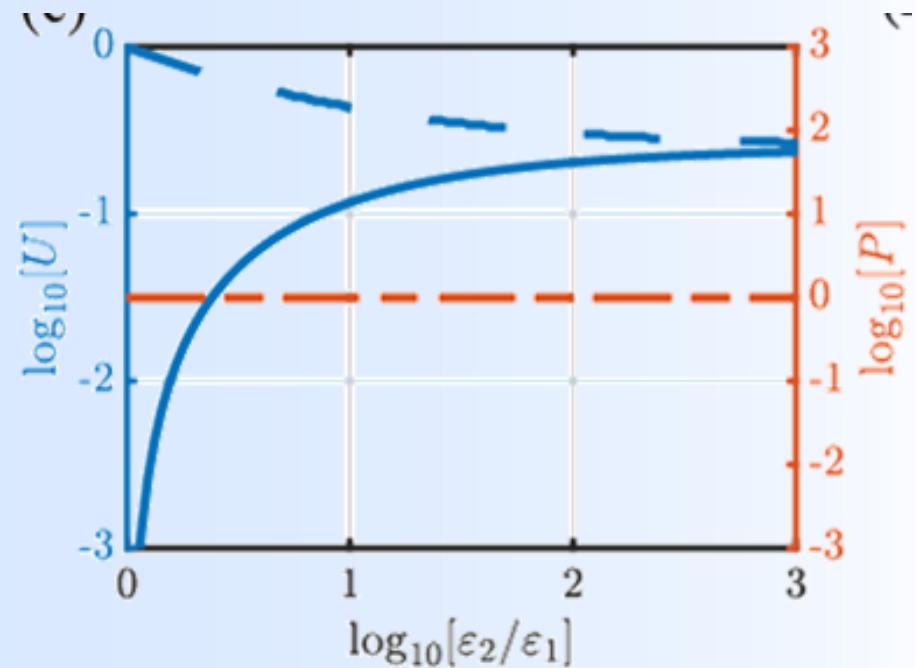
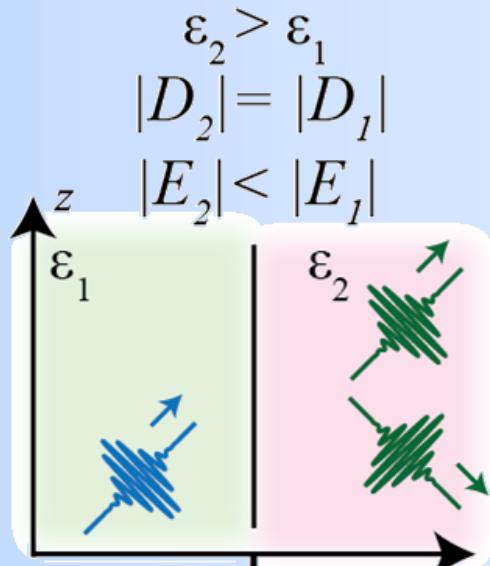


ENERGY REQUIREMENTS



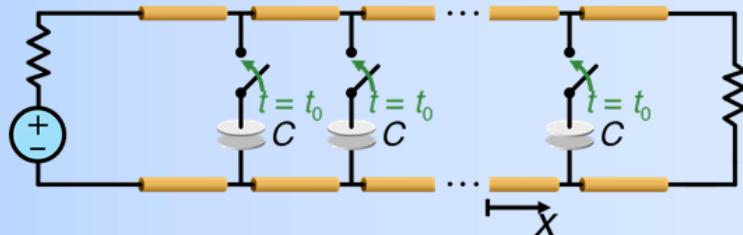
ENERGY REQUIREMENTS

Pockels/Kerr mod. (off)

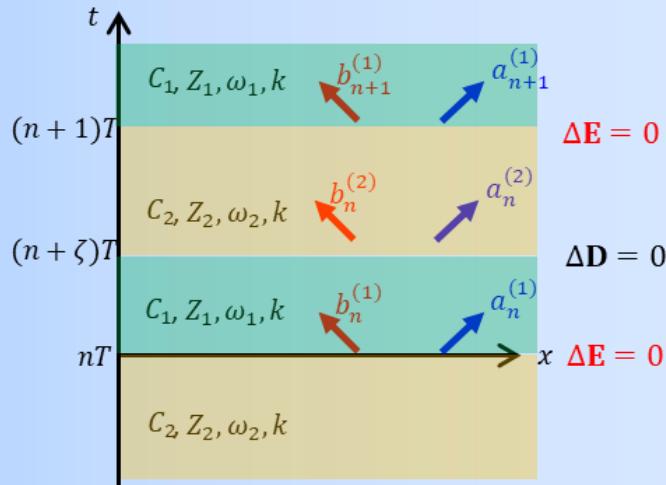


PASSIVE PHOTONIC TIME CRYSTAL

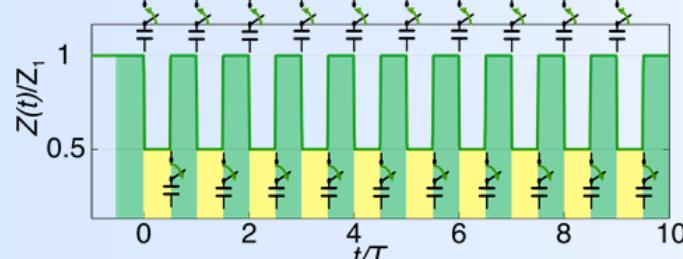
Transmission-line metamaterial



Temporal transfer matrix method



Switching scheme



$$V_n^{(1)}(x, t) = V_n^{(1)}(t) e^{-j k x} = \left[a_n^{(1)} e^{j \omega_1 (t - nT)} + b_n^{(1)} e^{-j \omega_1 (t - nT)} \right] e^{-j k x}$$

$$I_n^{(1)}(x, t) = I_n^{(1)}(t) e^{-j k x} = \frac{1}{Z_1} \left[a_n^{(1)} e^{j \omega_1 (t - nT)} - b_n^{(1)} e^{-j \omega_1 (t - nT)} \right] e^{-j k x}$$

$$\begin{bmatrix} V \\ I \end{bmatrix}_{n+1} = M \begin{bmatrix} V \\ I \end{bmatrix}_n = e^{j \Omega T} \begin{bmatrix} V \\ I \end{bmatrix}_n$$

$$\text{tr}(M) = 2 \cos \Omega T - (1 - z^2) e^{-j \Omega T}$$

where $z = \frac{Z_2}{Z_1} = \sqrt{\frac{C_1}{C_2}} < 1$ is the ratio of impedances



PASSIVE PHOTONIC TIME CRYSTAL

Comparison to *active* photonic time crystals (PTC)

	Passive PTC	Active PTC
Dispersion	$\text{tr}(M) = 2 \cos \Omega T - (1 - z^2)e^{-j\Omega T}$	$\text{tr}(M) = 2 \cos \Omega T$
Momentum balance	$\det(M) = \frac{c_1}{c_2} = z^2 < 1$ (Possible momentum decay)	$\det(M) = 1$

From the transfer matrix M , it is easy to find

$\text{tr}(M) \in \mathbb{R}$, for both passive and active PTCs

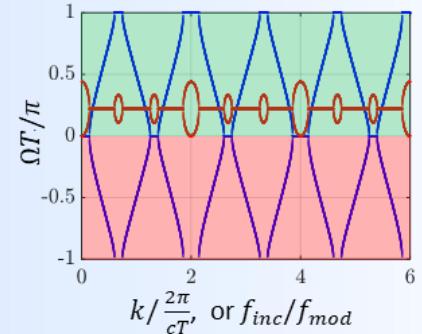
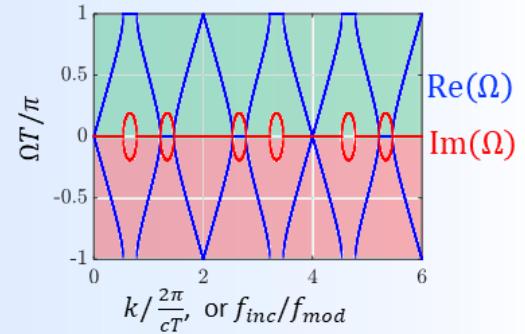


For an active PTC,

Ω is either **real** (in the bands) or **purely imaginary** (in the gaps).

For a passive PTC,

Ω is **complex almost everywhere** except for countably infinite momenta.



Passive PTC with stable k -gaps

Condition for the existence of stable k -gaps:

$$1 + \det(M) - |\text{tr}(M)| = 0,$$

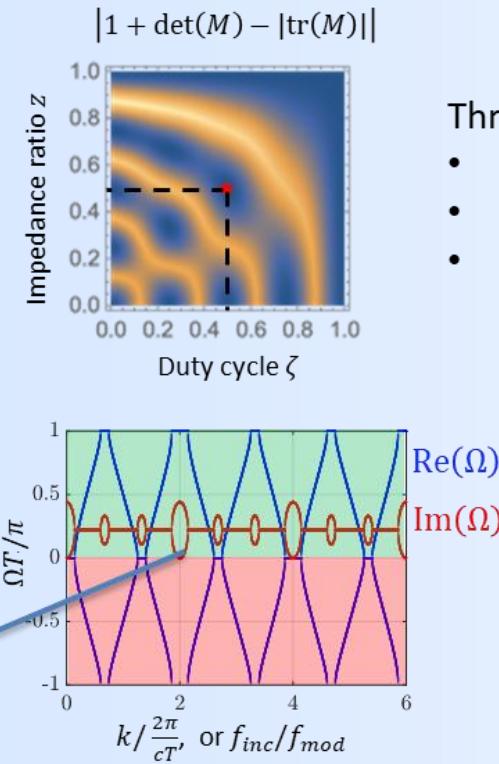
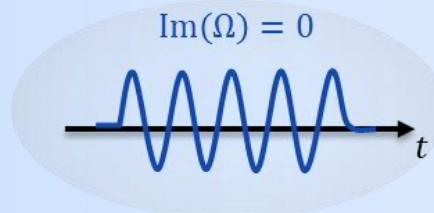
where the transfer matrix $M = \begin{bmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{bmatrix}$ with

$$M_{11} = z^2 \cos \theta_1 \cos \theta_2 + z \sin \theta_1 \sin \theta_2,$$

$$M_{12} = jz^2 \cos \theta_1 \cos \theta_2 - jz \sin \theta_1 \sin \theta_2,$$

$$M_{21} = \frac{j(1+z)}{2} \sin(\theta_1 + \theta_2) + \frac{j(1-z)}{2} \sin(\theta_1 - \theta_2),$$

$$M_{22} = \frac{1+z}{2} \cos(\theta_1 + \theta_2) + \frac{1-z}{2} \cos(\theta_1 - \theta_2).$$

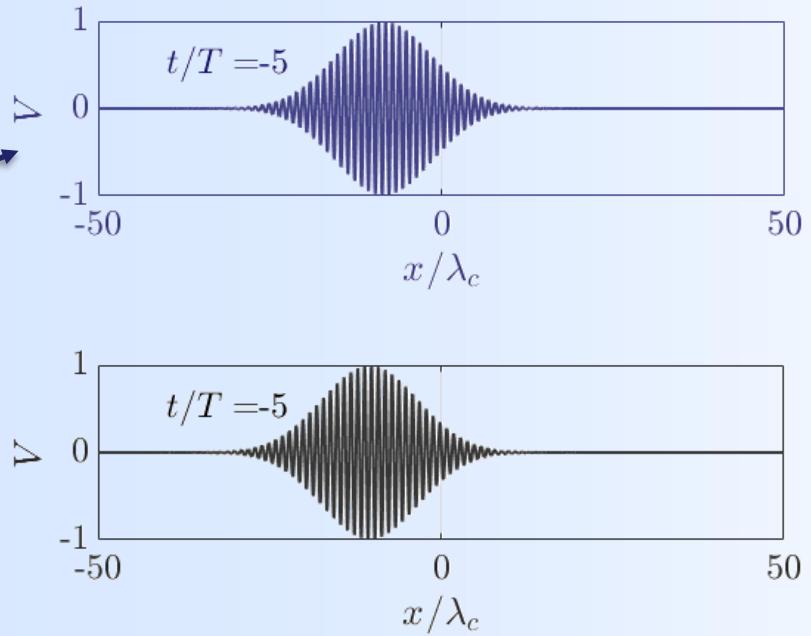
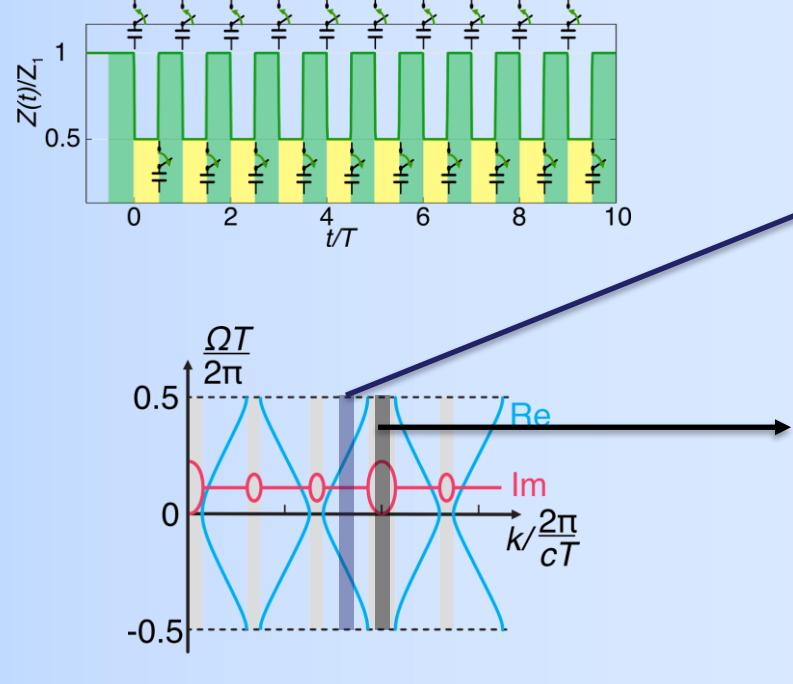


Three degrees of freedom:

- Modulation frequency;
- Duty cycle;
- Impedance contrast.

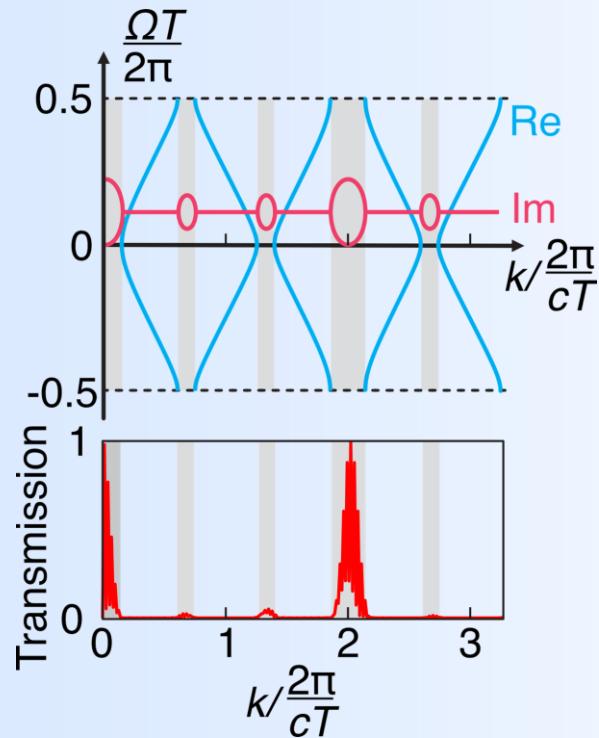
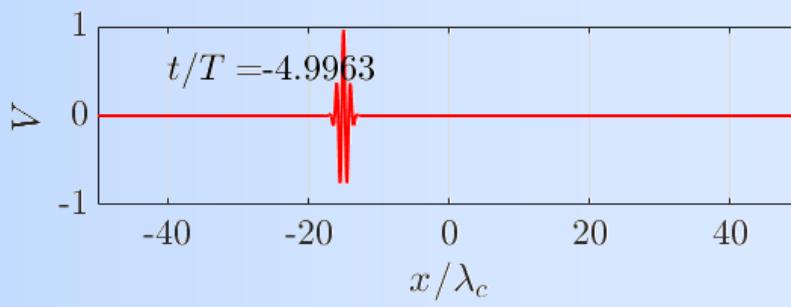
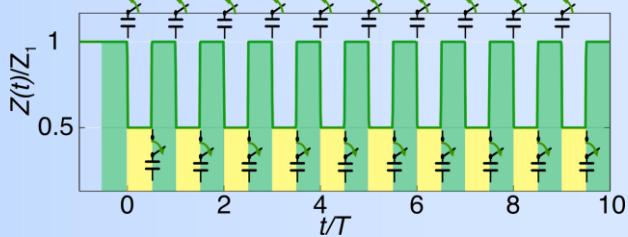


MOMENTUM SELECTIVITY

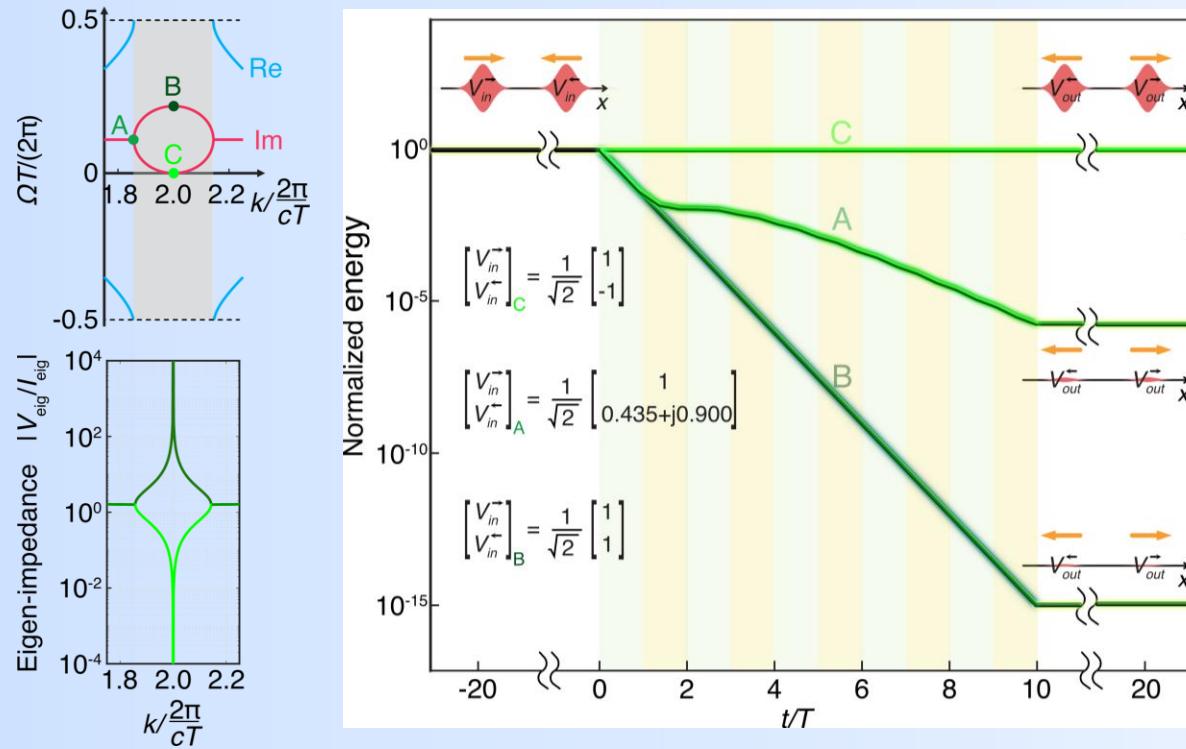


MOMENTUM SELECTIVITY

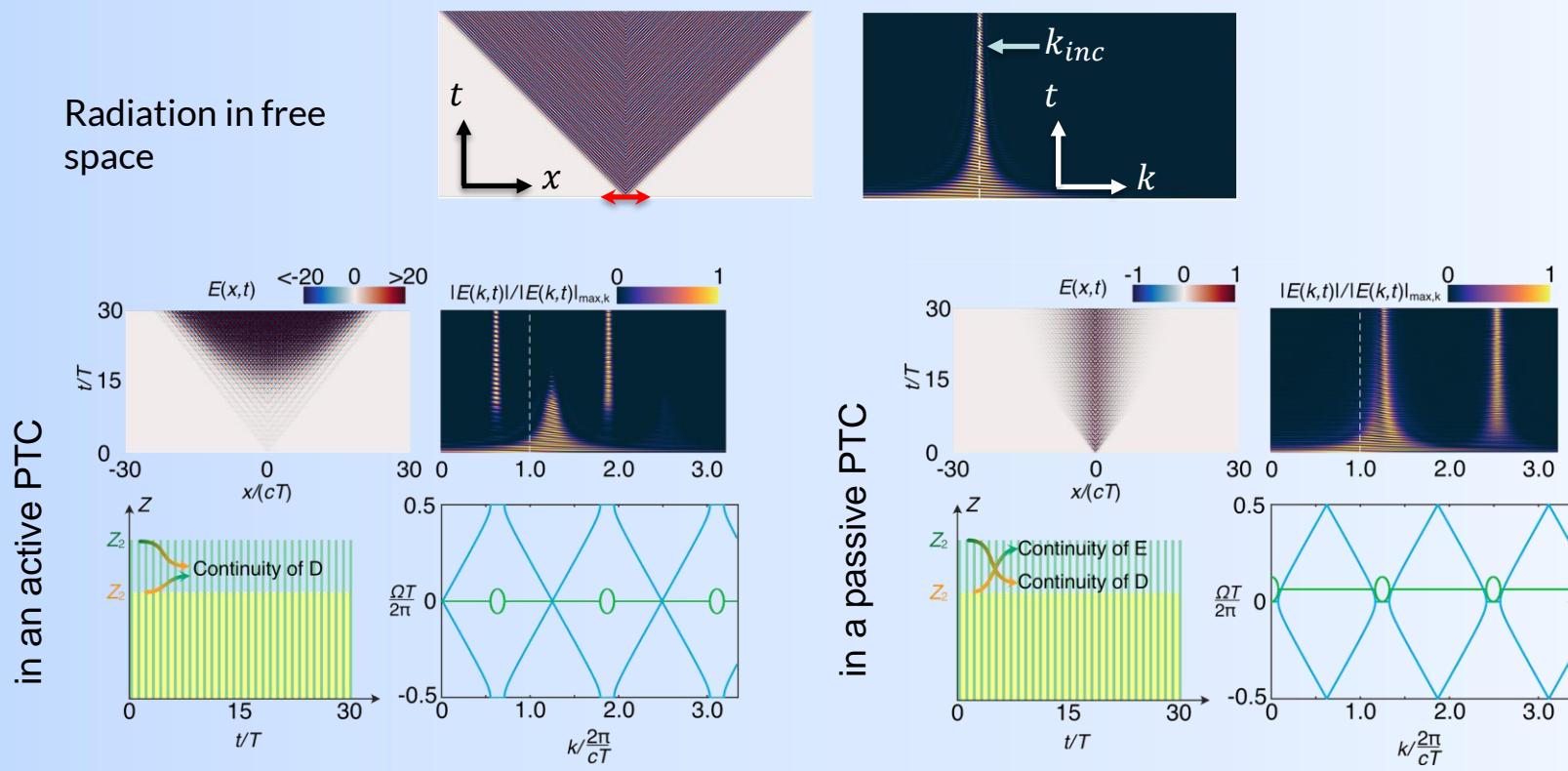
Broadband (short pulse) incidence



COHERENT PASSIVE K-GAPS



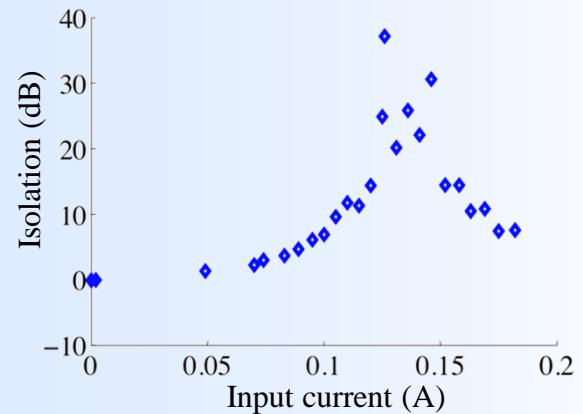
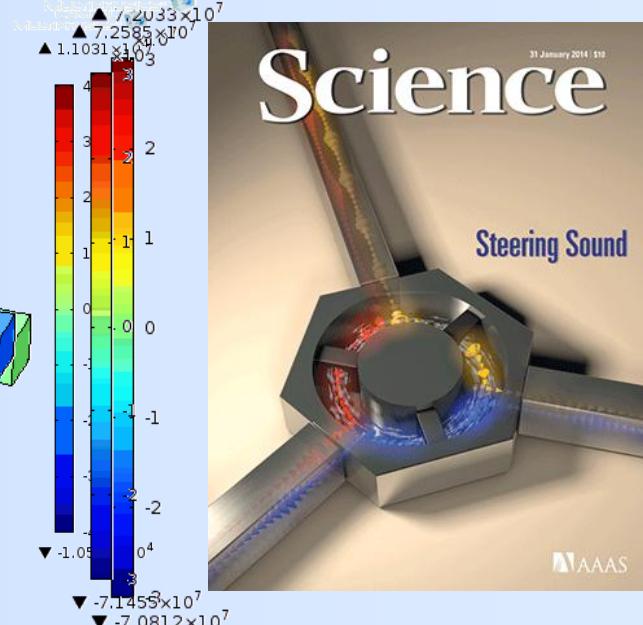
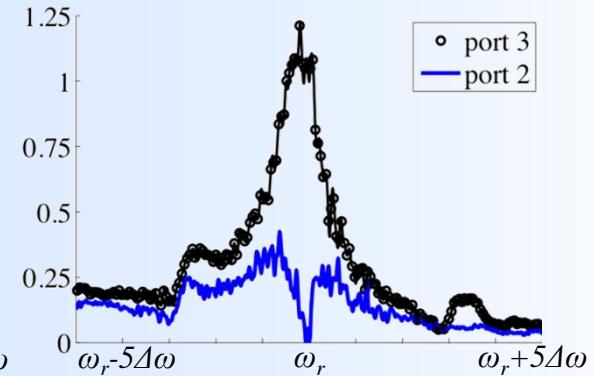
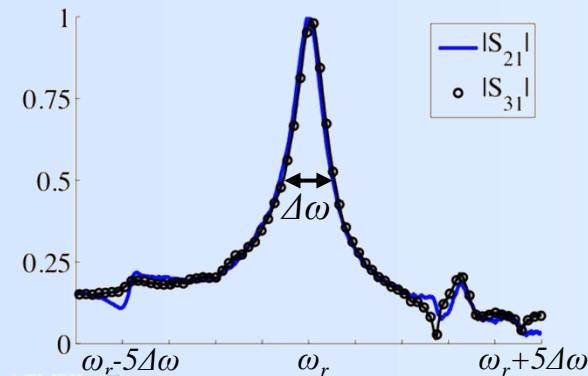
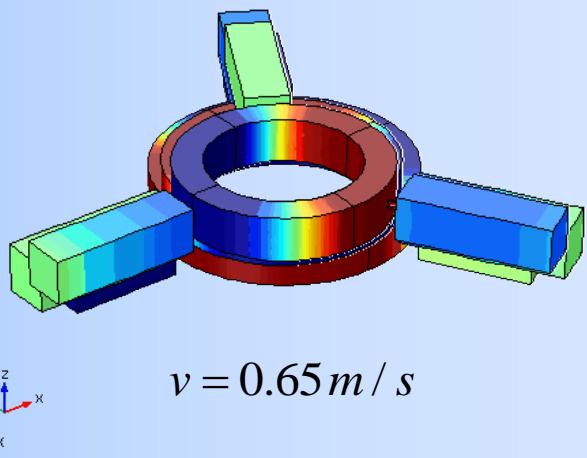
EMITTER COUPLED TO A PASSIVE TIME CRYSTAL



ANGULAR-MOMENTUM BIAS



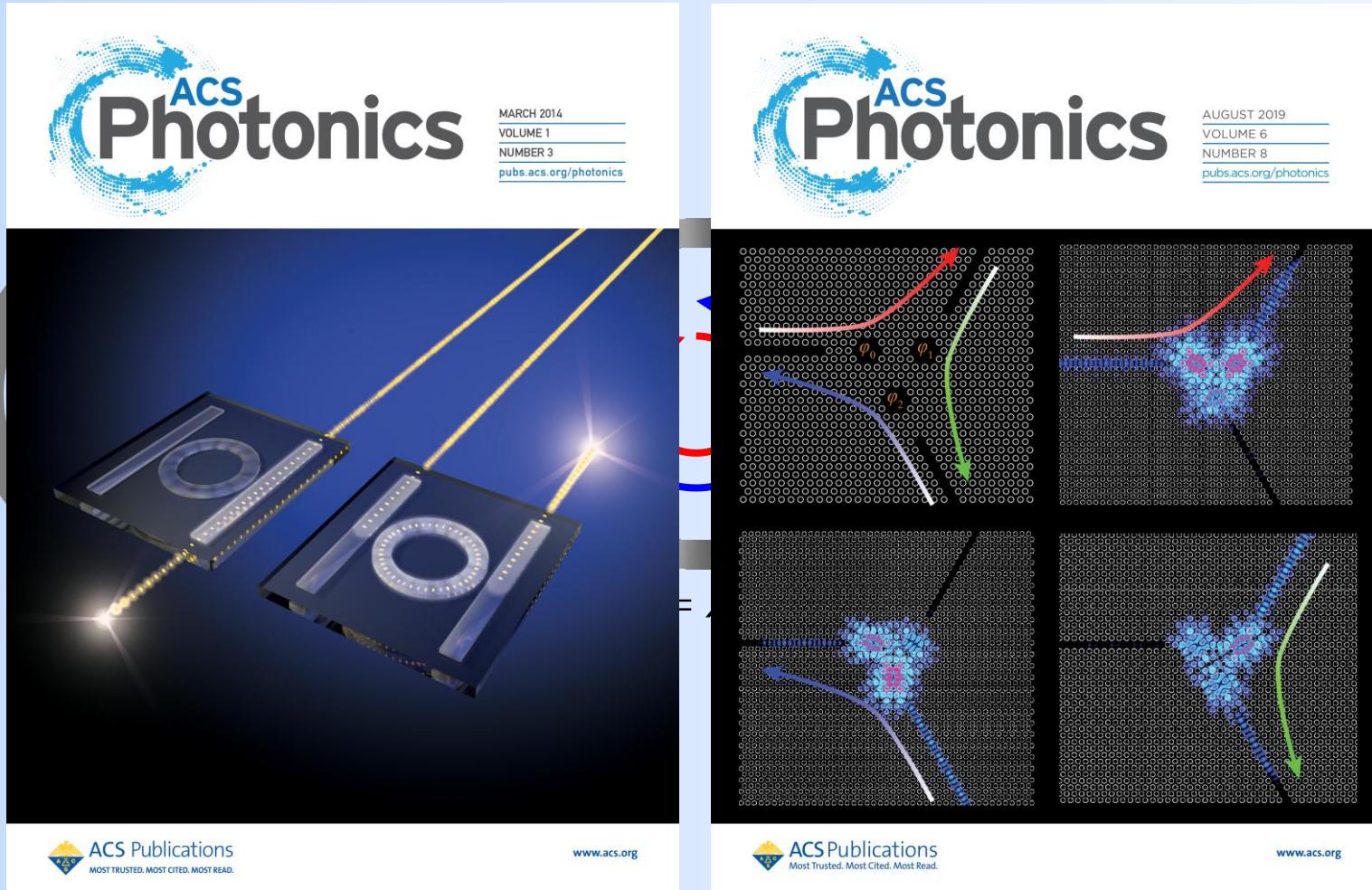
freq(153)=2955.5 Surface: Pressure (Pa)
freq(153)=2955.5 Surface: Pressure (Pa)
freq(58)=944 Surface: Pressure (Pa)



R. Fleury, D. L. Sounas, C. Sieck, M. Haberman, A. Alù, *Science* 343, 516 (2014)



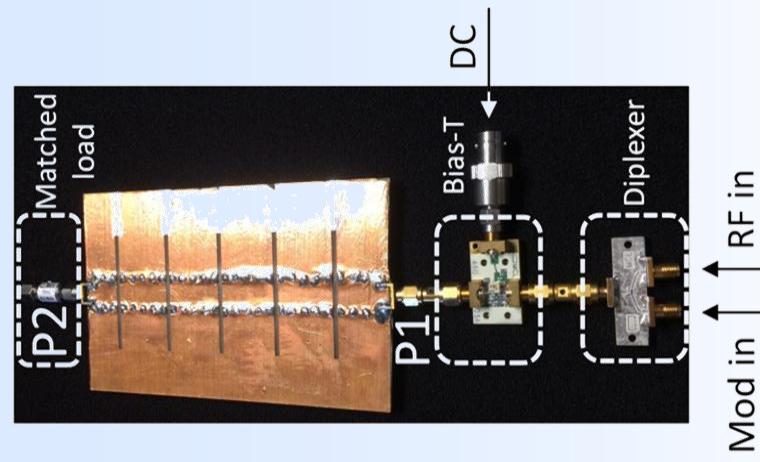
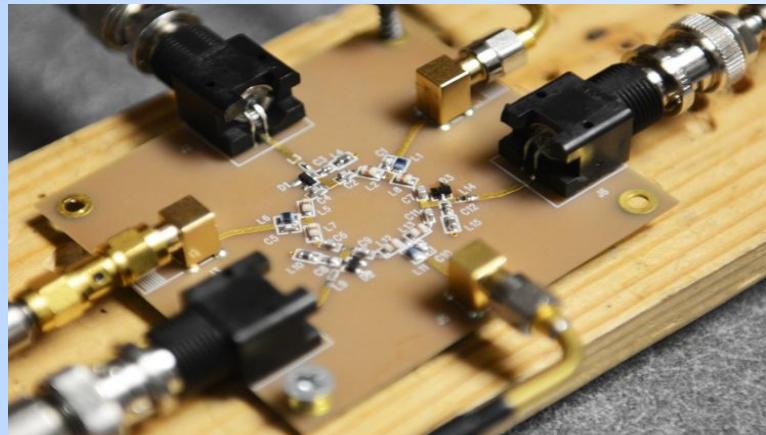
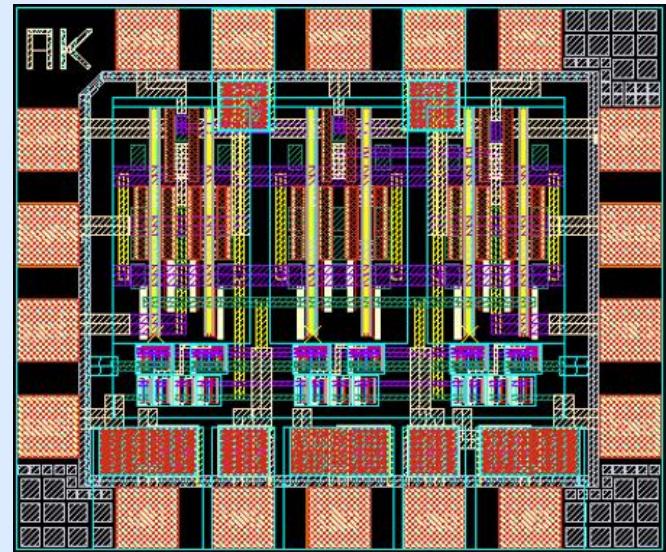
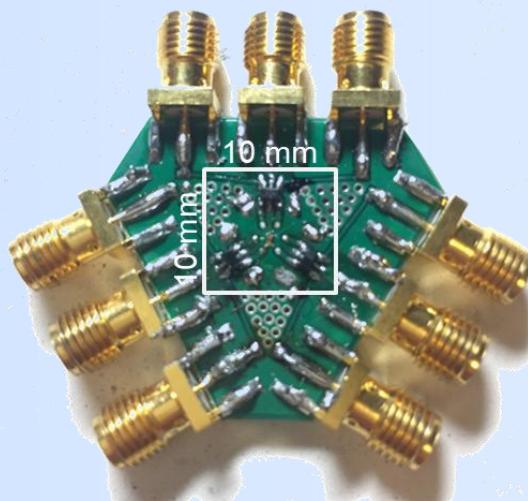
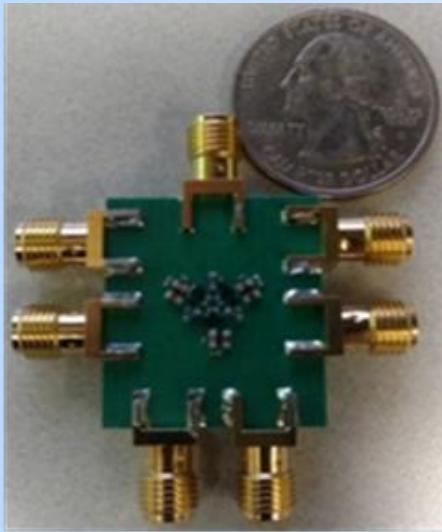
SYNTHETIC ANGULAR MOMENTUM WITH TIME MODULATION



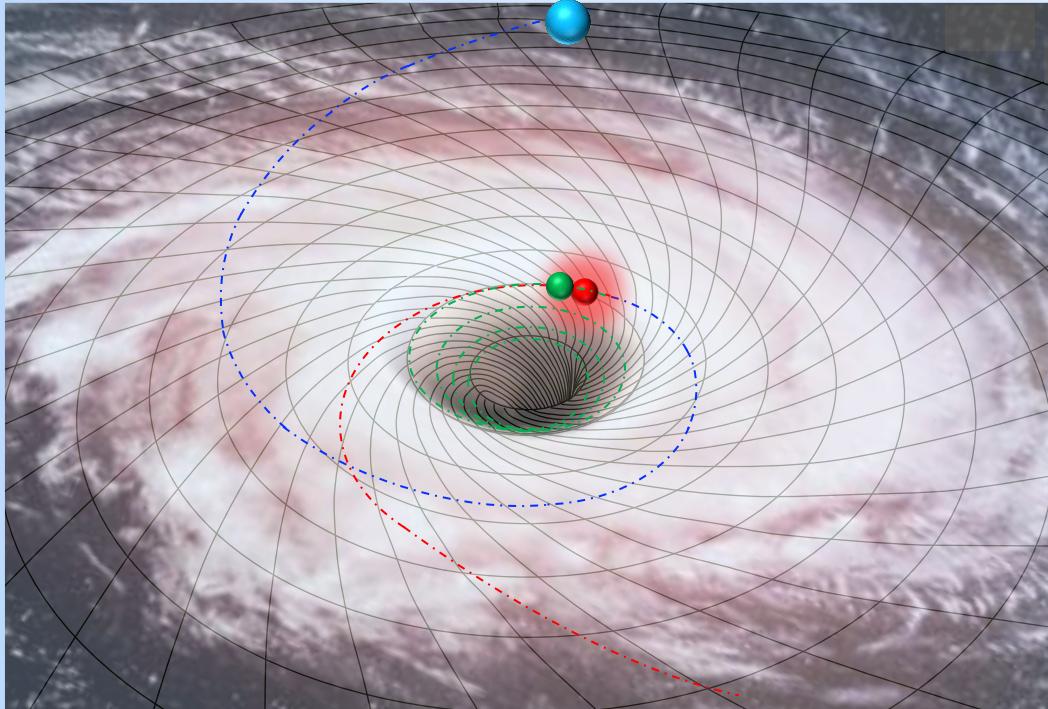
D. Sounas, A. Alù, *ACS Photonics* 1, 198 (2014)
A. Mock, D. Sounas, A. Alù, *ACS Photonics* 6, 2056 (2019)



COMPACT, MAGNET-LESS CIRCULATORS FOR RF AND LIGHT



EXTRACTING ENERGY FROM A ROTATING BLACK HOLE



Extraction of Rotational Energy from a Black Hole

THERE has been considerable interest recently in the question of the gravitational collapse of a massive body and of the possible astrophysical consequences of the existence of the "black hole" which general relativity predicts should sometimes be the result of such a collapse. In particular, the question has arisen whether the mass-energy content of a black hole could, under suitable circumstances, be a source of available energy. We now consider the extraction of rotational energy from a black hole, not least because the rotational energy (defined appropriately) of a black hole should, in general, be comparable with its total mass-energy¹.

Department of Mathematics,
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Received December 16, 1970.

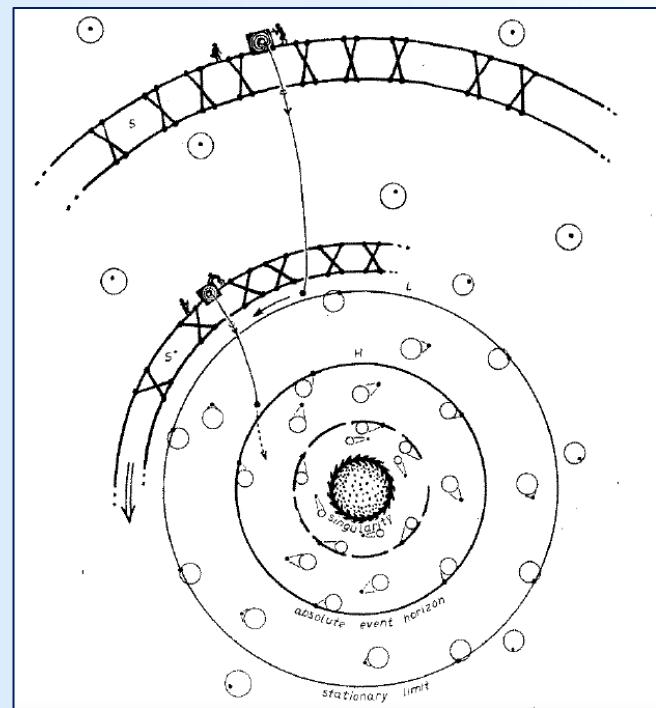
¹ Bardeen, J. M., *Nature*, **226**, 65 (1970).
² Israel, W., *General Relativity and Gravitation*, **2**, No. 1 (Plenum, in the press).

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R. PENROSE

R. M. FLOYD

Negative energy orbits



R. Penrose,
Riv. Nuovo Cim. Num. Spec. 1, 257 (1969)

ELECTROMAGNETIC PENROSE SUPER-RADIANCE

GENERATION OF WAVES BY A ROTATING BODY

Ya.B. Zel'dovich

Institute of Applied Mathematics, USSR Academy of Sciences

Submitted 9 July 1971

ZhETF Pis. Red. 14, No. 4, 270 - 272 (20 August 1971)

An axially-symmetrical body rotating inside a resonator cavity is capable of amplifying definite oscillation modes inside the resonator, transferring the rotation energy to these oscillations.

SOVIET PHYSICS JETP

VOLUME 35, NUMBER 6

DECEMBER, 1972

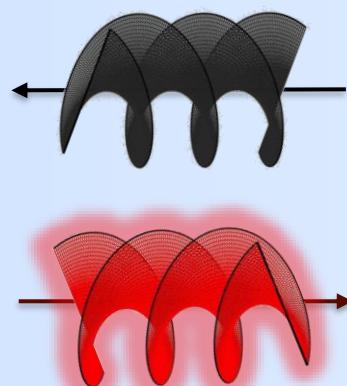
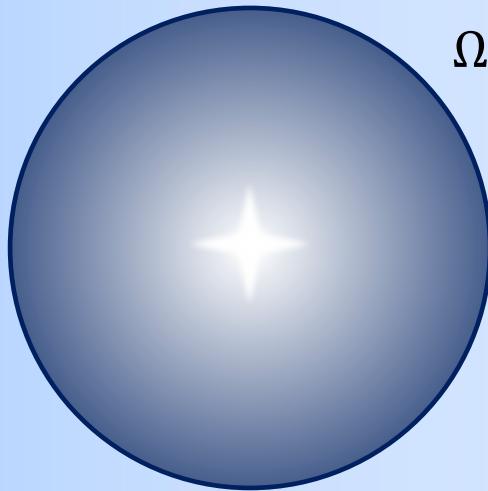
Amplification of Cylindrical Electromagnetic Waves Reflected from a Rotating Body

Ya. B. Zel'dovich

Institute of Applied Mathematics, USSR Academy of Sciences

Submitted December 10, 1971

Zh. Eksp. Teor. Fiz. 62, 2076-2081 (June, 1972)



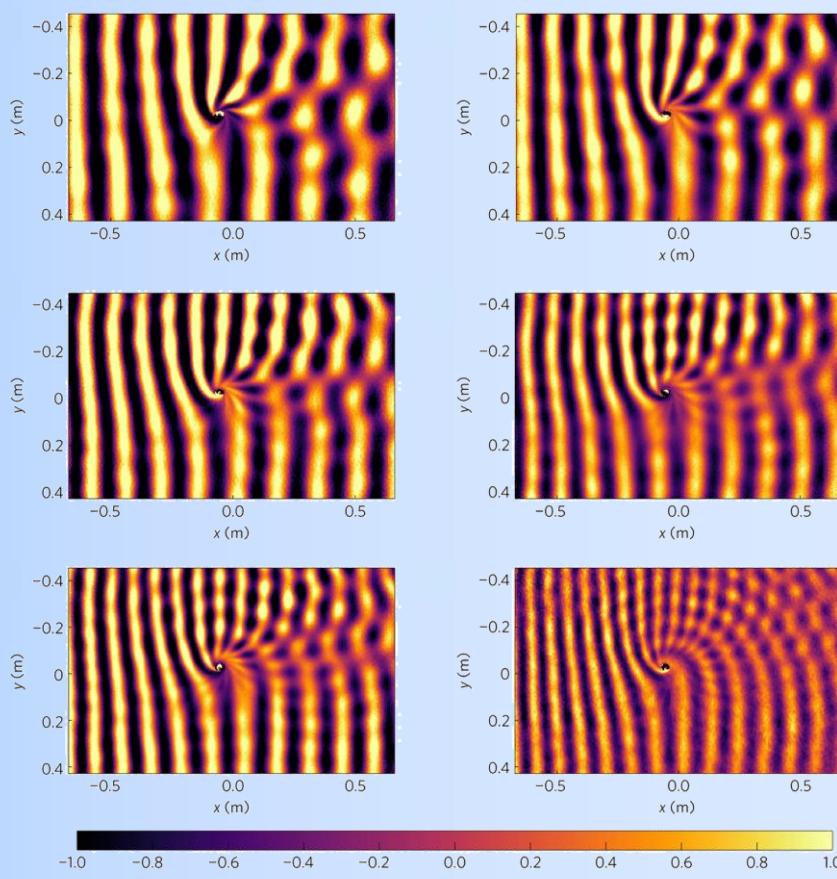
$$E(r, \varphi) = e^{-j\omega t} e^{jn\varphi} e^{-jkr} / \sqrt{r}$$

$$\omega(k) - v \cdot k < 0$$

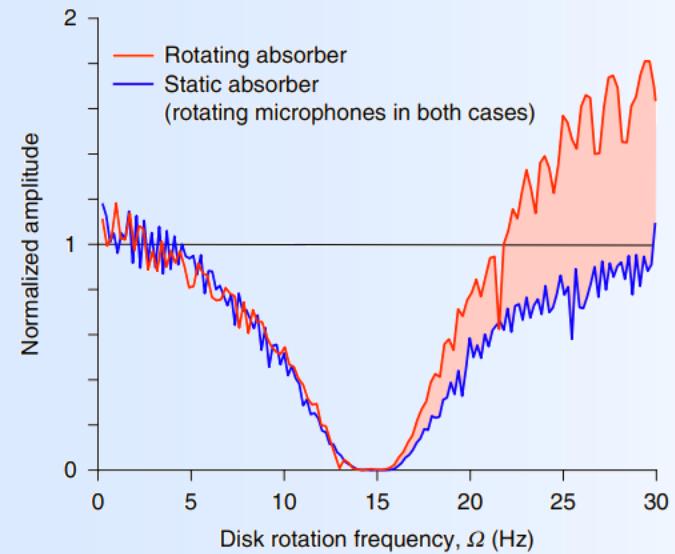
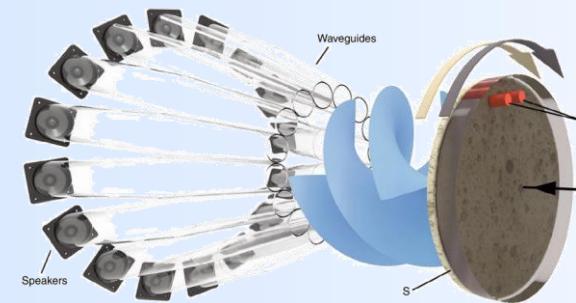


WAVE AMPLIFICATION FROM ROTATING BODIES

Rotational superradiant scattering in a *water vortex flow*



Amplification of **acoustic** waves from a rotating body

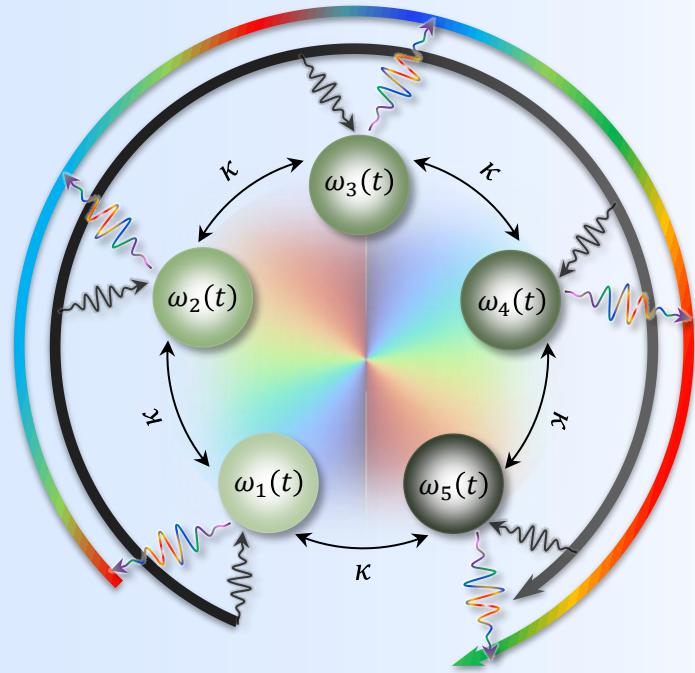
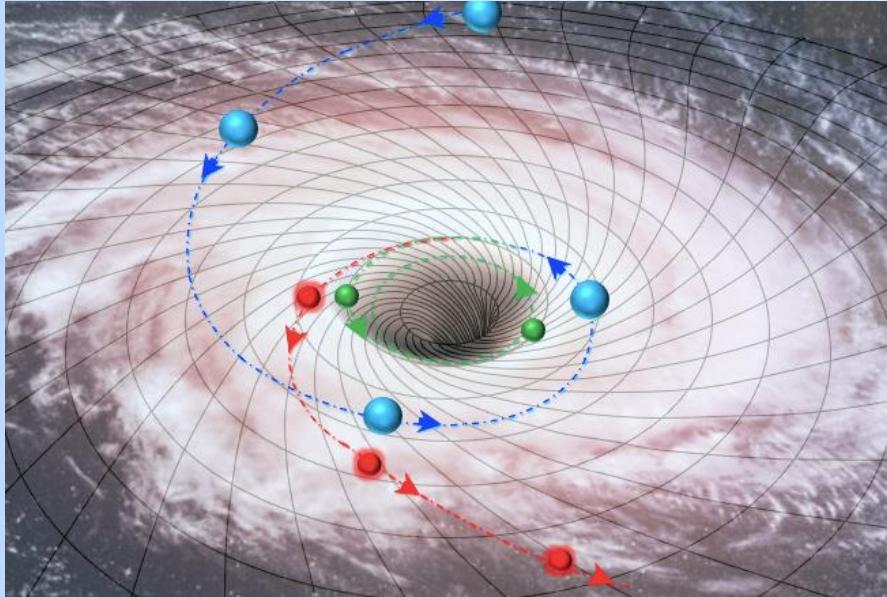


T. Torres et. al., *Nature Phys.* 13, 833–836 (2017)

M. Cromb, et al., *Nature Physics* 16, 1069 (2020)

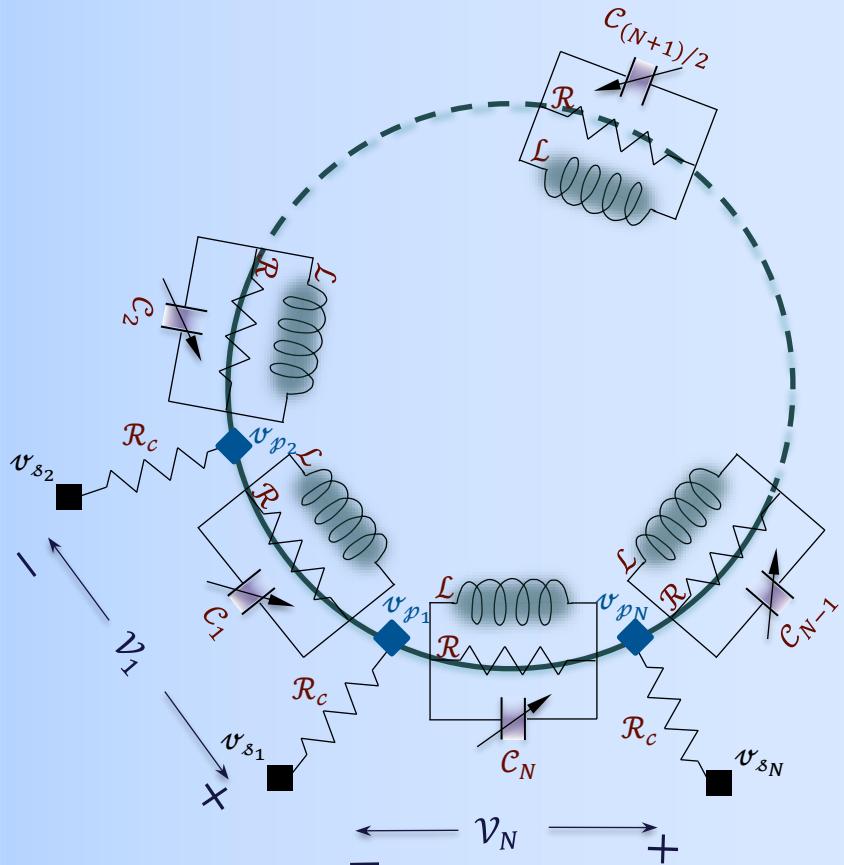


SYNTHETIC PENROSE SUPER-RADIANCE



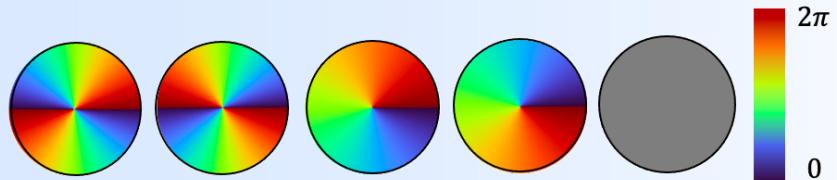
- Excitation with OAM signal
- Synthetic rotation based on Angular Momentum Biasing
- Monitoring the strength of output OAM signals

GOVERNING MATRIX EQUATION & HARMONIC ANALYSIS



$$\{\bar{\bar{\mathbf{A}}}\ \omega^2 + \bar{\bar{\mathbf{B}}}\ \omega + \bar{\bar{\mathbf{D}}}\} \bar{\mathbf{V}} + \frac{\delta}{2} \left\{ \omega^2 \bar{\bar{\mathbf{C}}} \bar{\mathbf{V}}^{(+)} + \omega^2 \bar{\bar{\mathbf{C}}}^* \bar{\mathbf{V}}^{(-)} \right\} = \bar{\bar{\mathbf{G}}} \bar{\mathbf{V}}_s$$

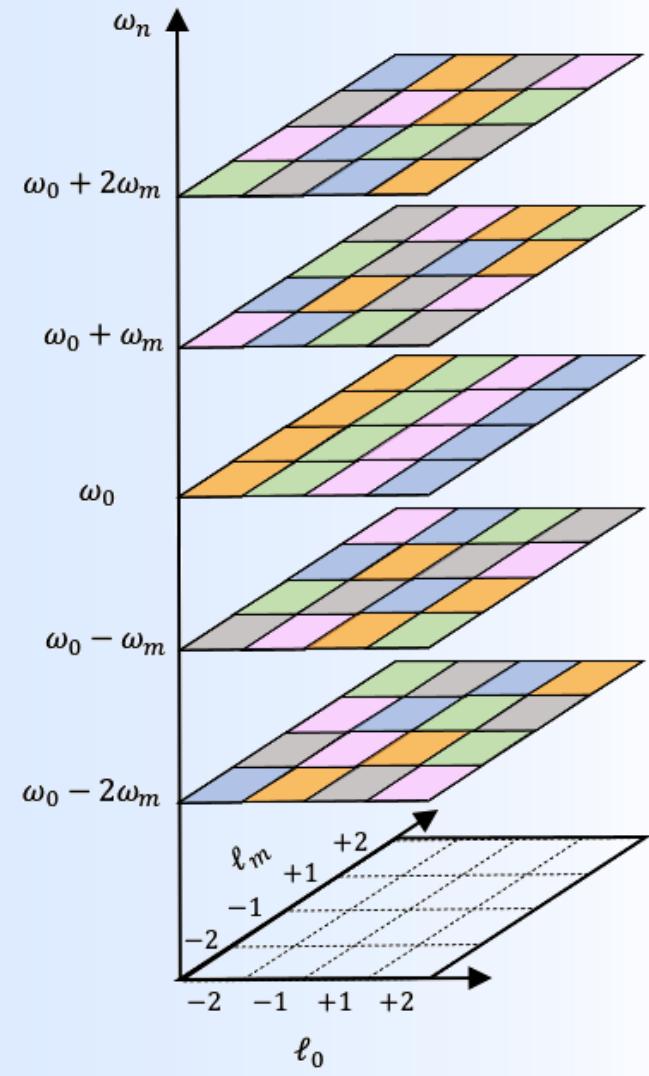
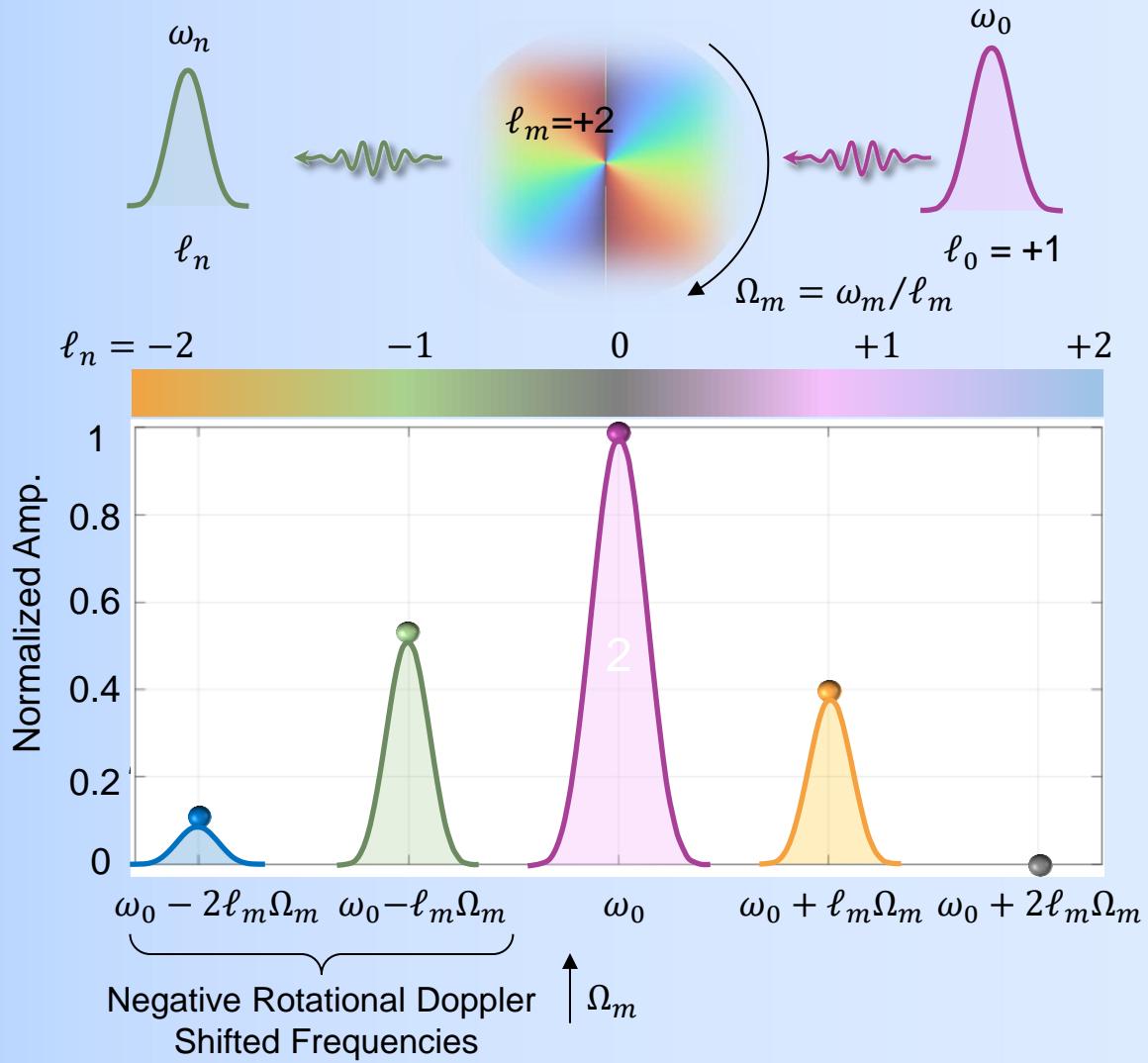
Eigenbasis of the static system



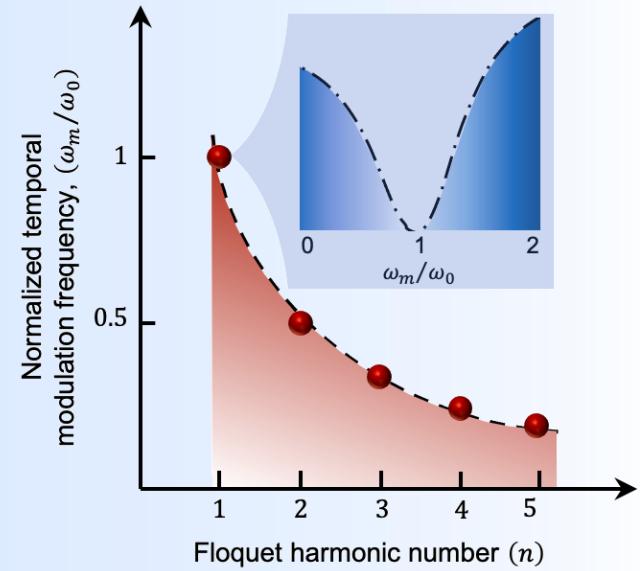
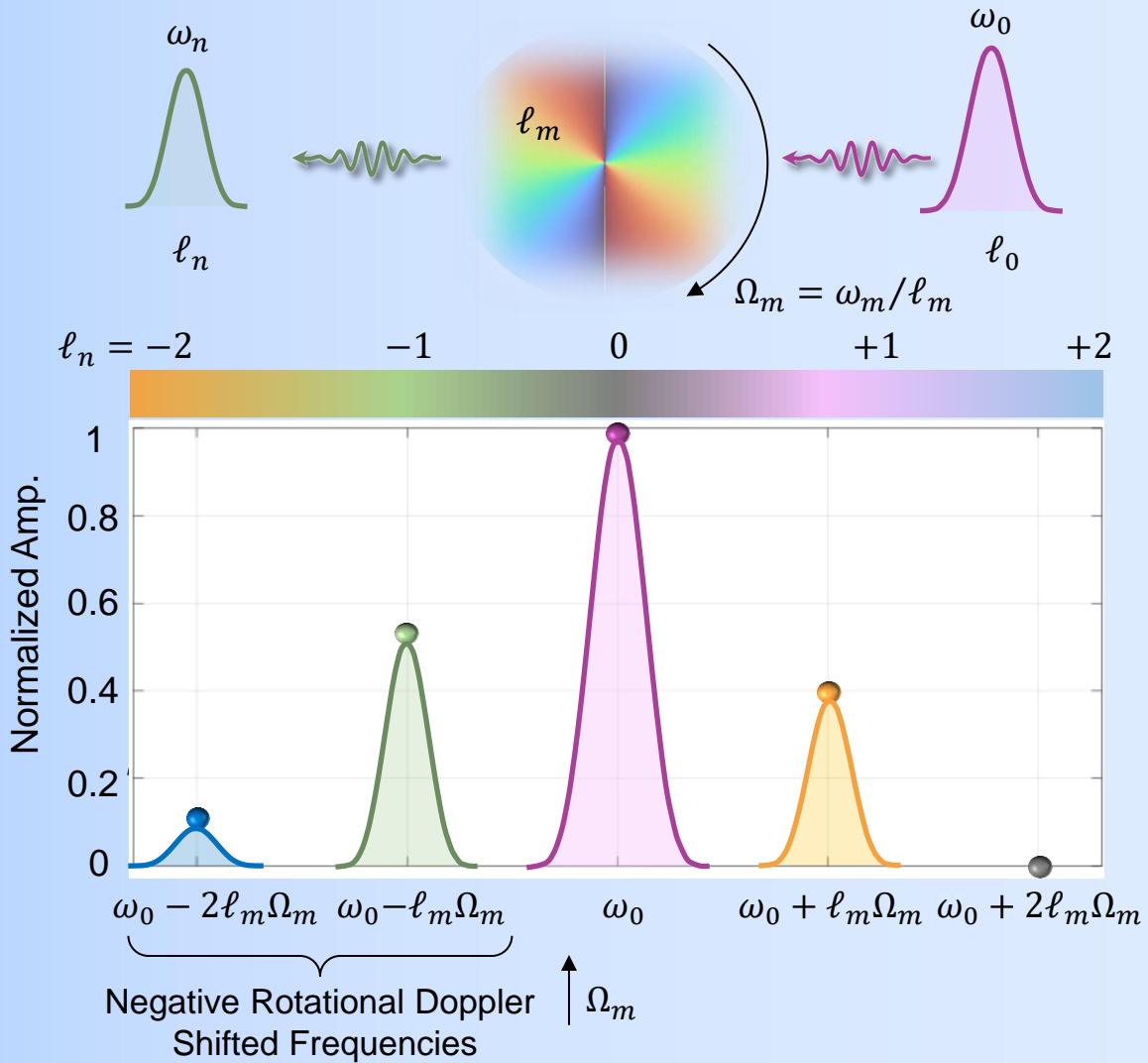
$$\bar{\bar{\mathbf{H}}} \bar{\mathbf{V}}_R + \frac{\delta}{2} \left\{ \omega^2 \bar{\bar{\mathbf{C}}}_{R(\ell_m)} \bar{\mathbf{V}}_R^{(+)} + \omega^2 \bar{\bar{\mathbf{C}}}^*_{R(\ell_m)} \bar{\mathbf{V}}_R^{(-)} \right\} = \bar{\bar{\mathbf{G}}}_R \bar{\mathbf{V}}_{s,R}$$

$$\bar{\bar{\mathbf{H}}} = \bar{\bar{\mathbf{A}}}_R \ \omega^2 + \bar{\bar{\mathbf{B}}}_R \ \omega + \bar{\bar{\mathbf{D}}}_R$$

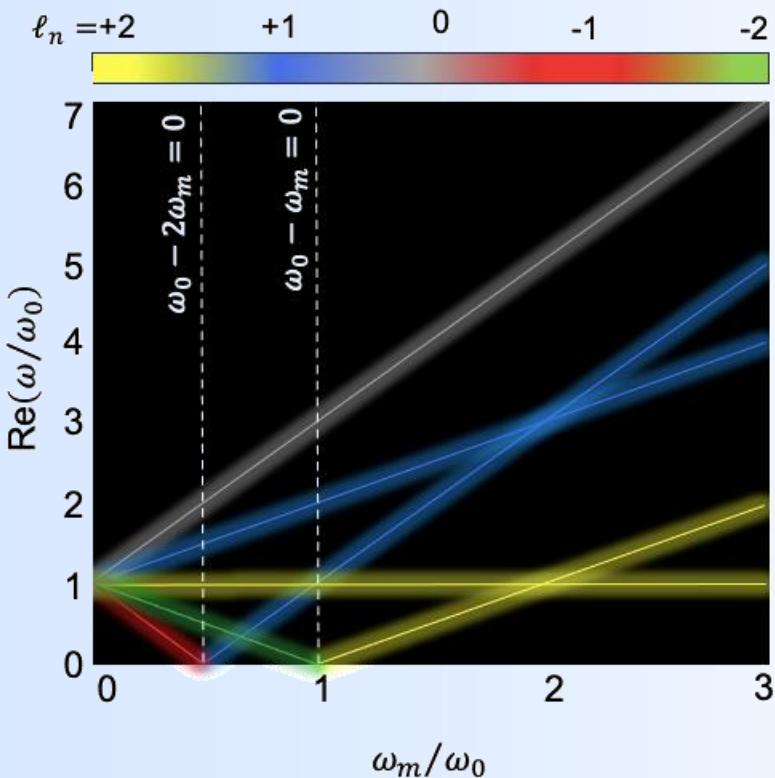
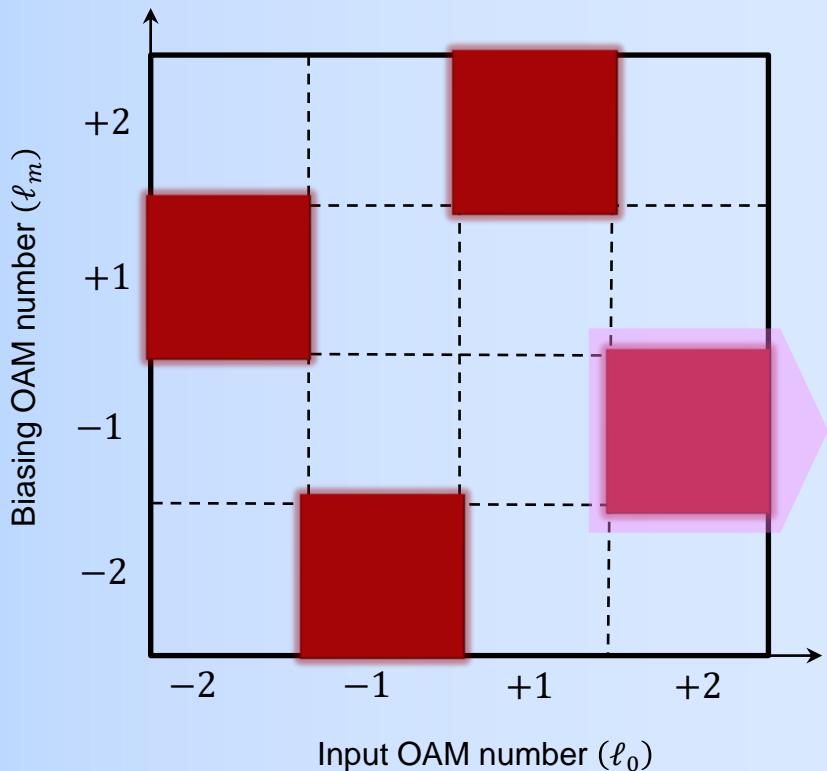
SYNTHETIC ROTATIONAL DOPPLER EFFECT



SYNTHETIC NEGATIVE ROTATION DOPPLER SHIFT



SYNTHETIC PENROSE SUPER-RADIANCE

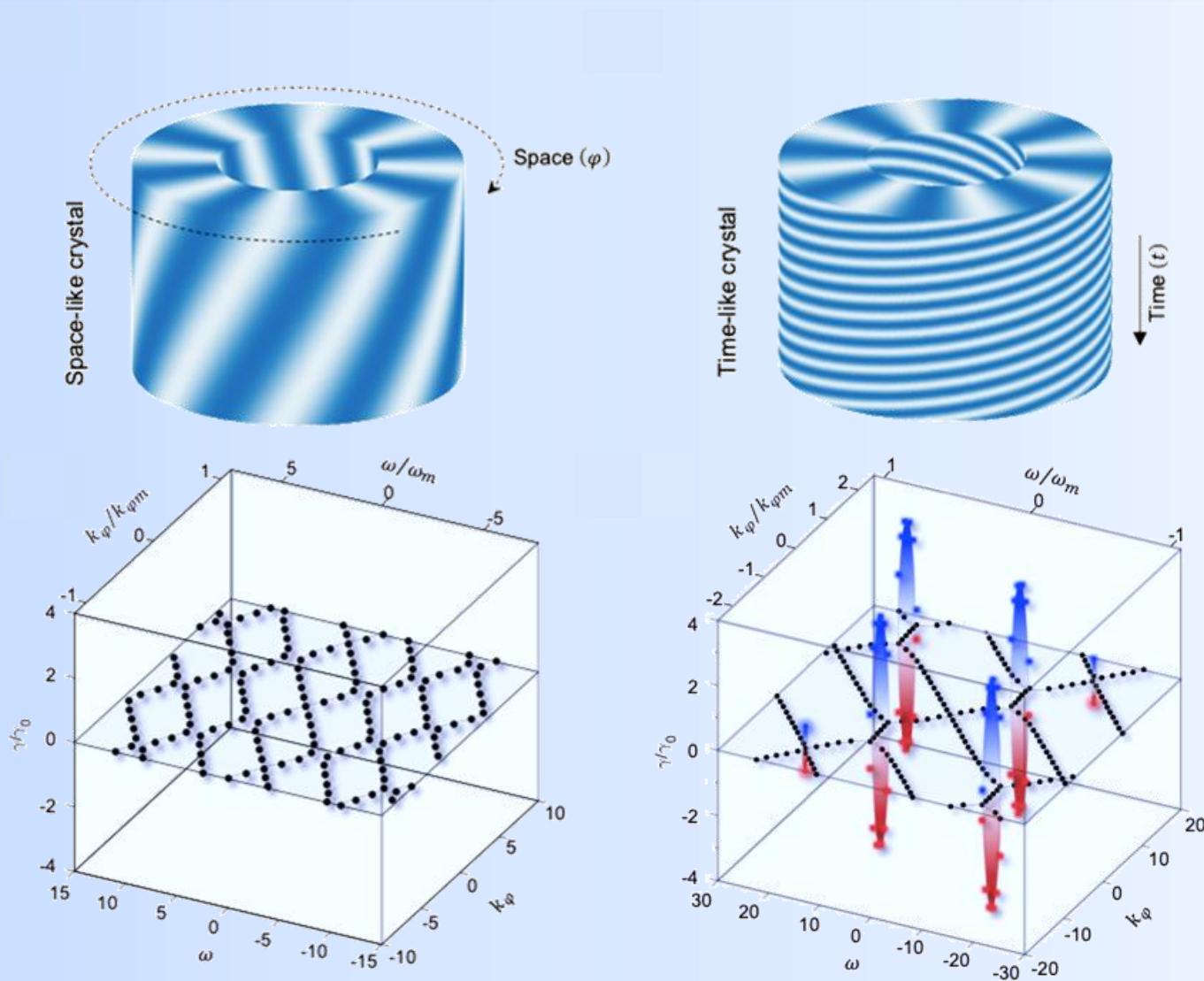


Phase matching between the

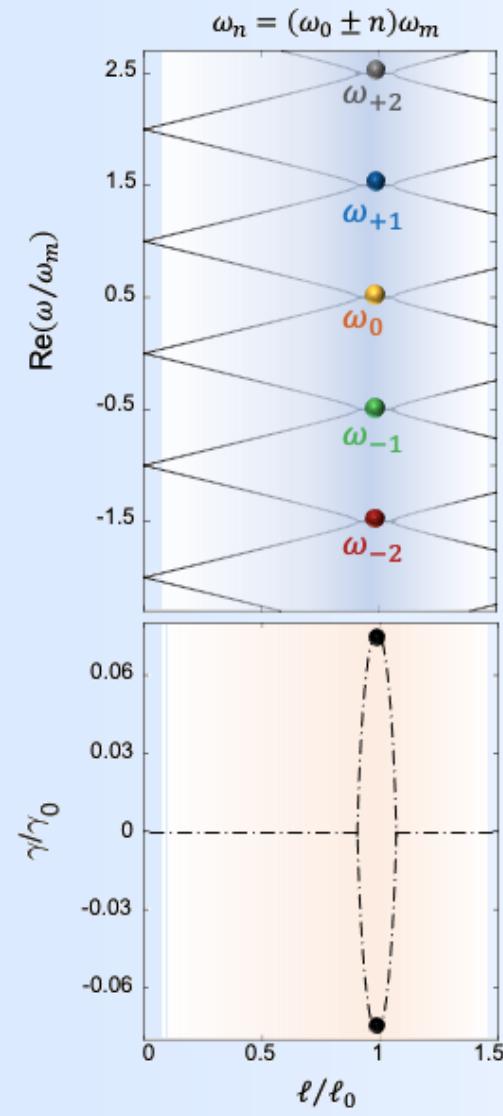
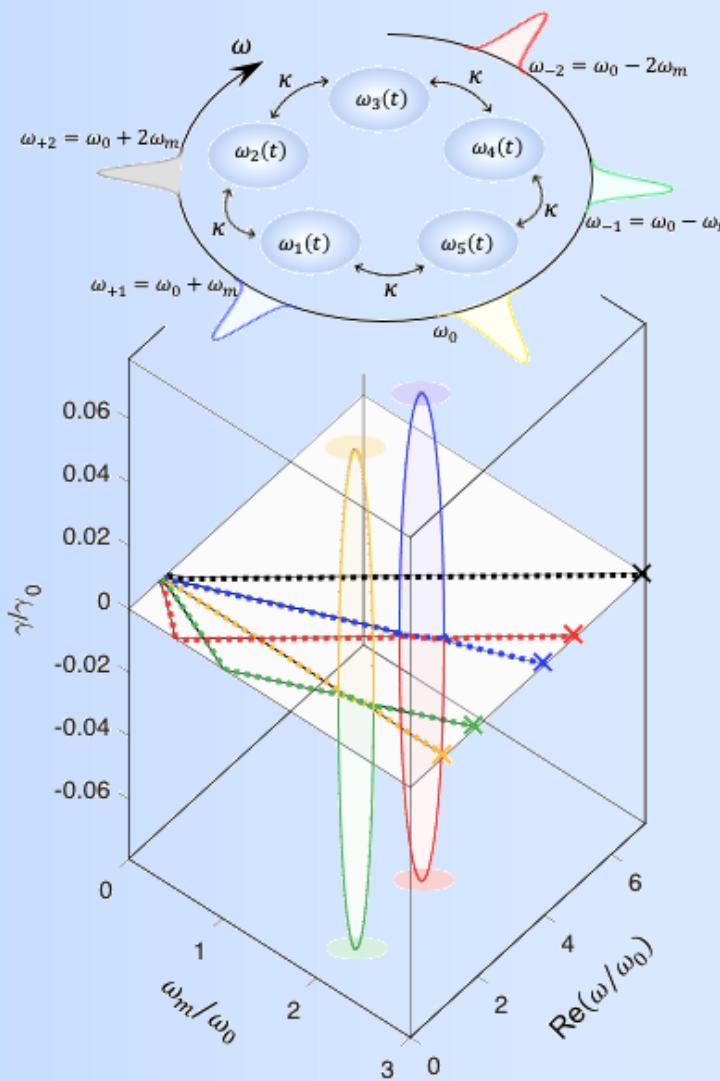
- Modulation and fundamental harmonic
- Interacting harmonics



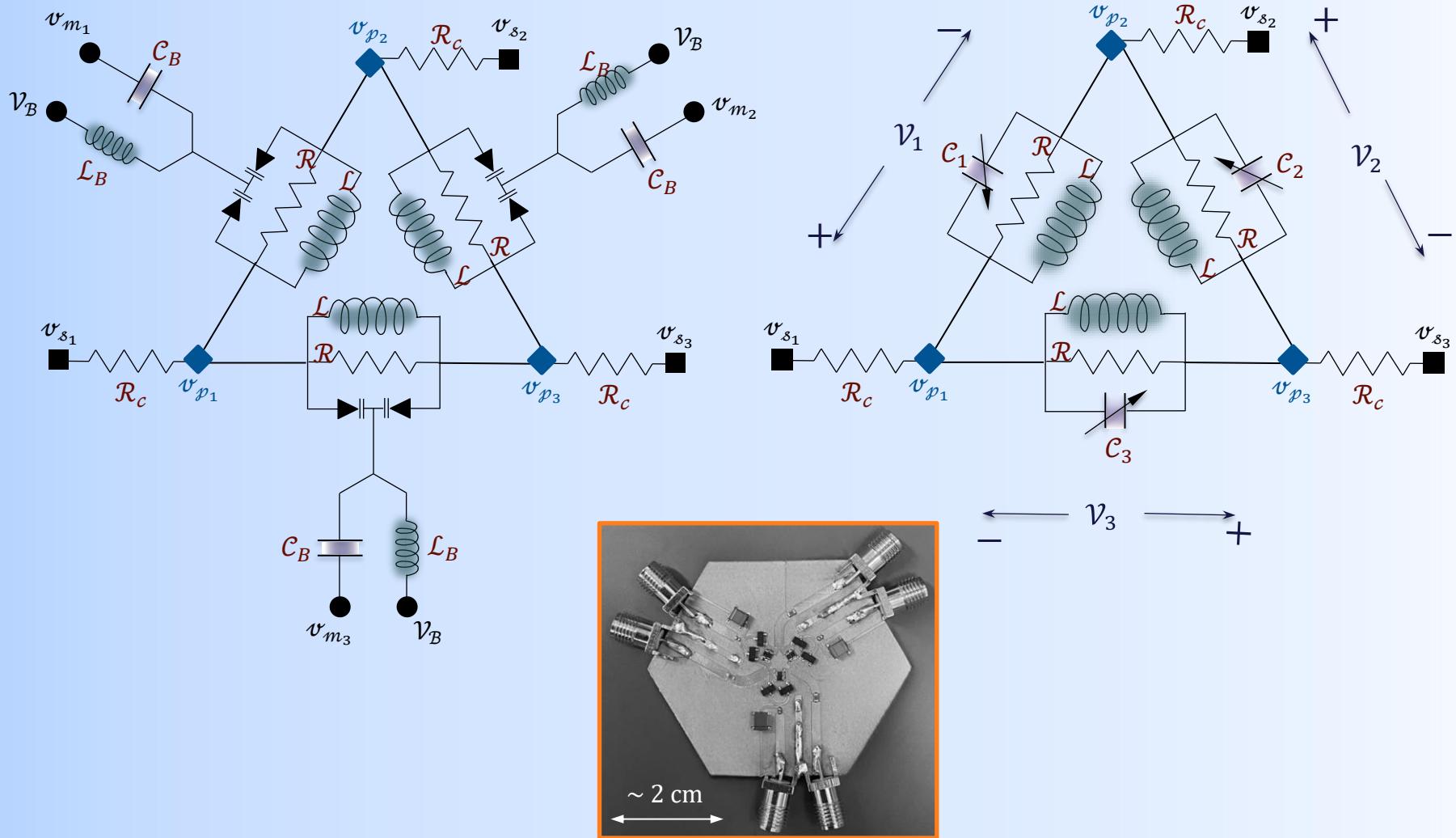
ANGULAR-MOMENTUM BANDGAPS



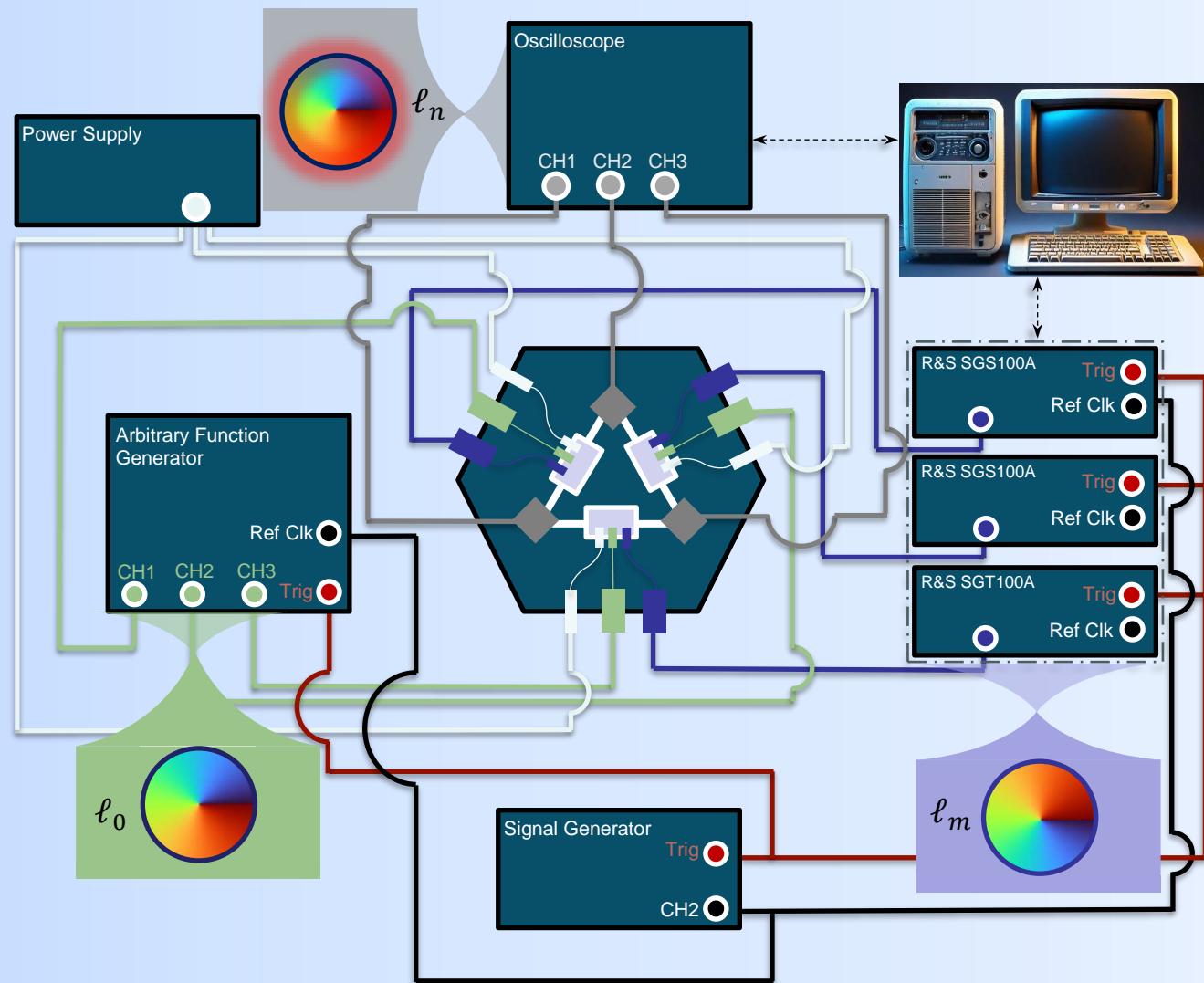
ANGULAR-MOMENTUM BANDGAPS



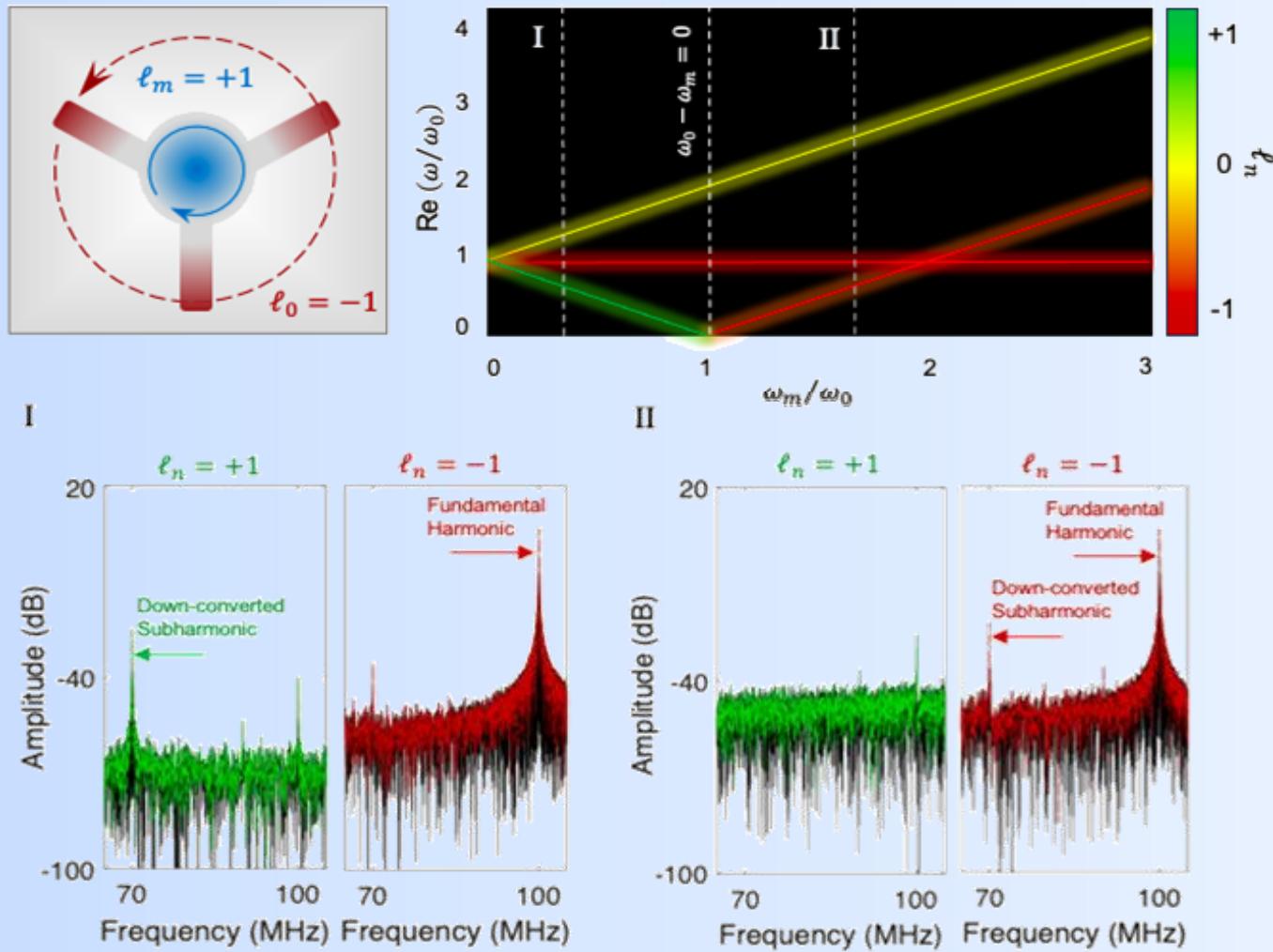
SYNTHETICALLY ROTATING CIRCUIT



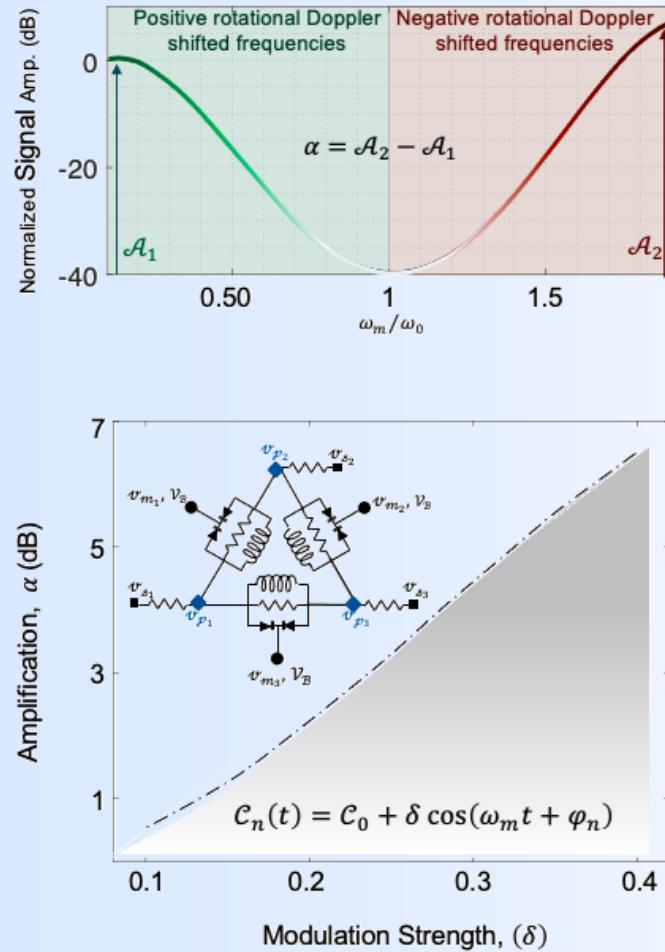
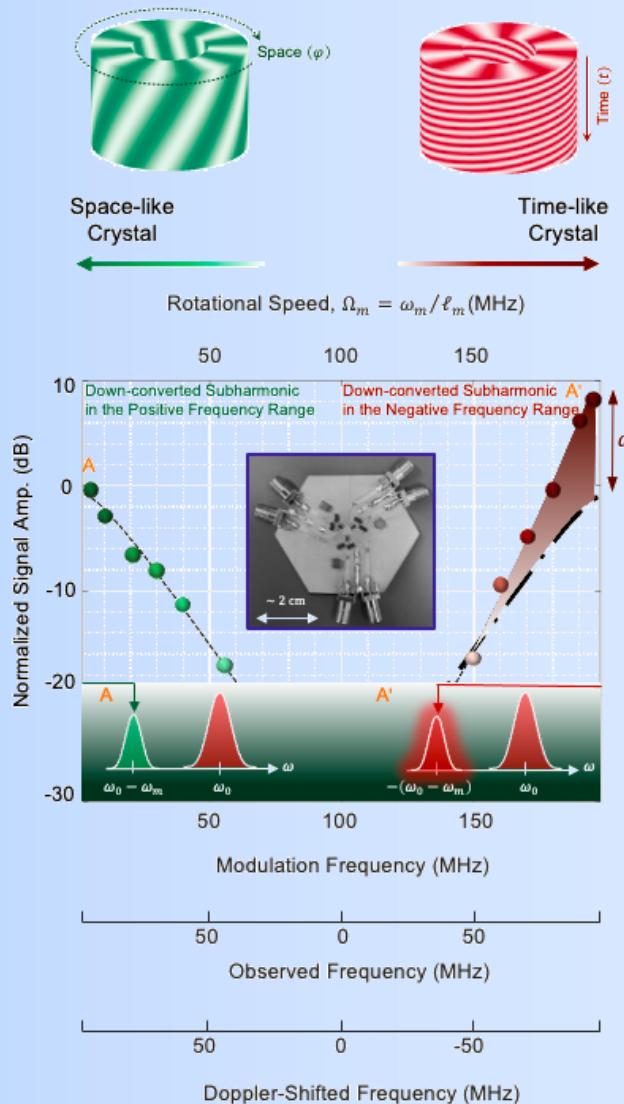
EXPERIMENTAL SETUP



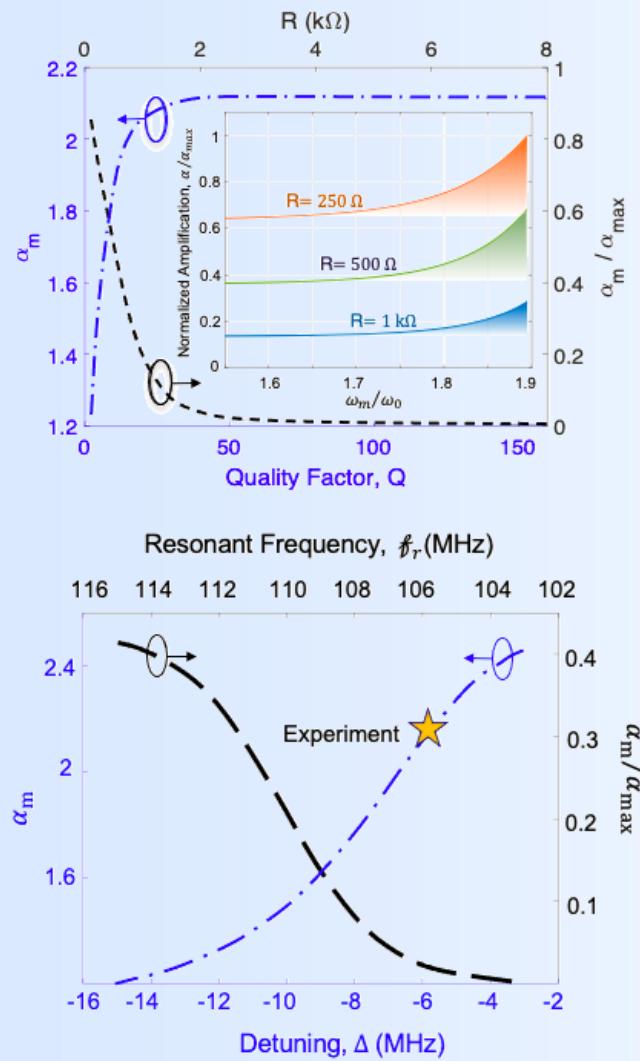
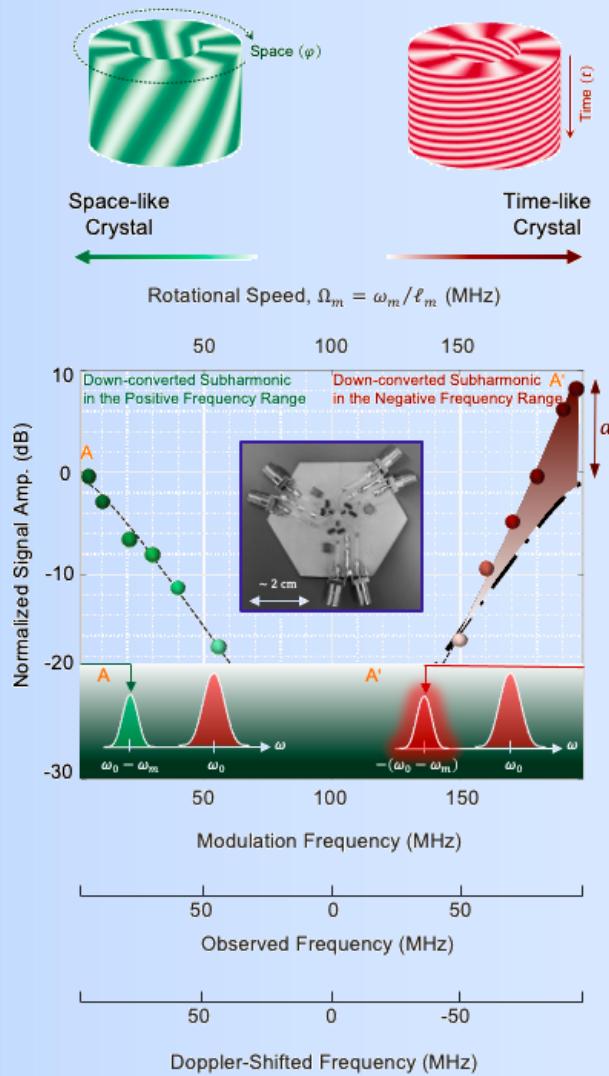
OBSERVATION OF SYNTHETIC PENROSE SUPER-RADIANCE



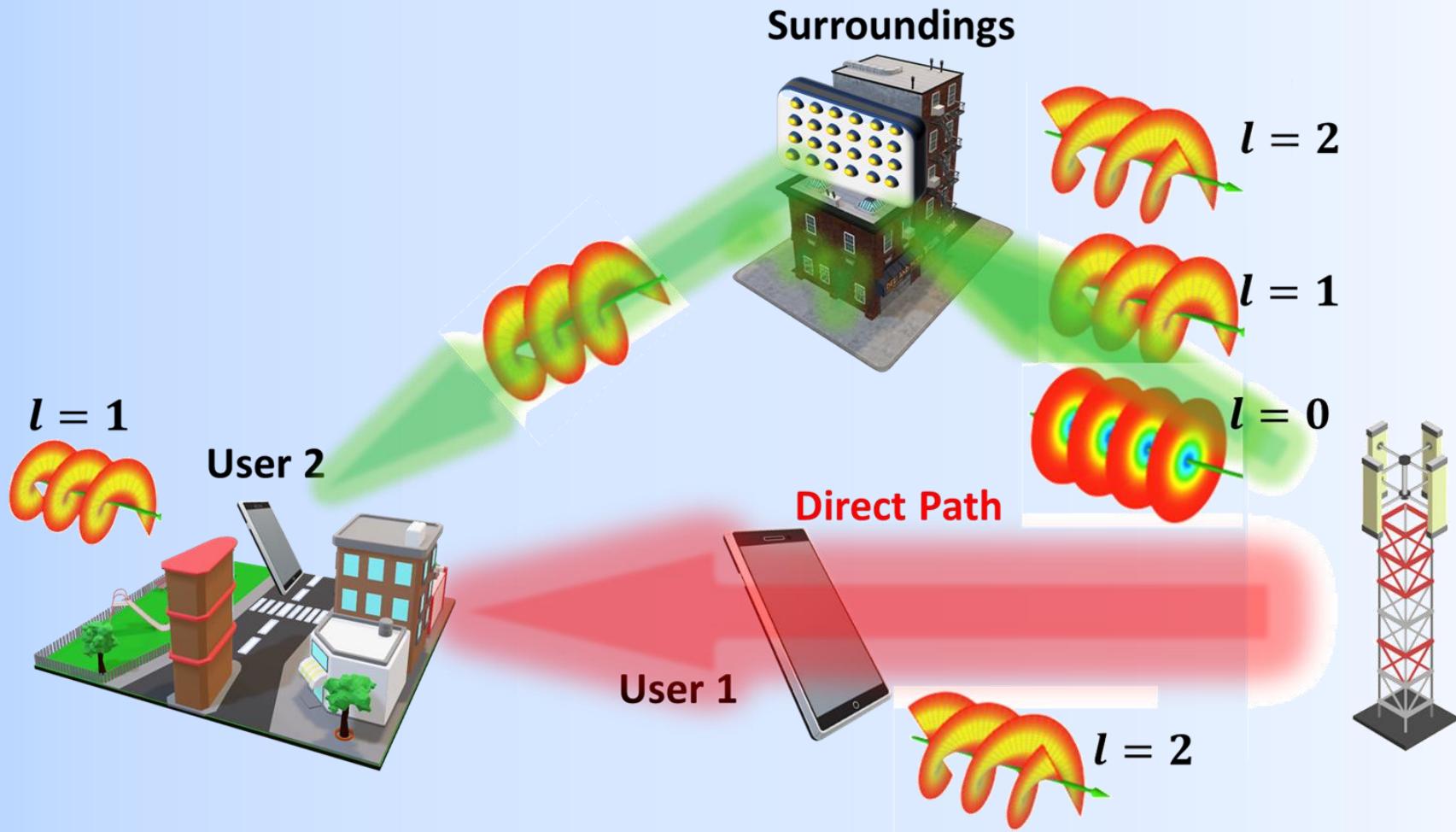
OBSERVATION OF SYNTHETIC PENROSE SUPER-RADIANCE



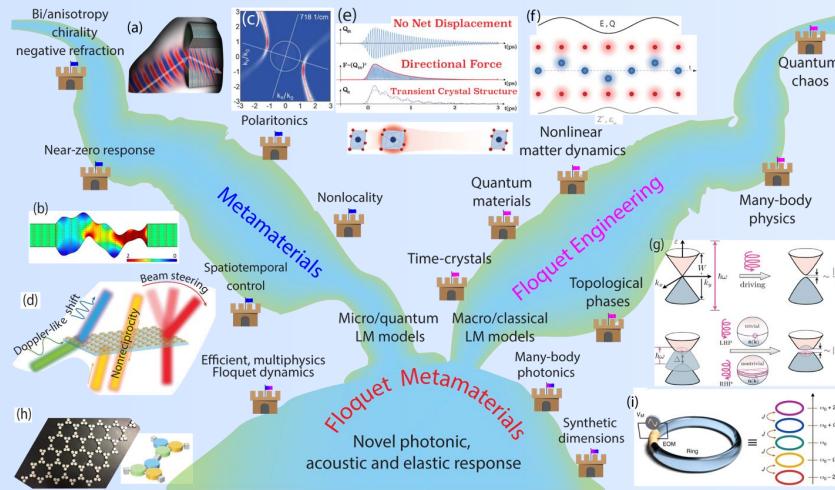
OBSERVATION OF SYNTHETIC PENROSE SUPER-RADIANCE



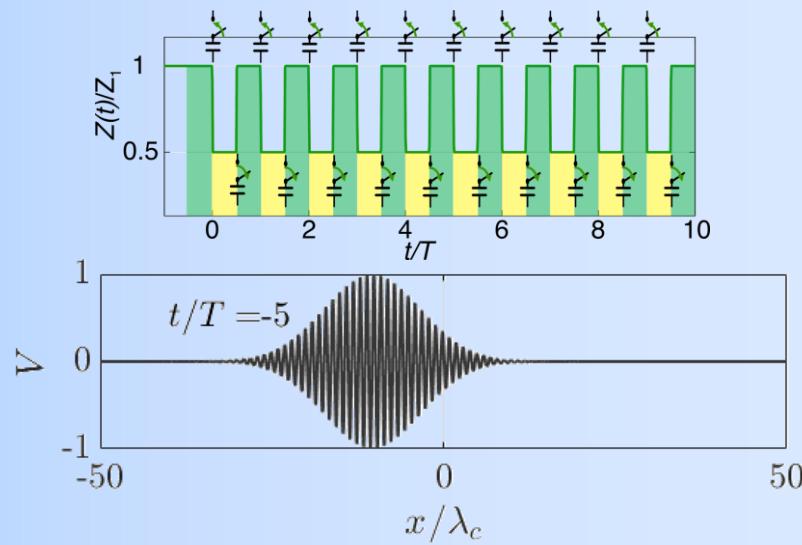
PENROSE SUPER-RADIANCE FOR 6G COMMUNICATIONS



TIME METAMATERIALS



Passive time crystals



Synthetic Penrose super-radiance

