

TOPOLOGICAL EM: ROBUST, BROADBAND ONE-WAY SIGNAL TRANSPORT

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The University of Texas at Austin

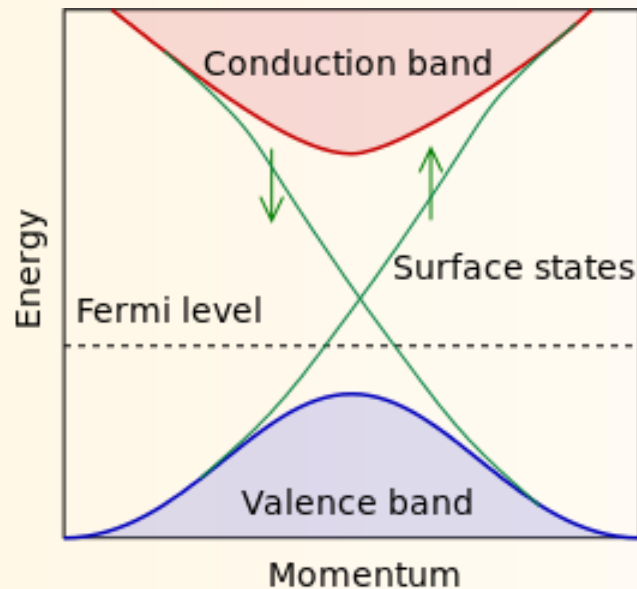
<http://users.ece.utexas.edu/~aalu>, alu@mail.utexas.edu

Supported by Dr. Arje Nachman's program with AFOSR grant No. FA9550-13-1-0204

ELECTRONIC TOPOLOGICAL INSULATORS

Symmetry-protected conducting surface states supported by topological order

Quantum Hall effect

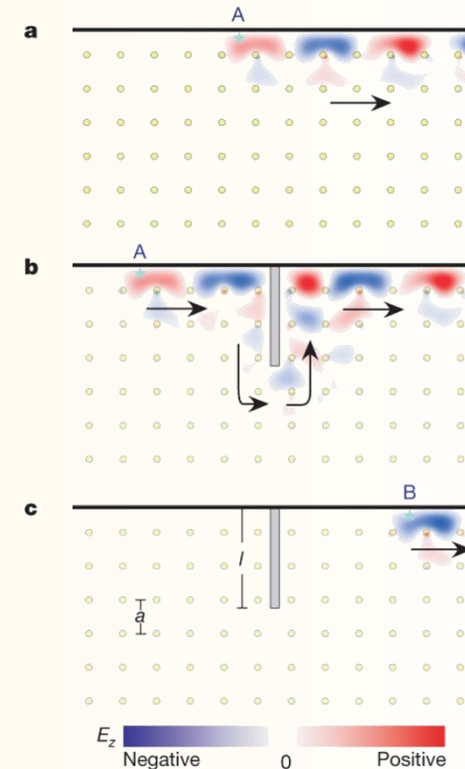
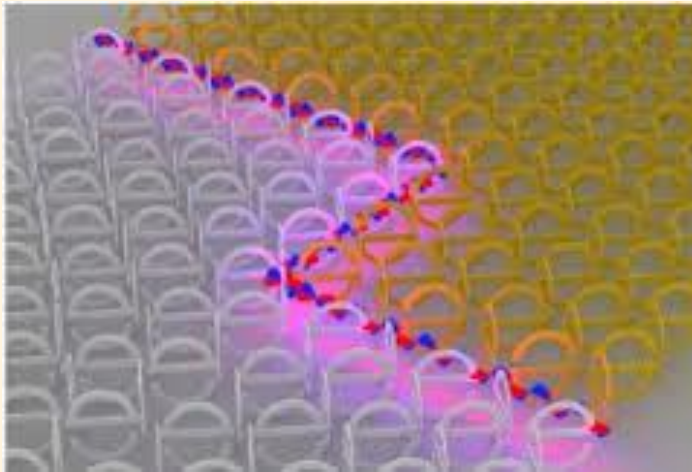


C. L. Kane, E. J. Mele, *Phys. Rev. Lett.* **95**, 146802 (2005)

M. Z. Hasan, C. L. Kane, *Rev. Mod. Phys.* **82**, 3045 (2010)

PHOTONIC TOPOLOGICAL INSULATORS

Weak/strong topological protection determined by time-reversal symmetry



Z Wang, Y. Chong, J. D. Joannopoulos, M. Soljačić, *Nature* **461**, 772 (2009)

A. P. Slobozhanyuk, A. N. Poddubny, A. E. Miroshnichenko, P. A. Belov, and Y. S. Kivshar, *Phys. Rev. Lett.* **114**, 123901 (2015)

NON-RECIPROCITY WITH MAGNETIC MATERIALS

Lorentz reciprocity theorem

$$\iiint \mathbf{J}_1 \cdot \mathbf{E}_2 dV = \iiint \mathbf{J}_2 \cdot \mathbf{E}_1 dV$$

$$\begin{aligned} \bar{\bar{\epsilon}} &= \bar{\bar{\epsilon}}^T \\ \bar{\bar{\mu}} &= \bar{\bar{\mu}}^T \end{aligned}$$

Time-invariant materials

Linear materials

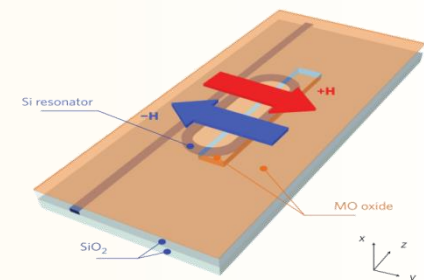
Static magnets



Weak effects →
Massive devices



Difficult to integrate



Bi, Nature Photon. 5, 758

BREAKING RECIPROCITY CONSTRAINTS

Lorentz reciprocity theorem

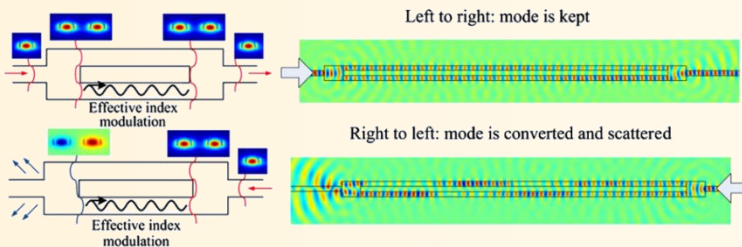
$$\iiint \mathbf{J}_1 \cdot \mathbf{E}_2 dV = \iiint \mathbf{J}_2 \cdot \mathbf{E}_1 dV$$

$$\bar{\bar{\epsilon}} = \bar{\bar{\epsilon}}^T$$

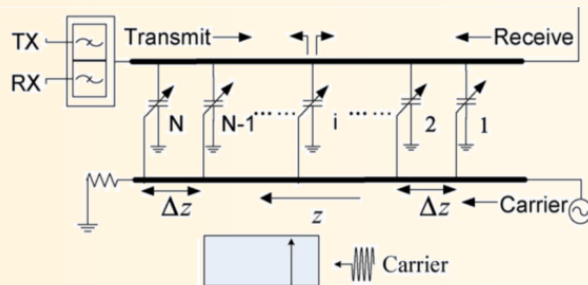
$$\bar{\bar{\mu}} = \bar{\bar{\mu}}^T$$

Time-invariant materials

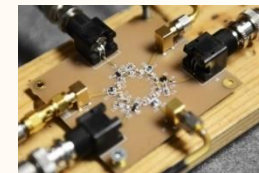
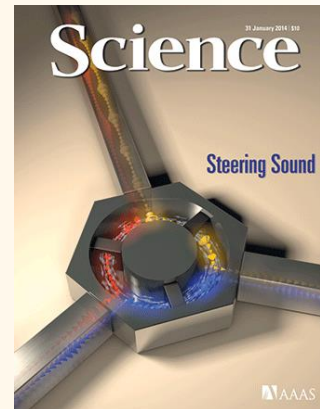
Linear materials



Lira, *PRL* **109**, 033901 (2012)



Qin, *IEEE TAP* **62**, 2260 (2014)



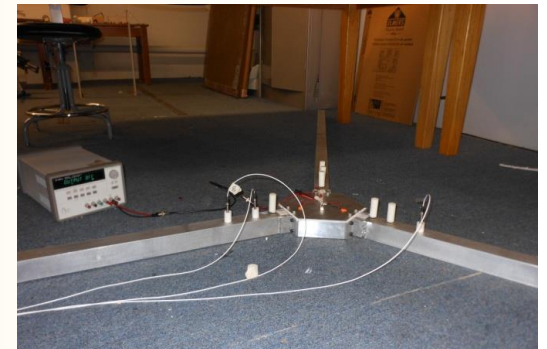
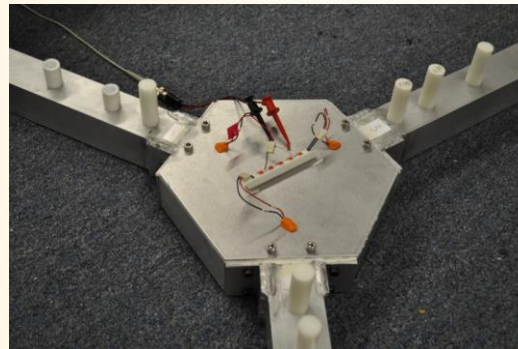
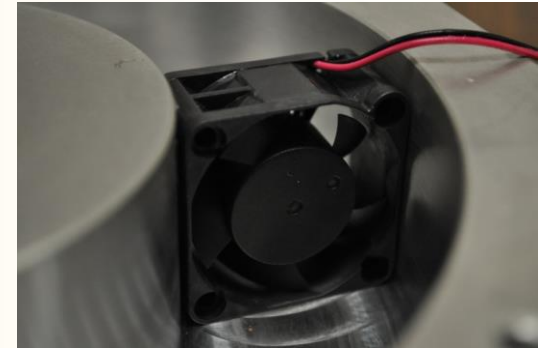
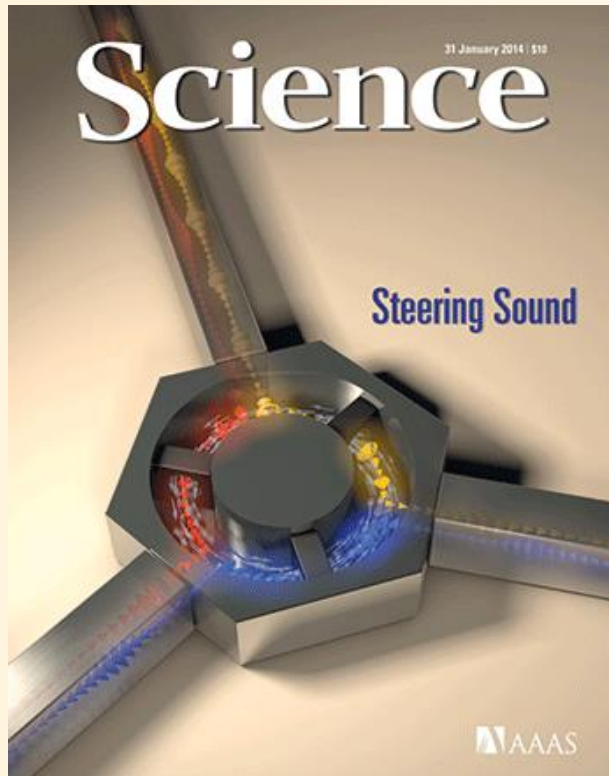
Sounas, *Nature Commun.* **4**, 2407 (2013)

D. L. Sounas, *ACS Photonics* **1**, 198 (2014)

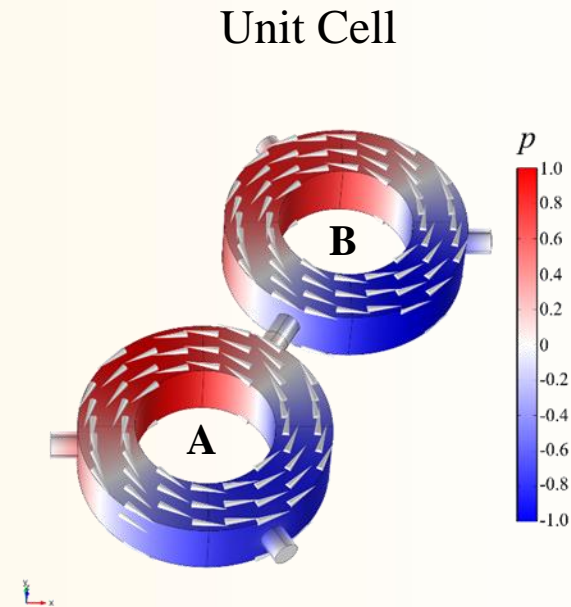
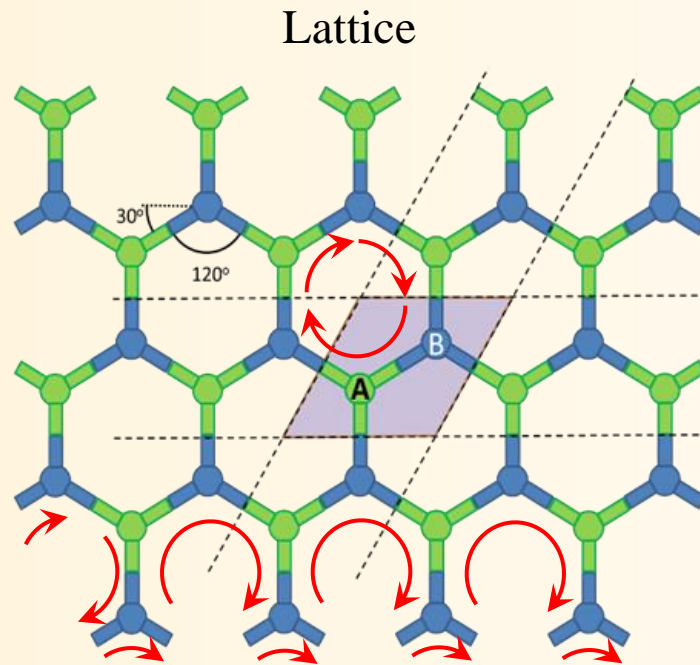
Fleury, *Science* **343**, 516 (2014)

Estep, *Nature Phys.* **10**, 923 (2014)

A CIRCULATOR FOR SOUND

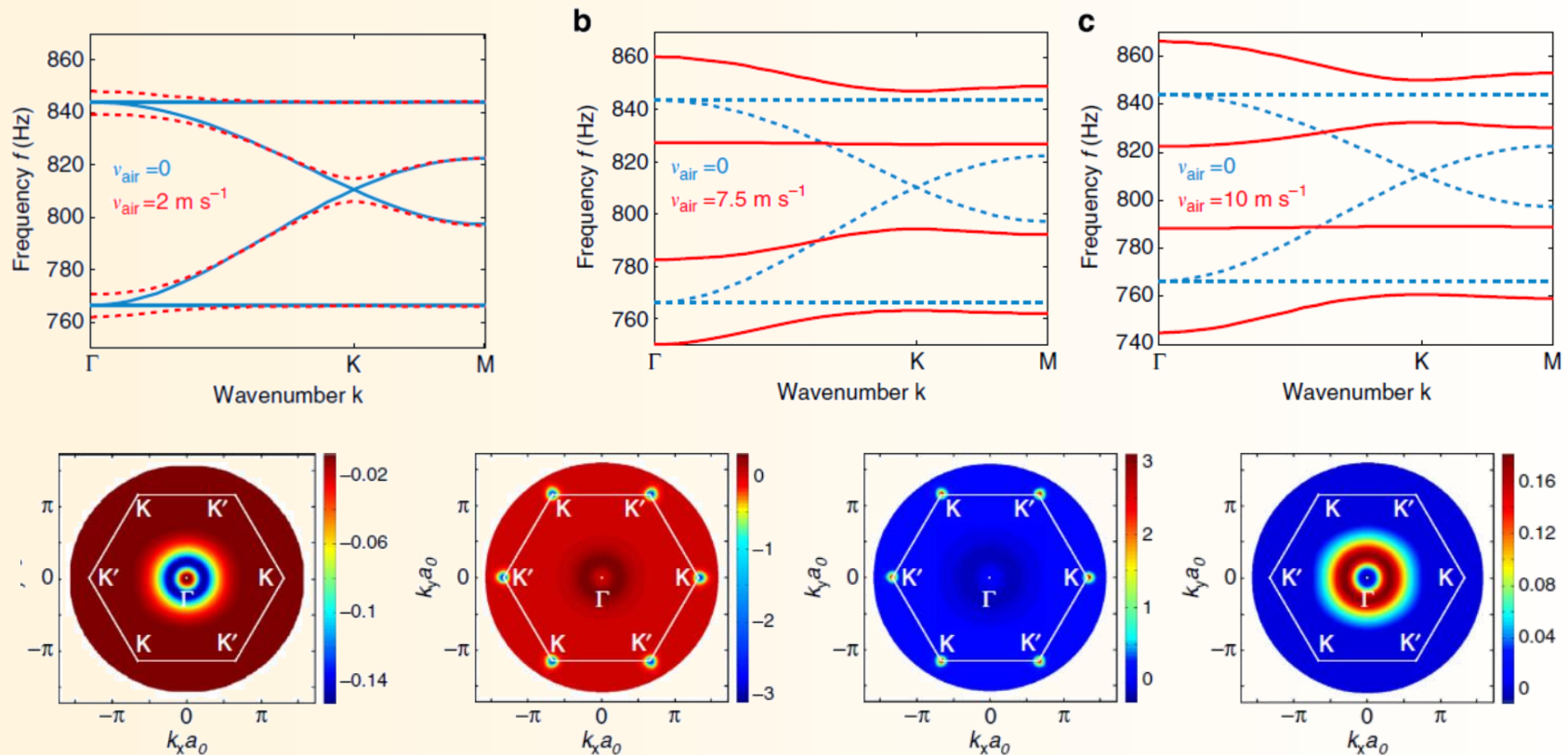


MOLDING THE TOPOLOGY OF THE BAND DIAGRAM



A. Khanikaev, R. Fleury, H. Mousavi, and A. Alù, *Nat. Comm.* **6**, 8260 (2015)

NON-TRIVIAL TOPOLOGY OF THE BAND STRUCTURE



$$C = \frac{1}{2\pi} \int_{\text{FBZ}} \Omega(\mathbf{k}) d^2\mathbf{k}$$

$$\Omega(\mathbf{k}) = \partial_{k_x} A_y - \partial_{k_y} A_x$$

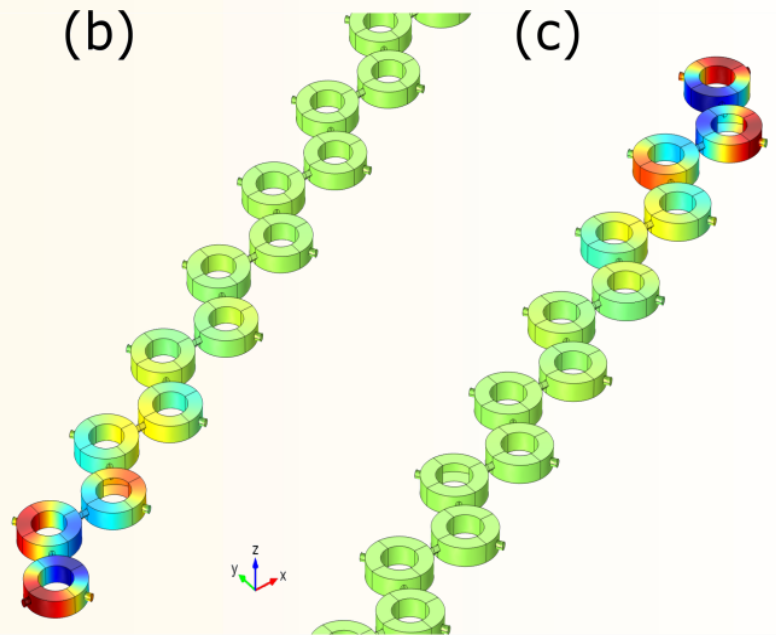
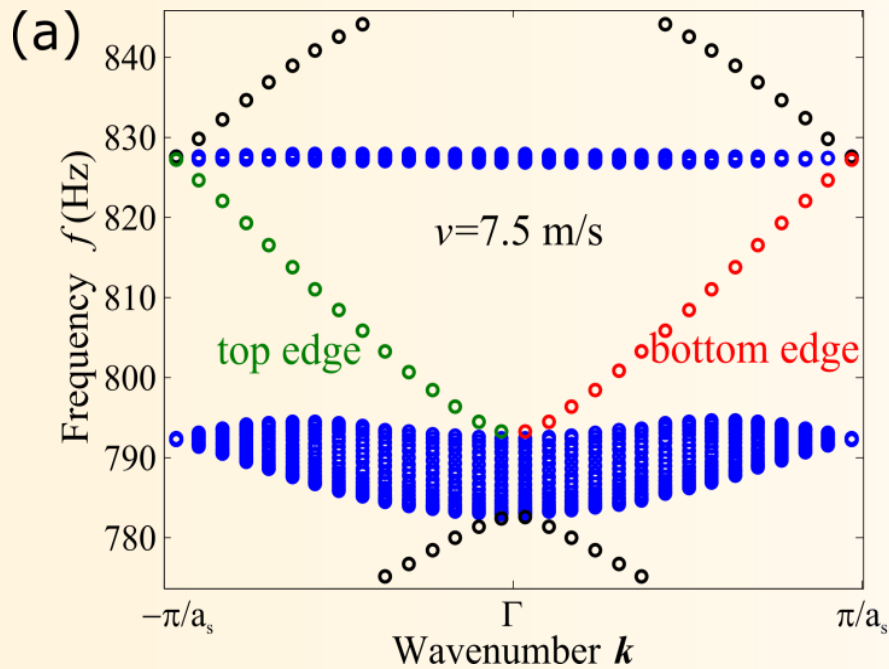
$$\mathbf{A}_{\mathbf{k}} = -i \langle p | \partial_{\mathbf{k}} | p \rangle$$

$C = \{-1, 0, 0, 1\}$ clockwise

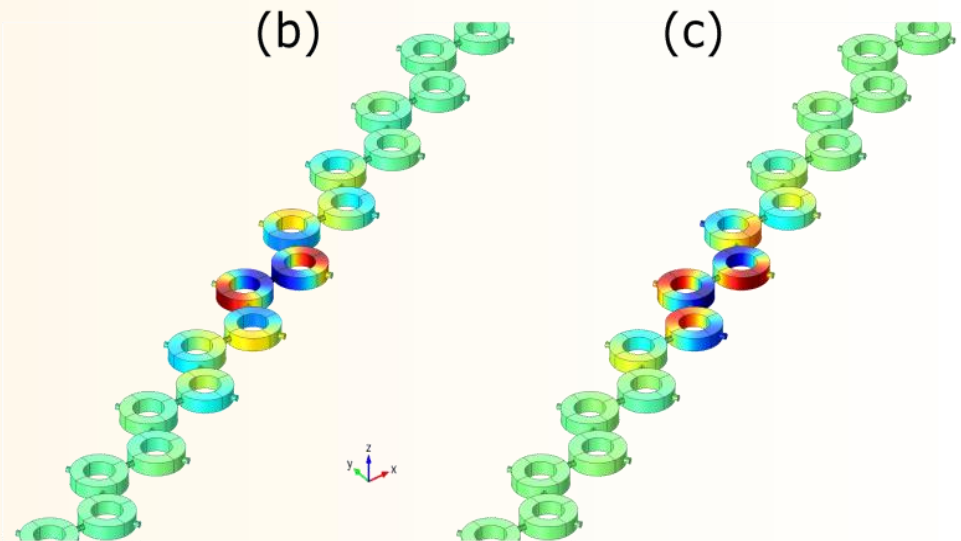
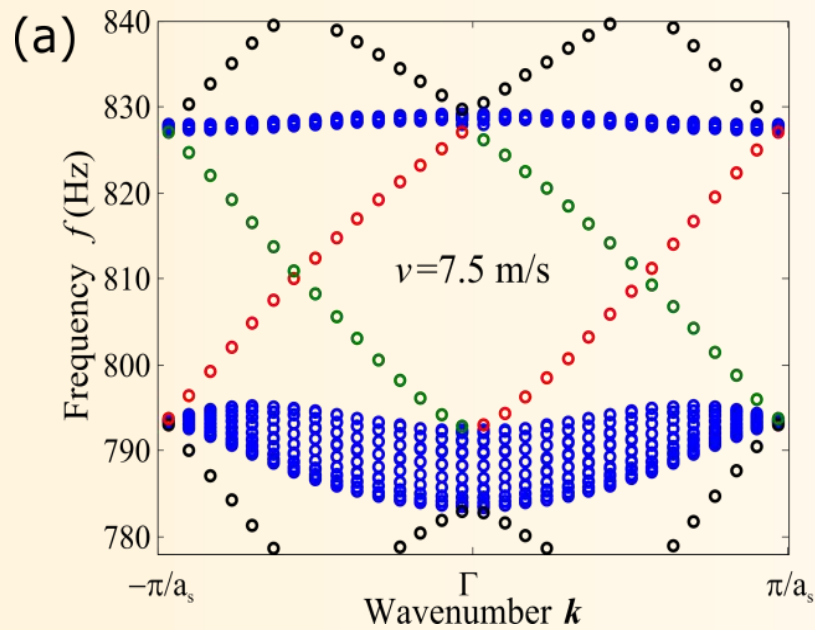
$C = \{1, 0, 0, -1\}$ counterclockwise

A. Khanikaev, R. Fleury, H. Mousavi, and A. Alù, *Nat. Comm.* **6**, 8260 (2015)

TOPOLOGICAL METAMATERIALS

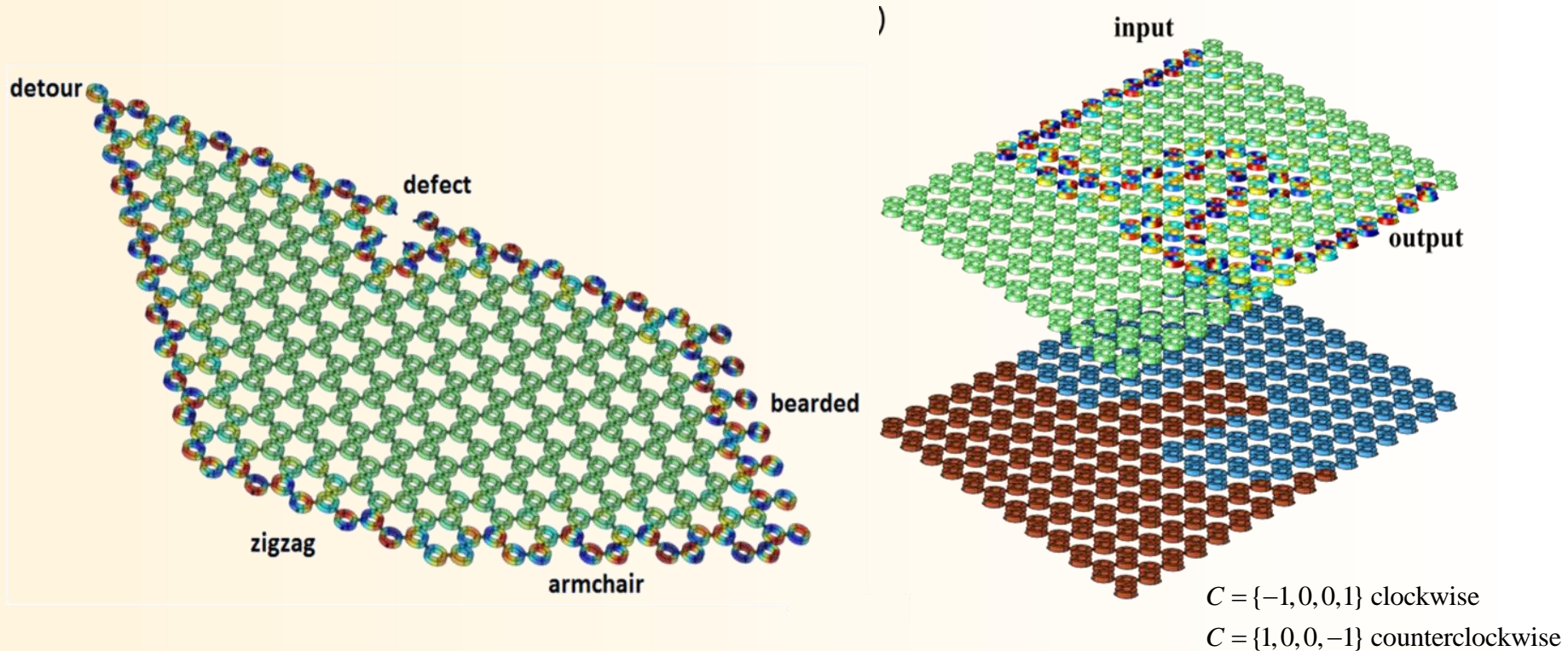


TOPOLOGICAL METAMATERIALS



TOPOLOGICAL METAMATERIALS

Reconfigurable waveguides, broadband isolators...



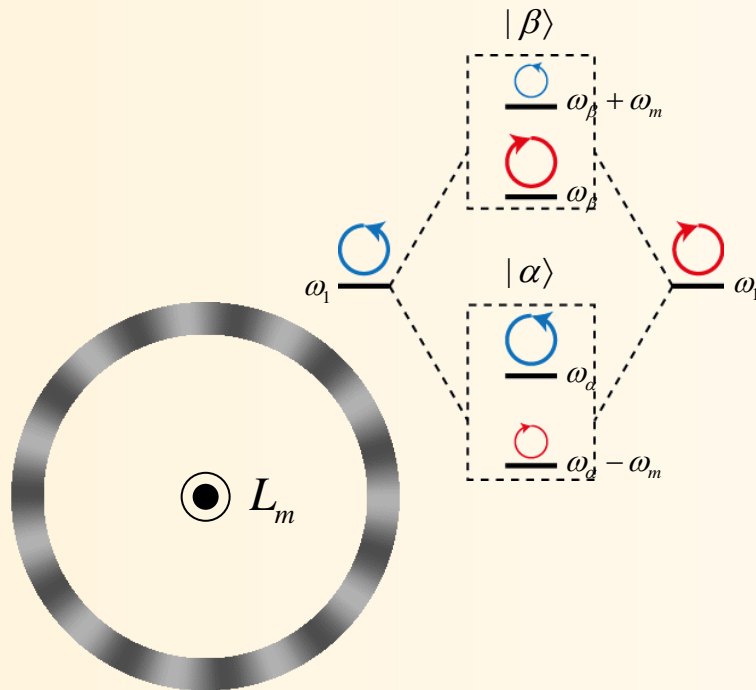
A. Khanikaev, R. Fleury, H. Mousavi, and A. Alù, *Nat. Comm.* **6**, 8260 (2015)

ANGULAR-MOMENTUM-BIAS IN NANOPHOTONICS

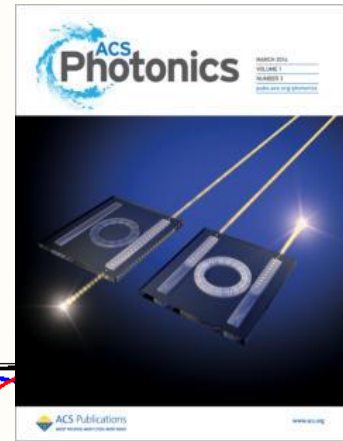
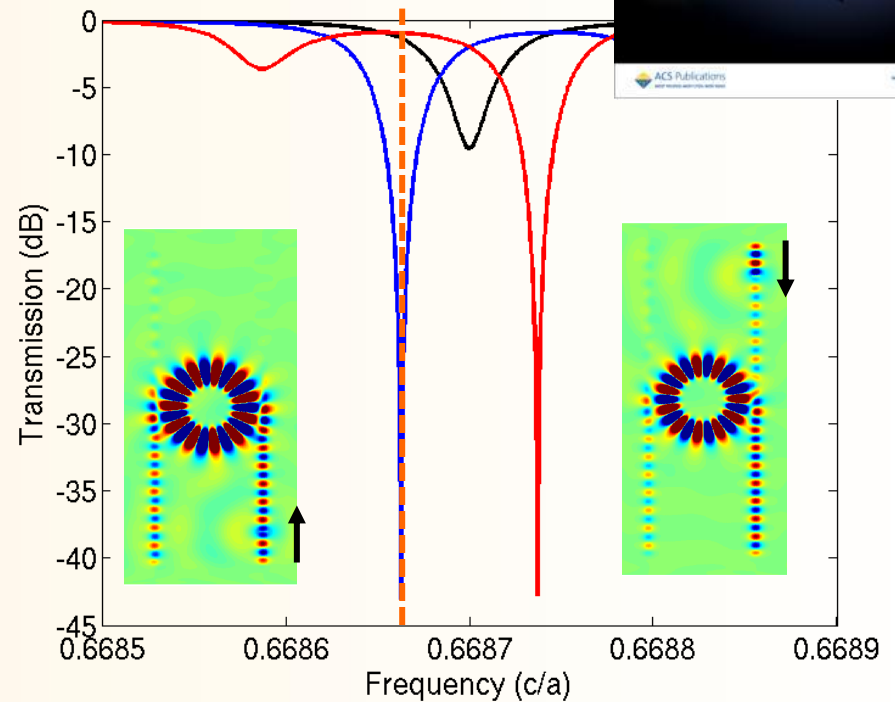
$$\Delta\epsilon_m = 5 \times 10^{-4} \epsilon !$$

$$Q \sim 7,000$$

$$Q\Delta\epsilon_m \approx 1$$

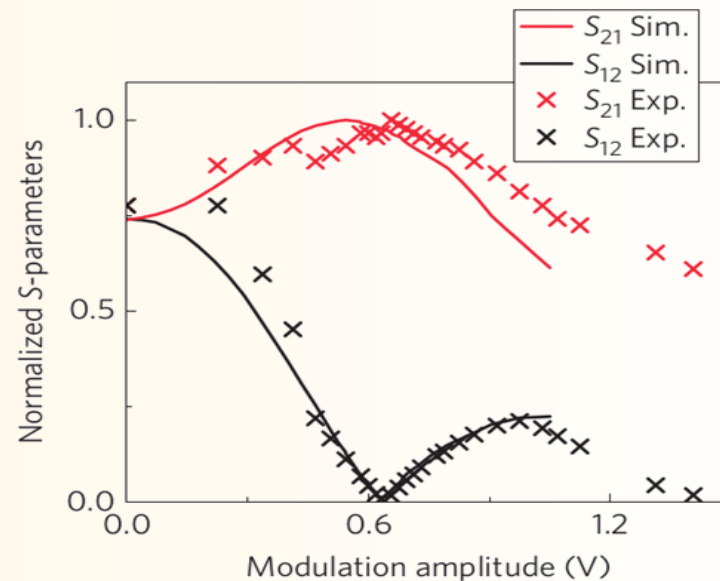
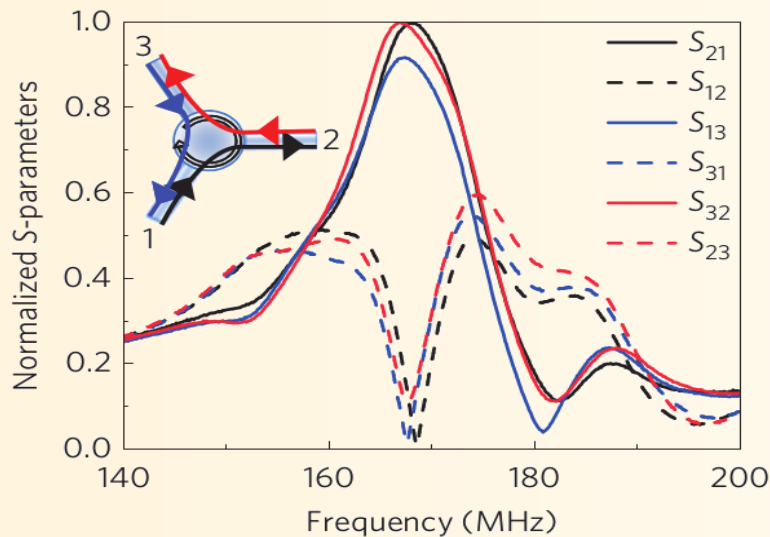
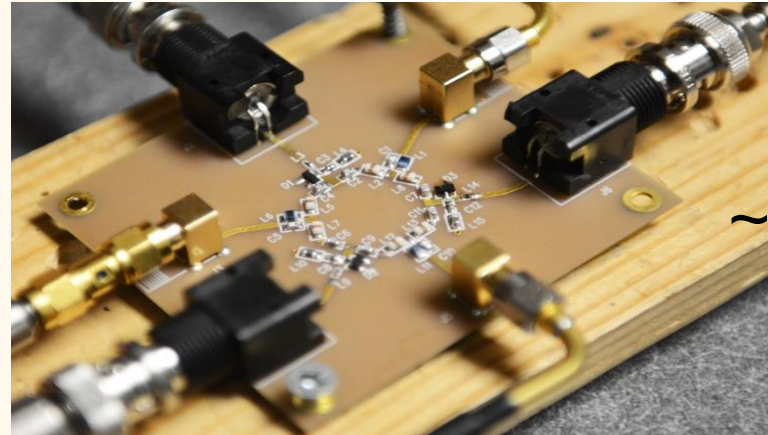
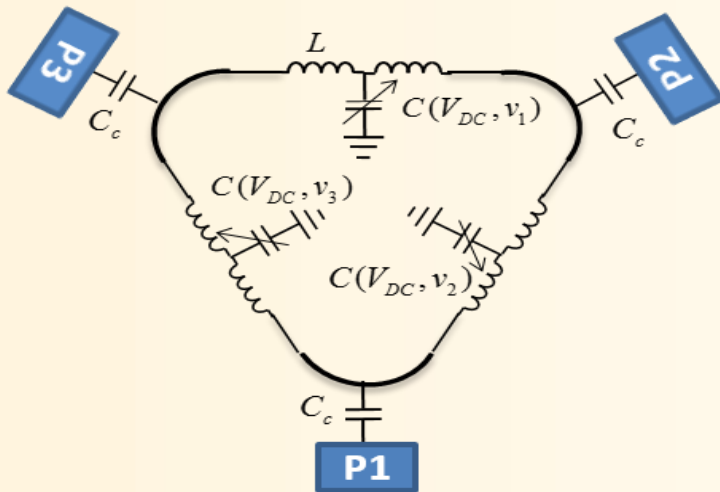


$$\Delta\epsilon(\varphi, t) = \Delta\epsilon_m \cos(\omega_m t - L_m \varphi)$$



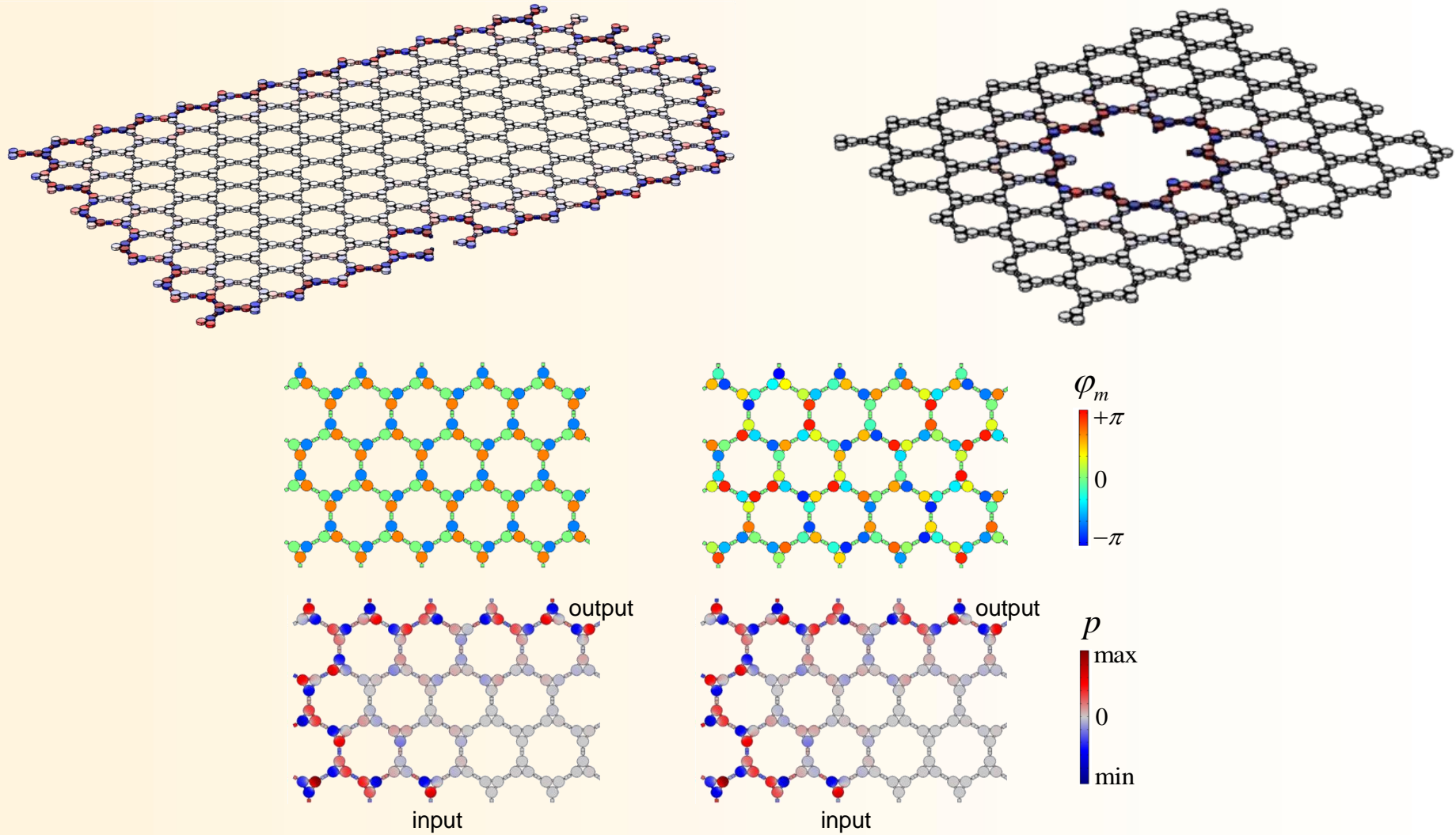
D. L. Sounas, A. Alù, *ACS Photonics* 1, 198 (2014)

RF MAGNET-LESS INTEGRATED CIRCULATOR



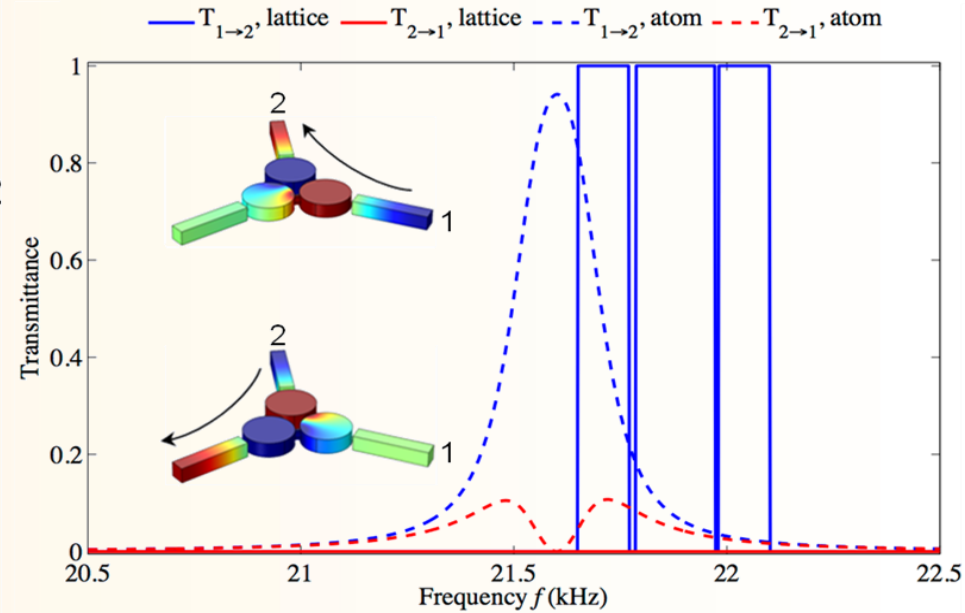
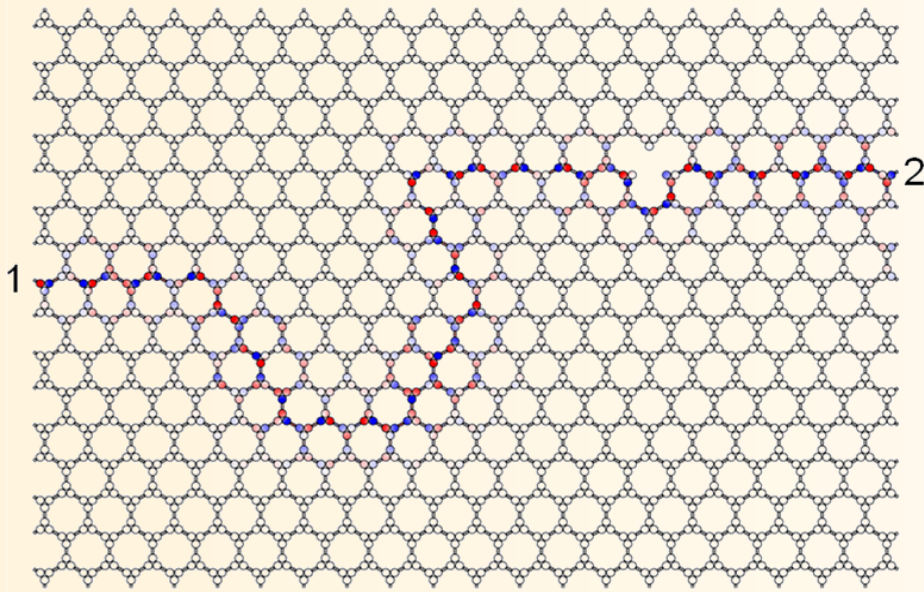
N. A. Estep*, D. L. Sounas*, J. Soric, and A. Alù, *Nature Phys.*, **10**, 923 (2014)

FLOQUET TOPOLOGICAL INSULATORS FOR LIGHT



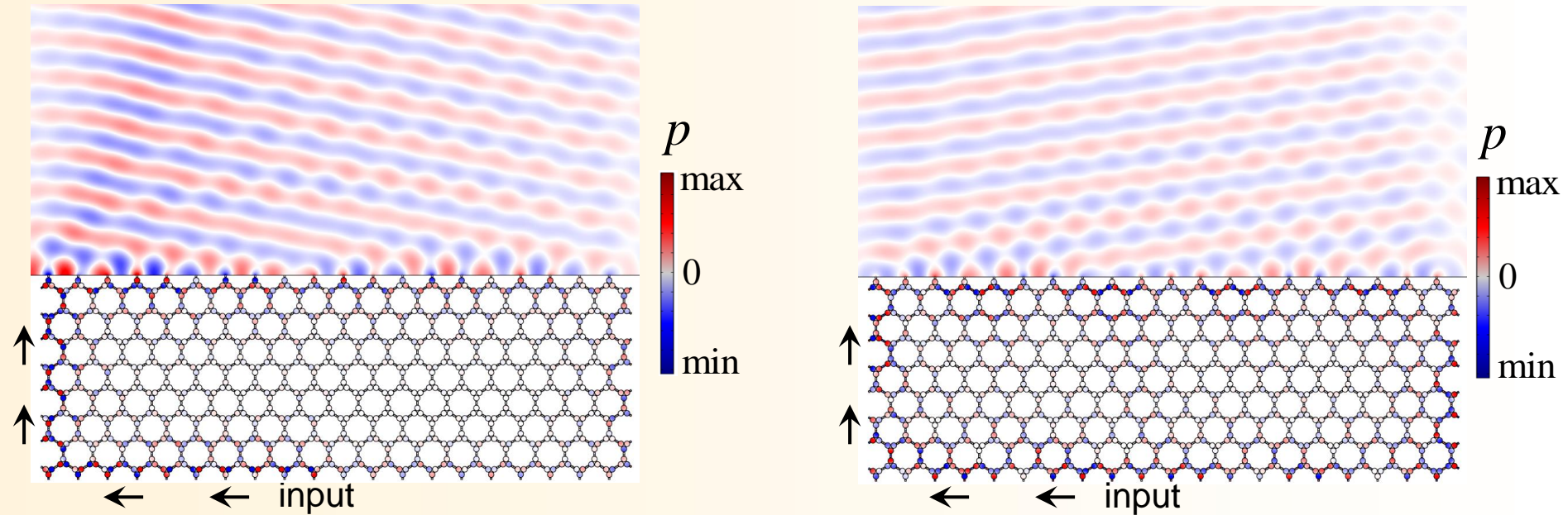
R. Fleury, A.B. Khanikaev and A. Alù, *Nature Communications*, **7**, 11744 (2016)

CONTINUOUS BANDWIDTH OF ISOLATION



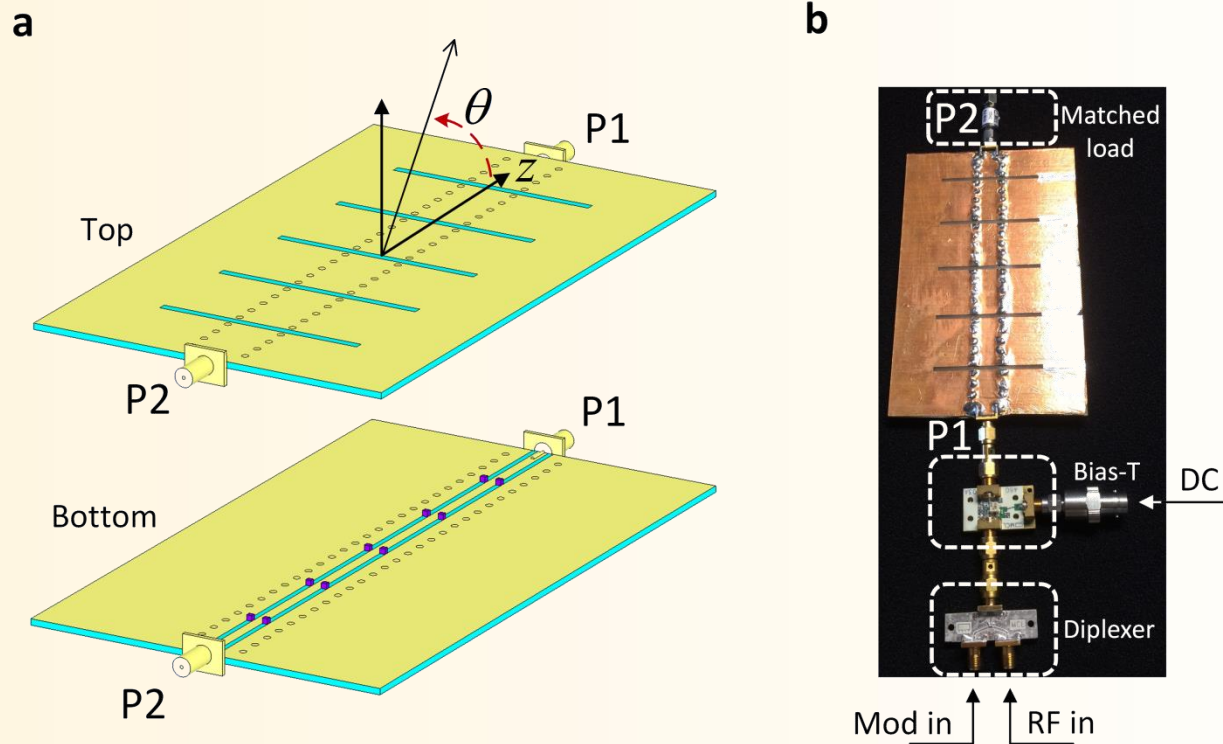
R. Fleury, A.B. Khanikaev and A. Alù, *Nature Communications*, **7**, 11744 (2016)

TOPOLOGICALLY PROTECTED LEAKY-WAVE RADIATION



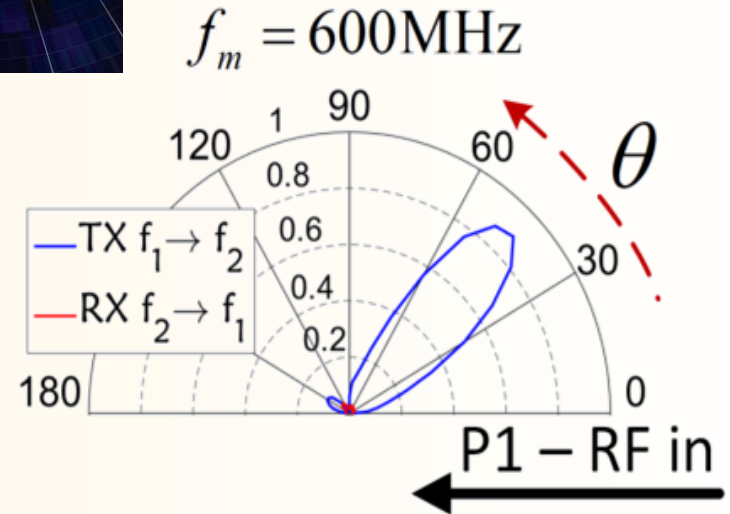
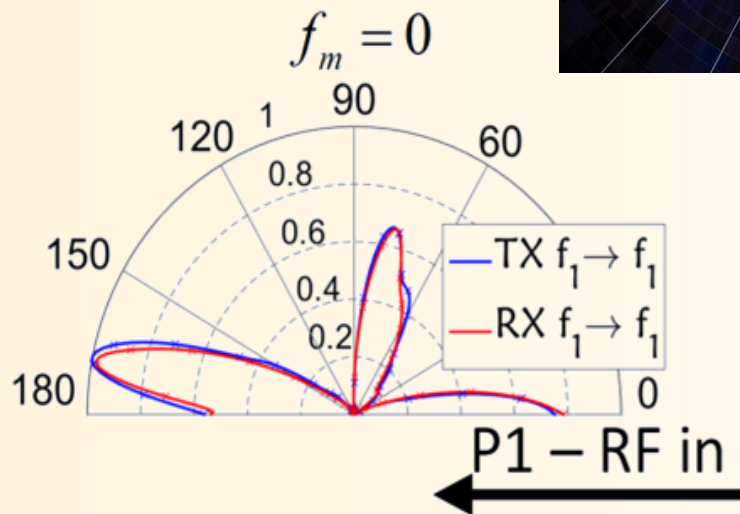
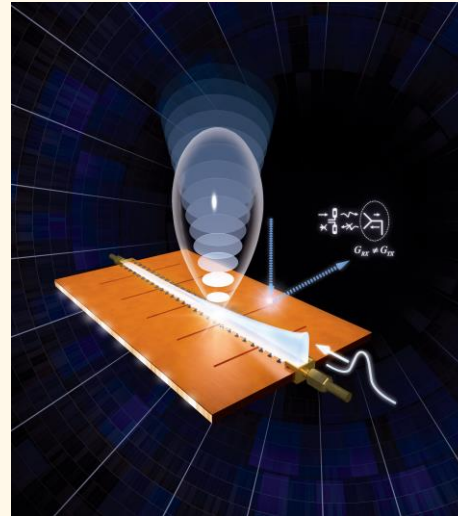
R. Fleury, A.B. Khanikaev and A. Alù, *Nature Communications*, **7**, 11744 (2016)

NON-RECIPROCAL METASURFACES BASED ON ST MODULATION



Y. Hadad, J. C. Soric, and A. Alù, *PNAS* **113**, 33471 (2016)

NON-RECIPROCAL METASURFACES BASED ON ST MODULATION



Y. Hadad, J. C. Soric, and A. Alù, *PNAS* **113**, 33471 (2016)

NON-RECIPROCITY BASED ON NON-LINEAR EFFECTS

Lorentz reciprocity theorem

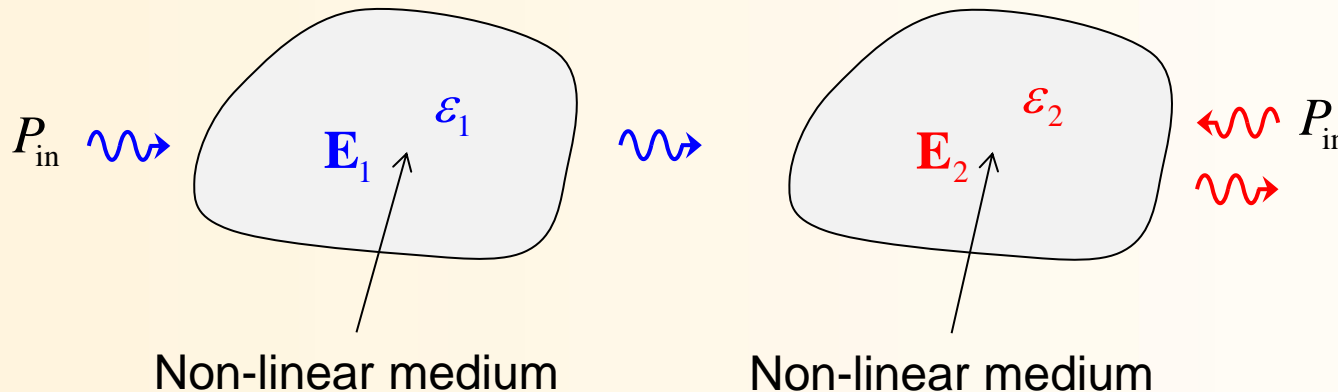
$$\iiint \mathbf{J}_1 \cdot \mathbf{E}_2 dV = \iiint \mathbf{J}_2 \cdot \mathbf{E}_1 dV$$

$$\bar{\bar{\epsilon}} = \bar{\bar{\epsilon}}^T$$

$$\bar{\bar{\mu}} = \bar{\bar{\mu}}^T$$

Time-invariant materials

Linear materials



Chi-3 non-linearity

$$\Delta\epsilon_{\text{NL}} = \chi^{(3)} |\mathbf{E}|^2$$

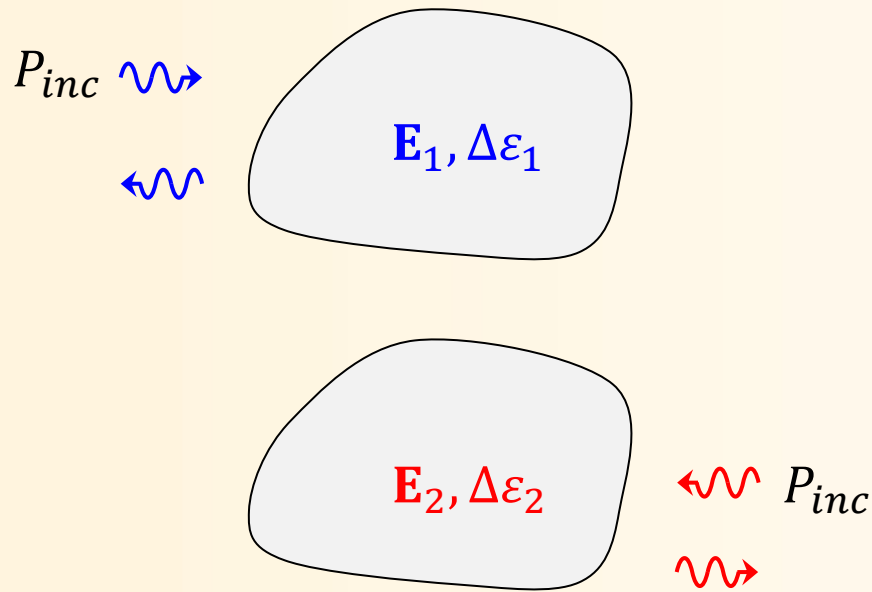
P. Saboo, J. Joseph, *Appl. Opt.* **52**, 8252–8257 (2013)

L. Fan, et al., *Opt. Lett.* **38**, 1259–1261 (2013)

Y. Shi, Z. Yu, S. Fan, *Nature Photon.* **9**, 388–392 (2015)

A. M. Mahmoud, A. Davoyan, N. Engheta, *Nature Comm.* **6**, 8359 (2015)

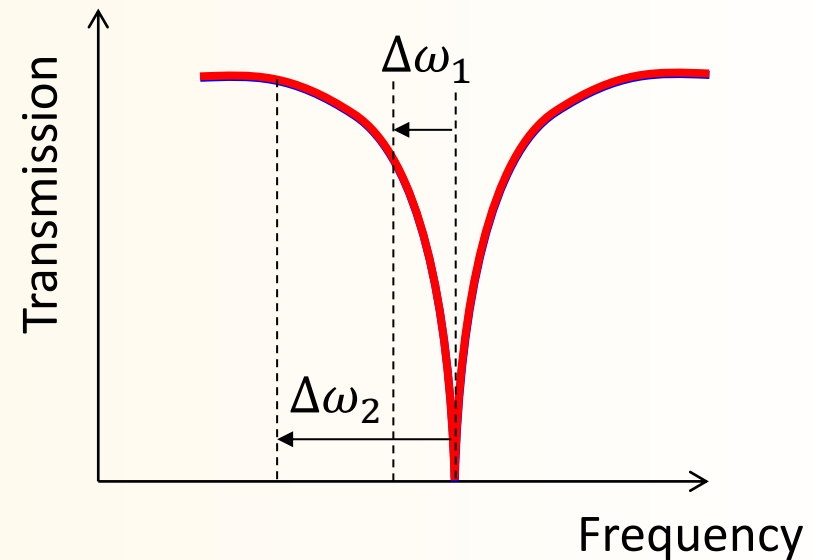
BASIC OPERATION – SINGLE RESONANCE



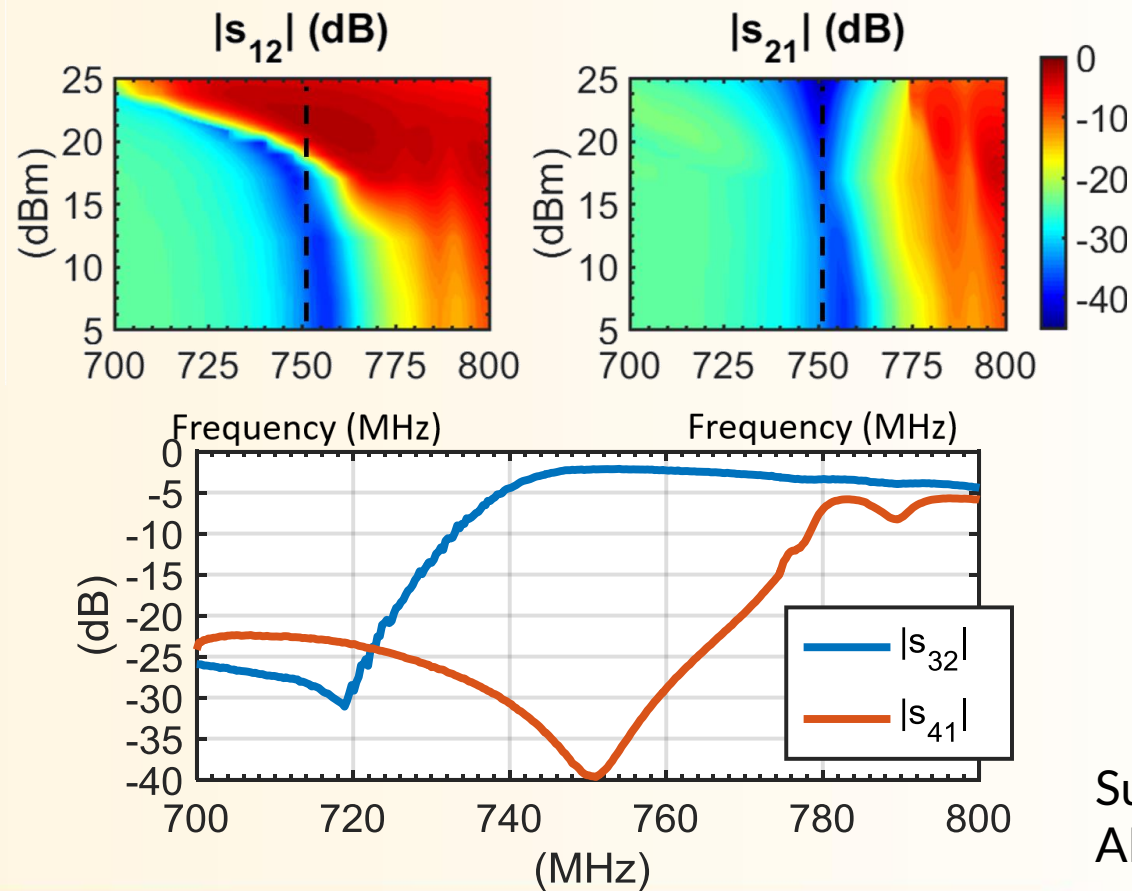
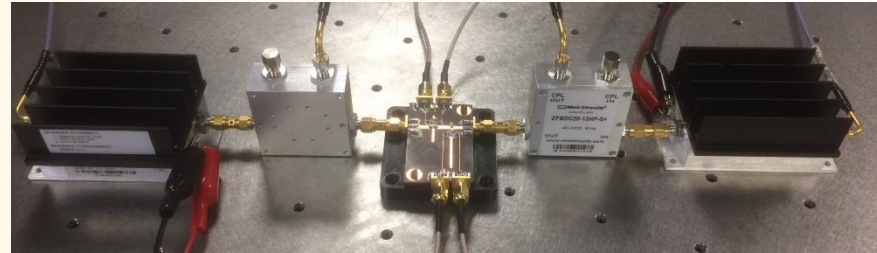
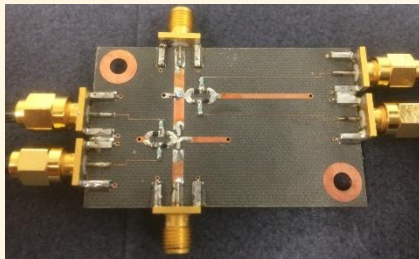
$$\kappa = \frac{|\mathbf{E}_2|^2}{|\mathbf{E}_1|^2} > 1$$

$$\epsilon_{NL} = \epsilon + \chi^{(3)}|\mathbf{E}|^2$$

$$\Delta\omega = -\omega_0 \frac{\int \Delta\epsilon |\mathbf{E}_0|^2 dV}{\int \epsilon |\mathbf{E}_0|^2 dV}$$

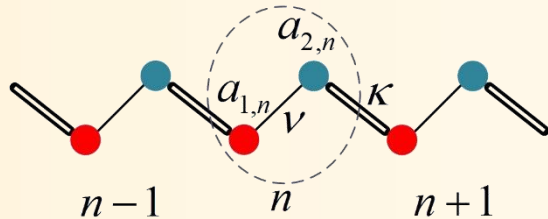


EXPERIMENTAL RESULTS AT RF



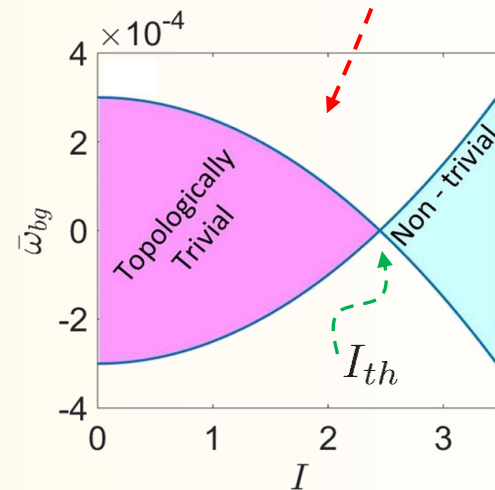
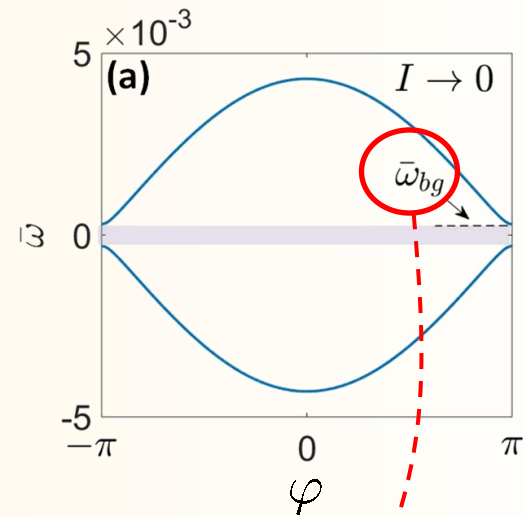
Supported by
AFOSR (A. Sayir)

NONLINEARITY-INDUCED TOPOLOGICAL TRANSITIONS

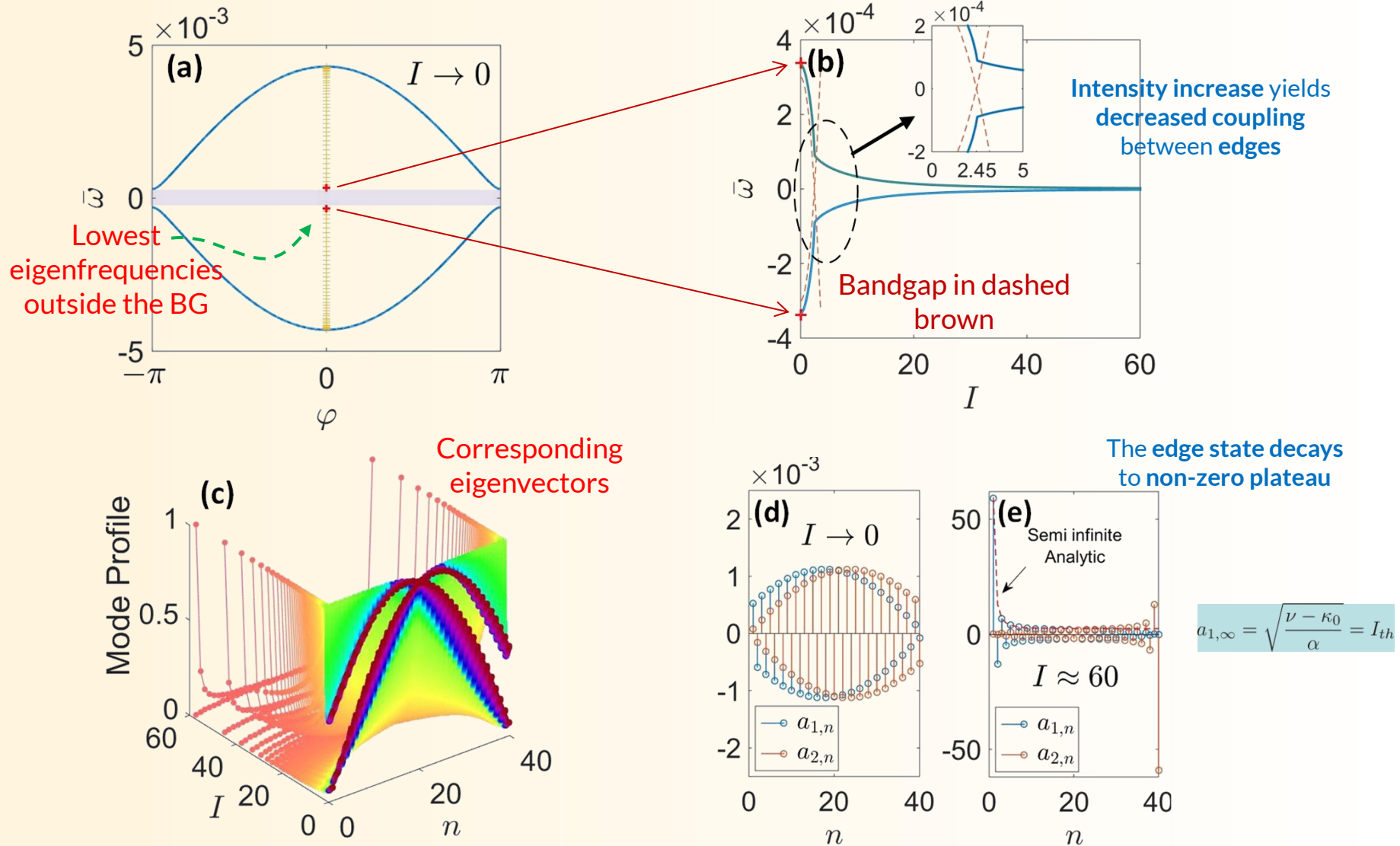


$$i \frac{d\Psi_n}{dt} = \Omega \Psi_n + \mathbf{K}_m \Psi_{n-1} + \mathbf{K}_p \Psi_{n+1}$$

$$\bar{\omega}_{bg} = \pm |\nu - \kappa_0 - \alpha I^2| \xrightarrow{\nu > \kappa_0} I_{th} = \sqrt{\frac{\nu - \kappa_0}{\alpha}}$$

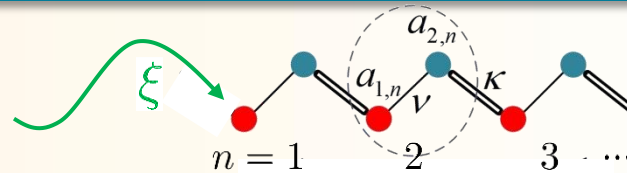


FINITE ARRAY: 40 ELEMENT CHAIN

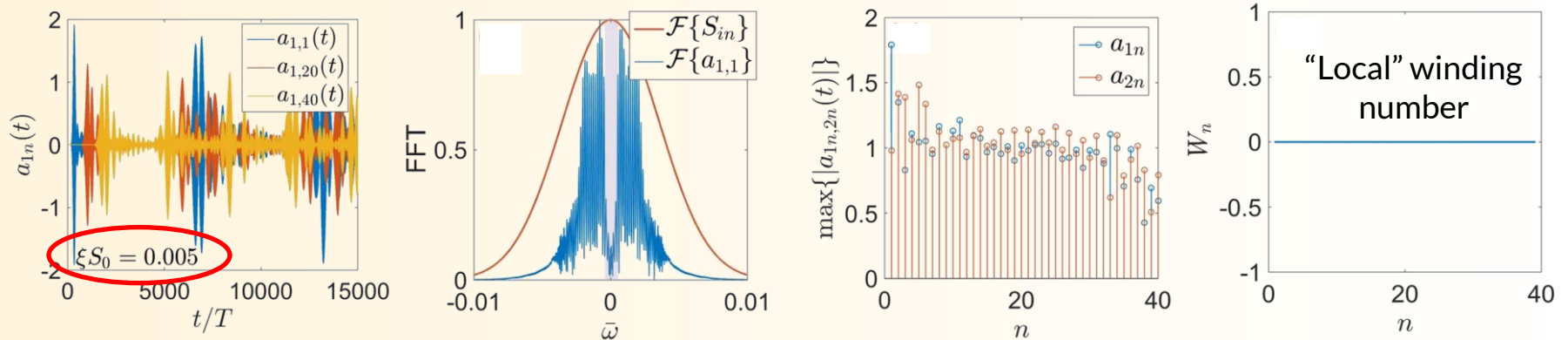


NONLINEARITY-INDUCED TOPOLOGICAL TRANSITIONS

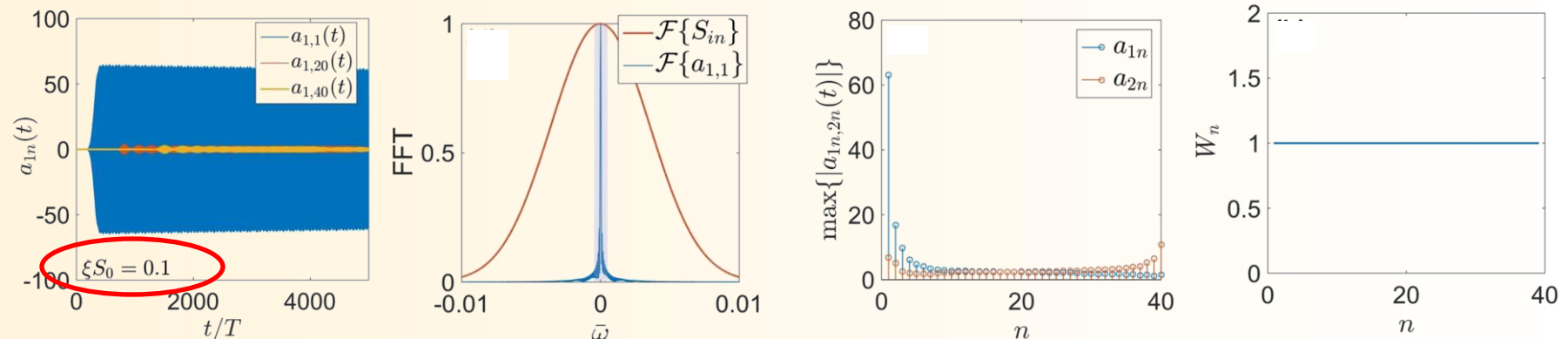
$$S_{in} = S_0 e^{i\omega t} e^{-(t-t_0)^2/T^2}$$



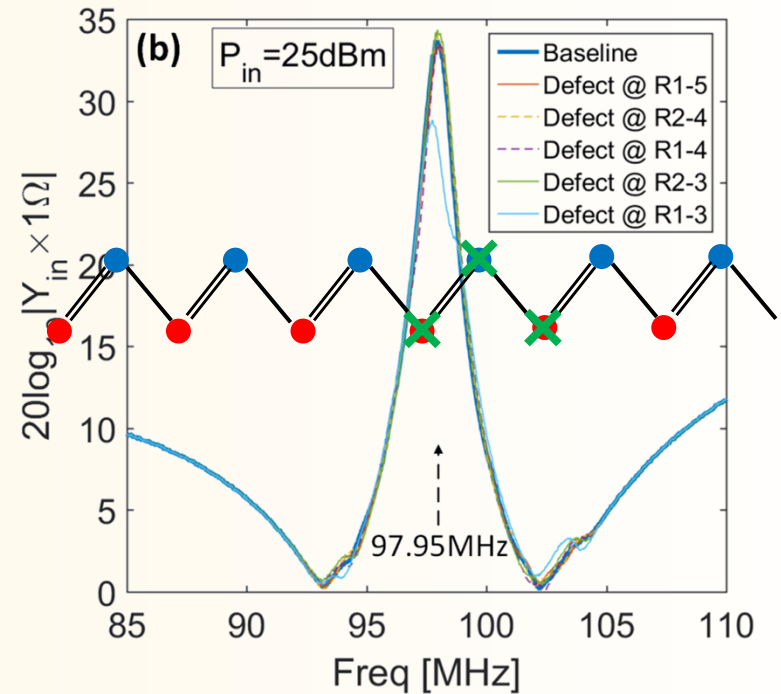
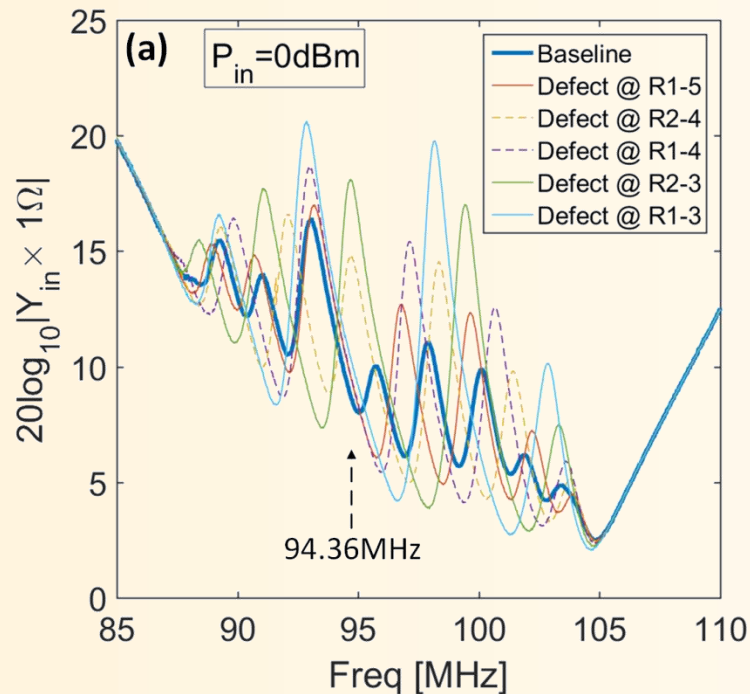
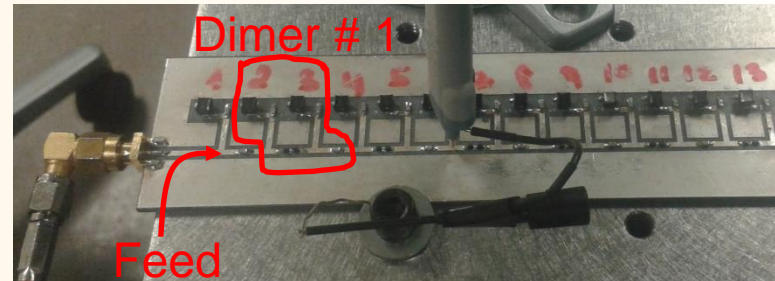
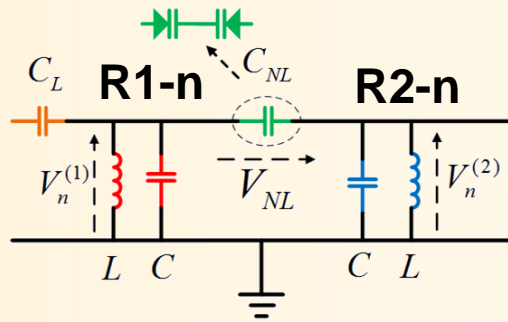
Low input intensity:



High input intensity:



NONLINEARITY-INDUCED TOPOLOGICAL TRANSITIONS



EQUIVALENT CONTINUUM MODEL

Stored energy
in dimer

$$q_n(t) = \sqrt{|a_n^{(1)}|^2 + |a_n^{(2)}|^2}$$

Relative phase
in dimer

$$\alpha_n(t) = \tan^{-1} \left(j \frac{a_n^{(2)}}{a_n^{(1)}} \right)$$



Continuum
limit

Moving
frame $\xi = x - ct$

$$q_n(t), \alpha_n(t) \mapsto q(x, t), \alpha(x, t) \mapsto q(\xi), \alpha(\xi)$$

Continuum model:

$$\frac{d}{d\xi} \begin{pmatrix} q \\ \alpha \end{pmatrix} = F(q) \begin{pmatrix} \kappa \Delta q \cos 2\alpha \\ -c - \kappa \Delta \sin 2\alpha \end{pmatrix}$$

Fixed points $\underline{x}^* = (q^*; \alpha^*)$

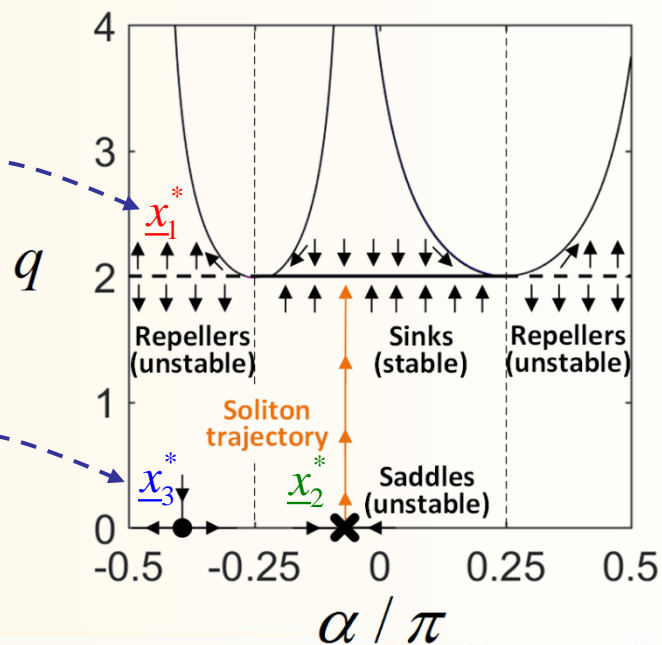
- line of unstable/stable nodes

$$\underline{x}_1^* \rightarrow q_1^* = (v_0 - \kappa) / (\kappa - v_\infty), \quad \forall \alpha_1^*$$

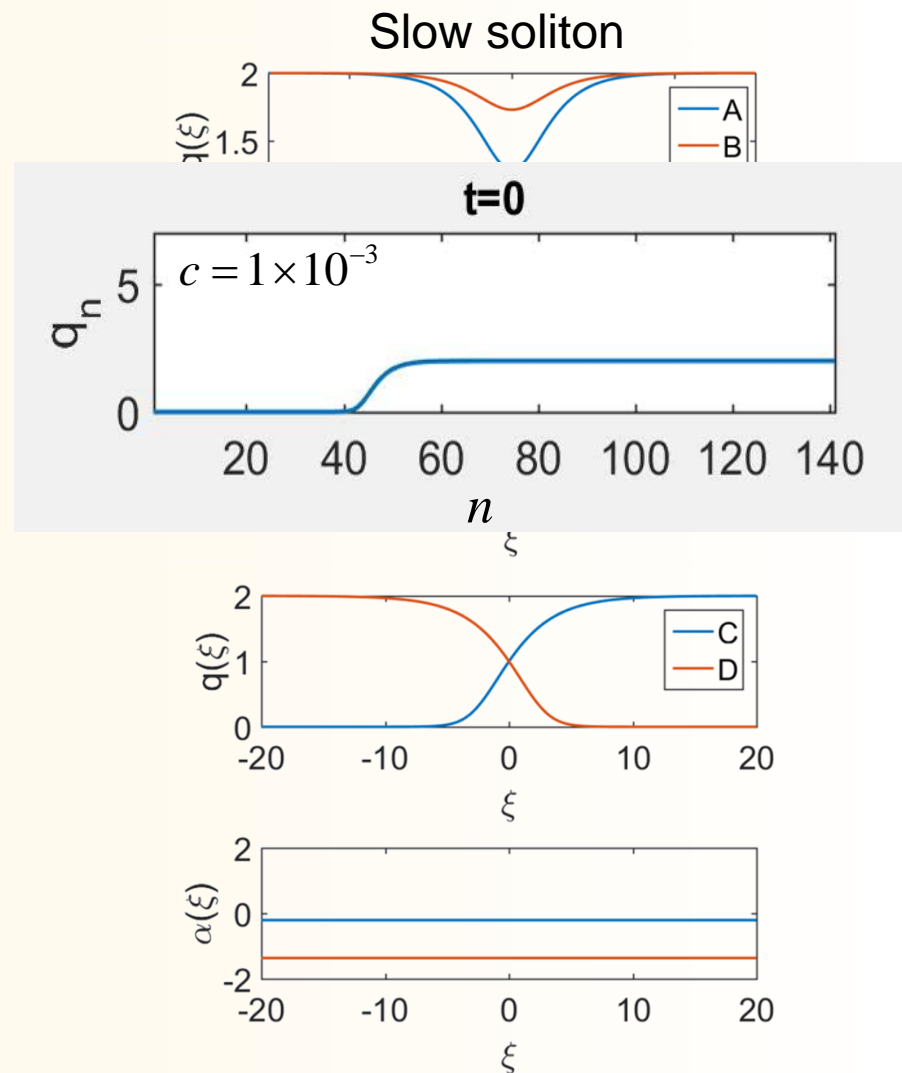
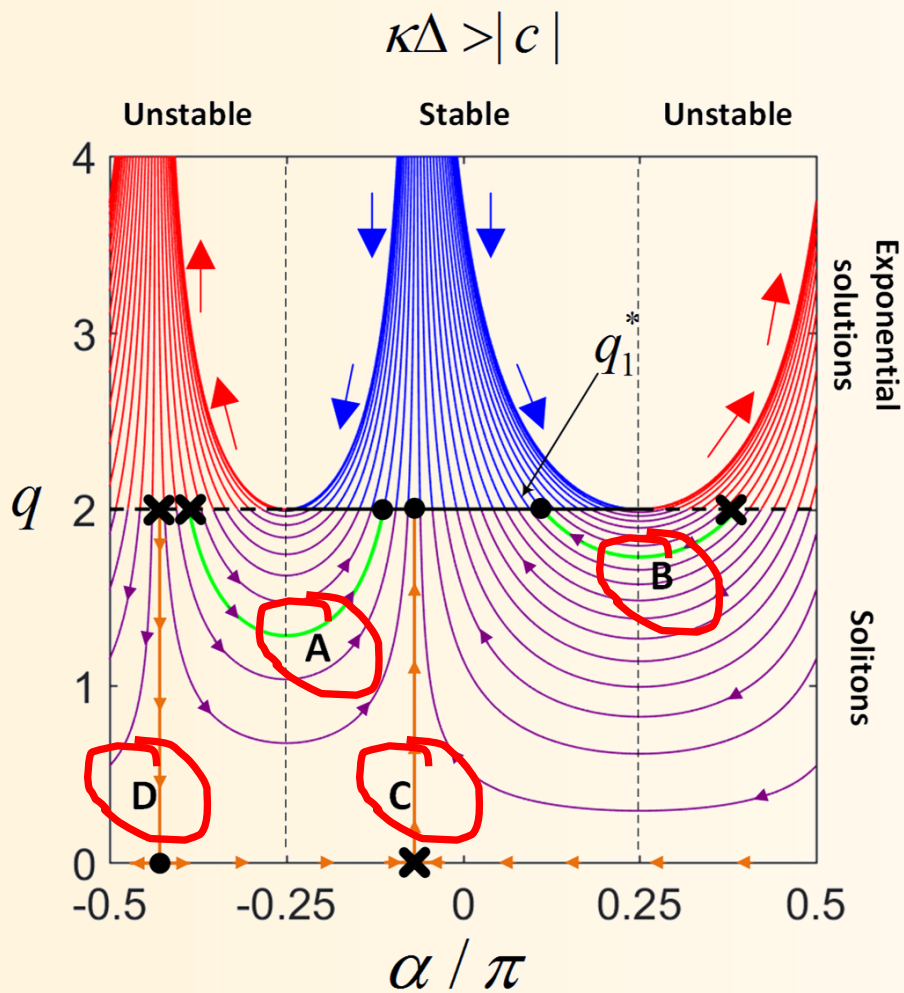
- Saddle (unstable) nodes (only if $|c| < \kappa \Delta$)

$$\underline{x}_2^* \rightarrow q_2^* = 0, \quad \alpha_2^* = -0.5 \sin^{-1} [c / \kappa \Delta],$$

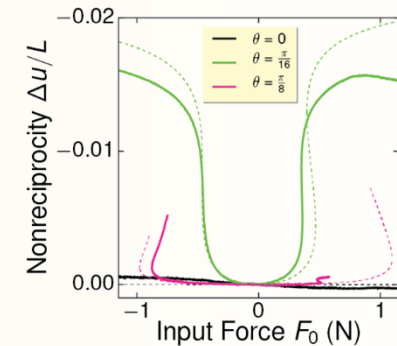
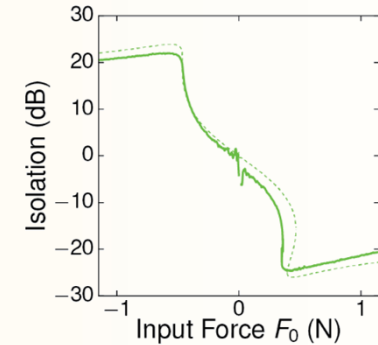
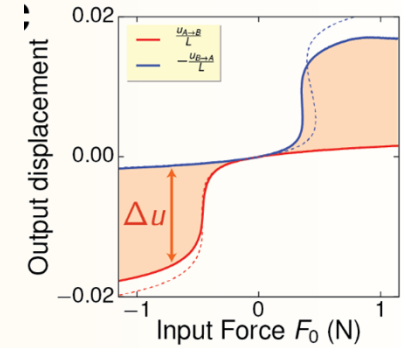
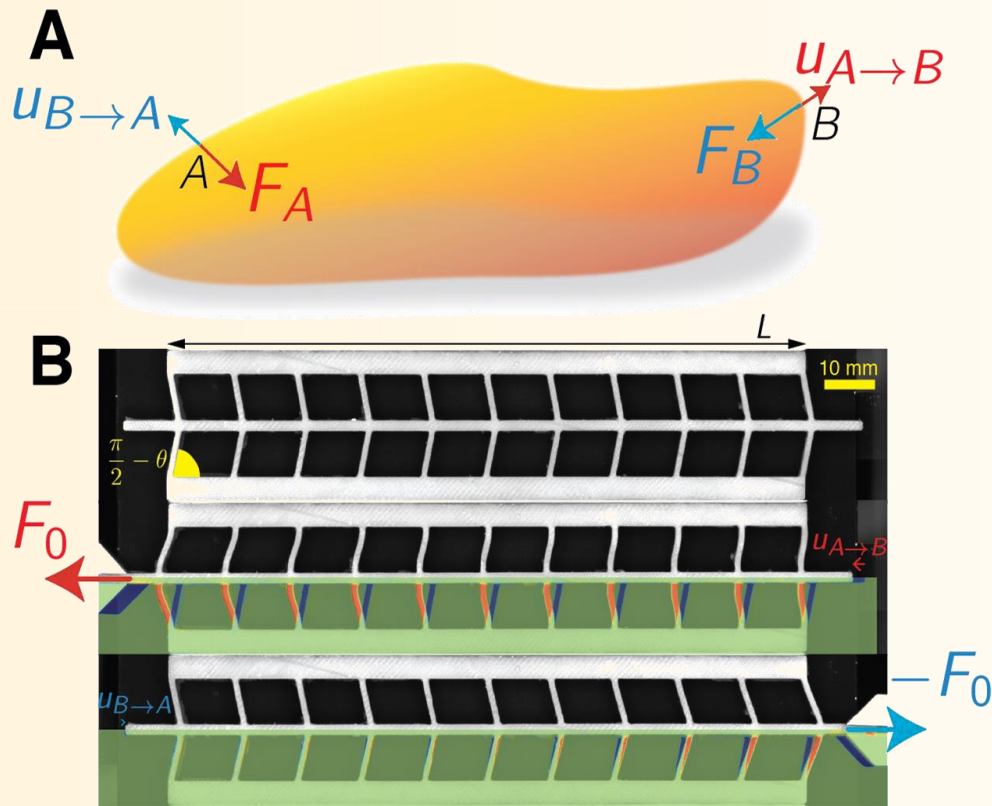
$$\underline{x}_3^* \rightarrow q_3^* = 0, \quad \alpha_3^* = -\pi / 2 - \alpha_2^*$$



PHASE DIAGRAMS AND SOLITONS

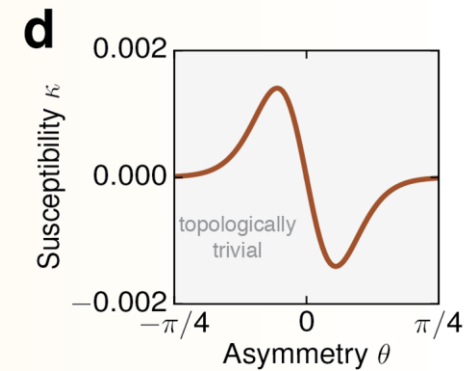
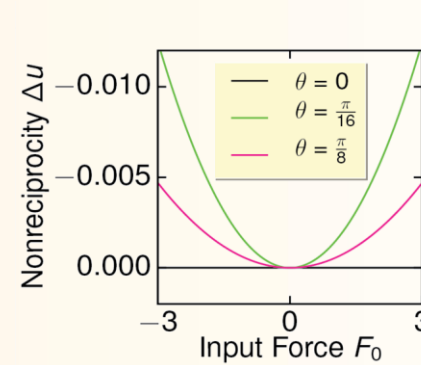
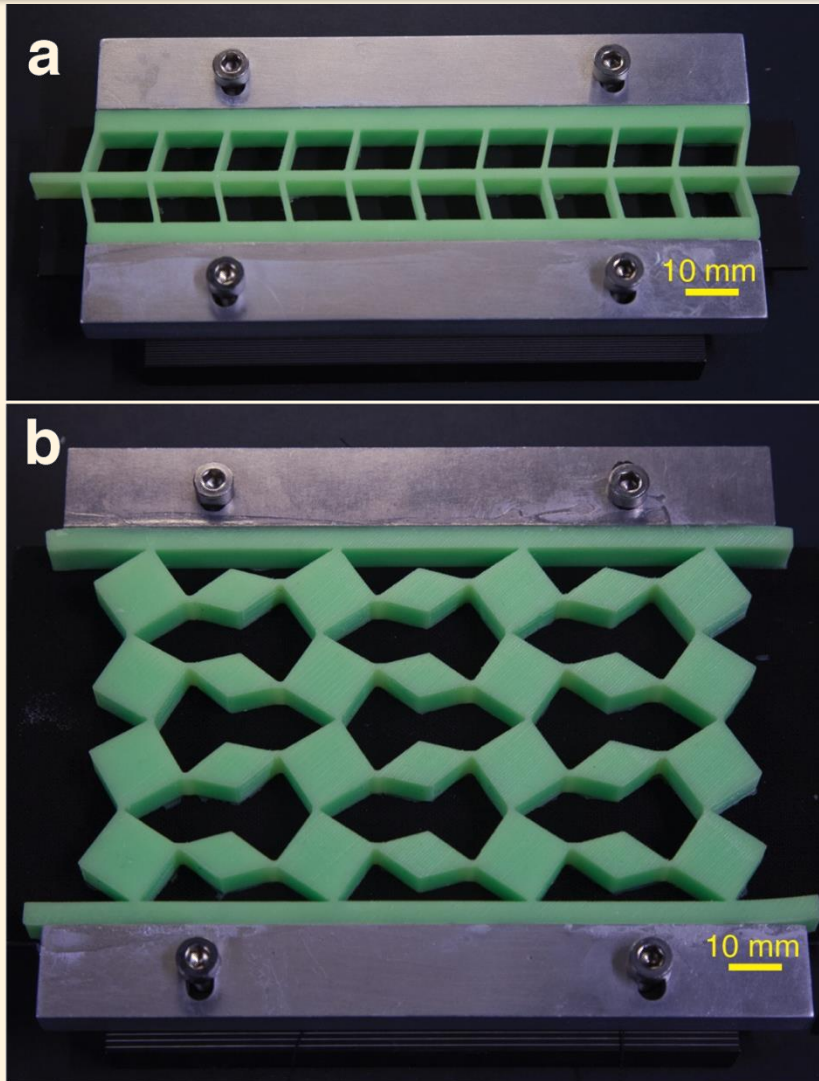


NON-RECIPROcity IN STATICS

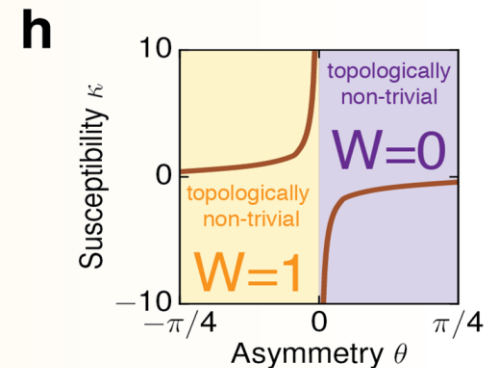
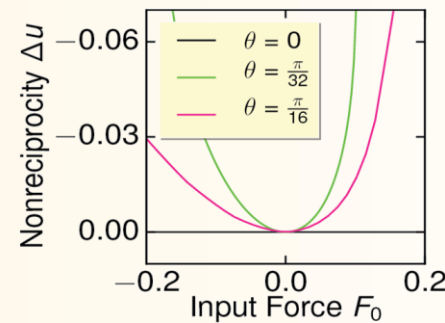


C. Coulais, D. Sounas, A. Alù, *Nature*, in press (2017)

NON-RECIPROcity IN STATIC TOPOLOGICAL METAMATERIALS

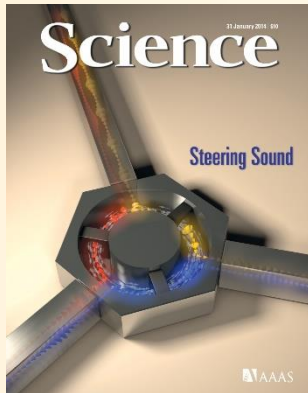


$$\Delta u = \kappa(\theta) F_0^2$$

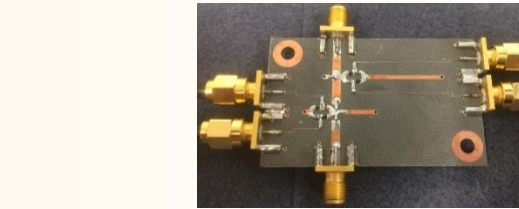


C. Coulais, D. Sounas, A. Alù, *Nature*, in press (2017)

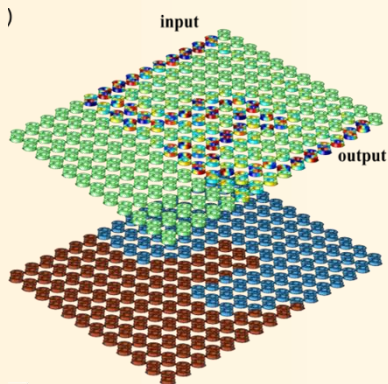
TOPOLOGICAL ELECTROMAGNETICS



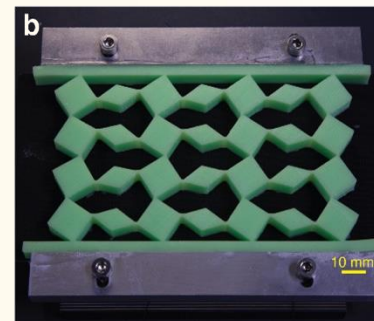
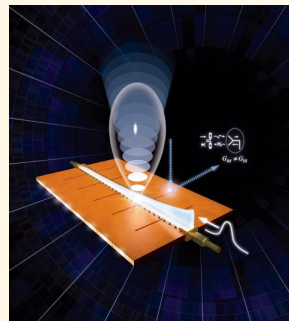
Magnetic-free, linear nonreciprocity at the subwavelength scale: angular-momentum biased meta-atoms



Nonlinearity and asymmetry to build bias-free isolators and non-reciprocal devices



Topological protection in momentum-biased arrays for reconfigurable, broadband isolation and one-way signal transport

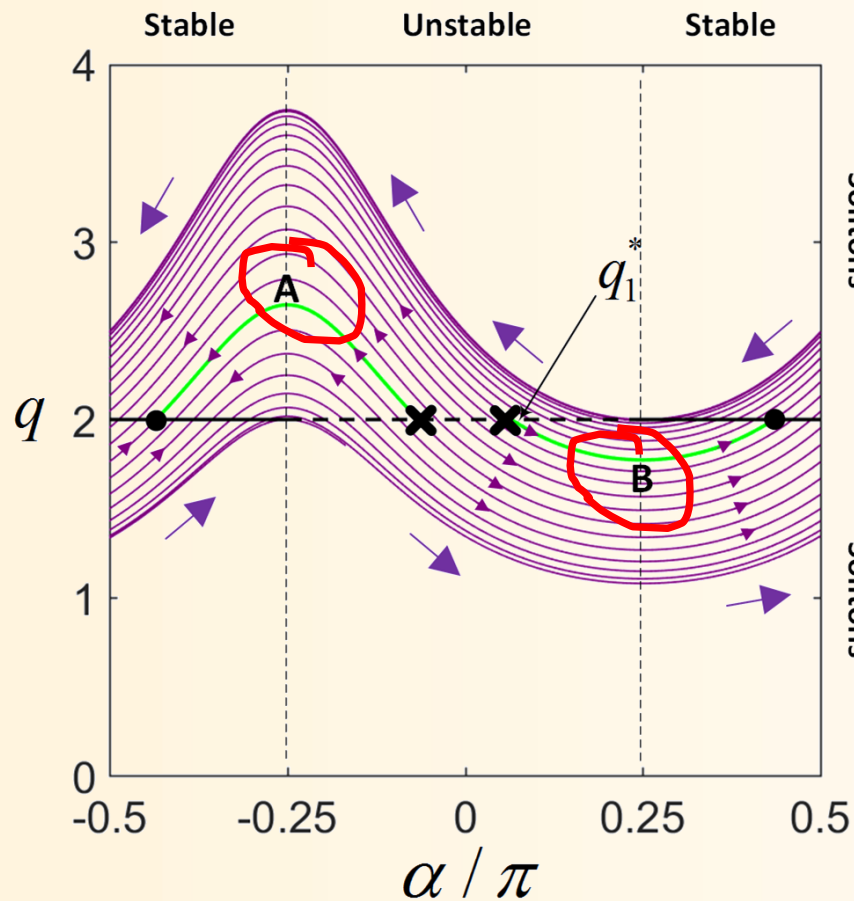


Intensity-induced topological protection in nonlinear arrays

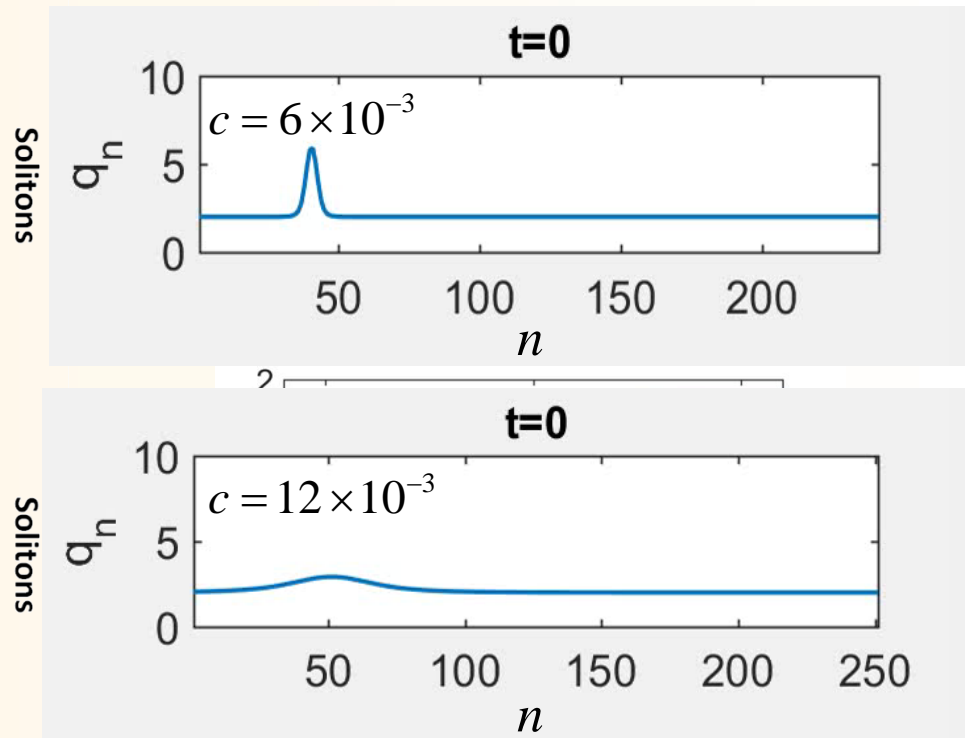


PHASE DIAGRAMS AND SOLITONS

$$\kappa\Delta < |c|$$



Fast soliton



EDGE STATES

Limit case of moving solitons with ZERO velocity

