

Turbulent and non-turbulent optical angle-of-arrival and irradiance fluctuations in the stably stratified atmosphere

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Introduction

- Motivation: why a propagation testbed?
- Equipment and experimental setup
- Retrieval of temporal dT/dz fluctuations
- Retrieval of optical **turbulence** intensity (C_n^2 , C_T^2)
- Retrieval of **transverse wind** velocities
- Summary and a conclusion

Motivation:

Why a propagation testbed?

- Need to test hypotheses and approximations (e.g., homogeneity, isotropy and stationarity of turbulence; Taylor hypothesis; Markov approximation; geometrical-optics approximation; negligibility of fluctuating aerosol concentrations)
- Need to compare optical observations/retrievals with accurate, precise, fast-response, simultaneous, in-situ observations along the propagation path.

Equipment and experimental setup

Equipment:

- Telescopes
- CCD cameras
- Ultrasonic anemometers/thermometers (“sonics”)
- Test-light arrays

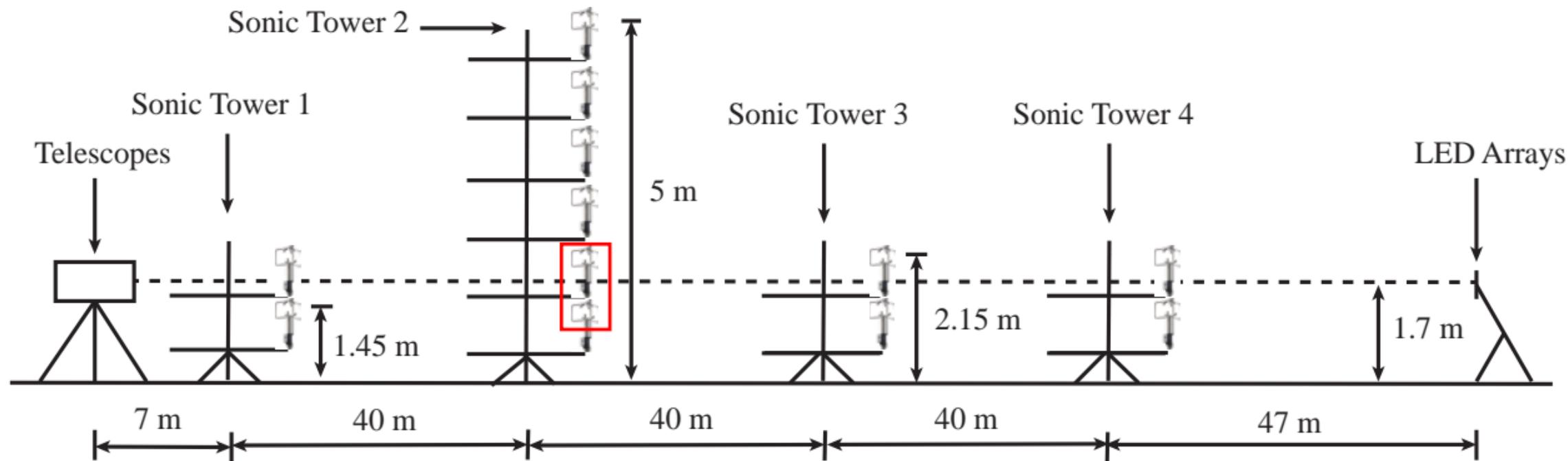
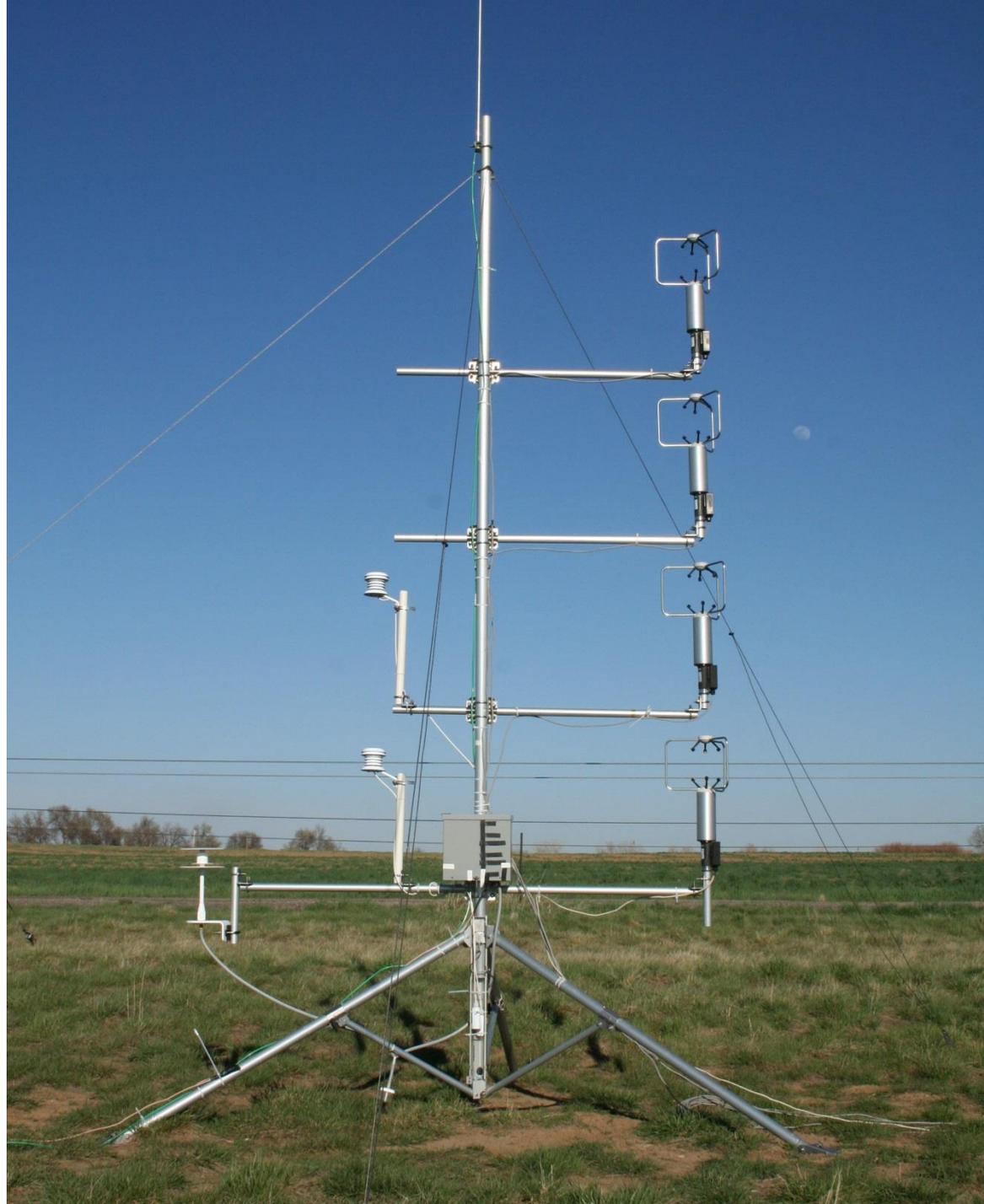


Figure 1: *Experimental setup of several optical propagation experiments conducted in August, 2012 at the Boulder Atmospheric Observatory near Erie, CO. Only data from the two sonics marked by the red box have been used for the present study.*

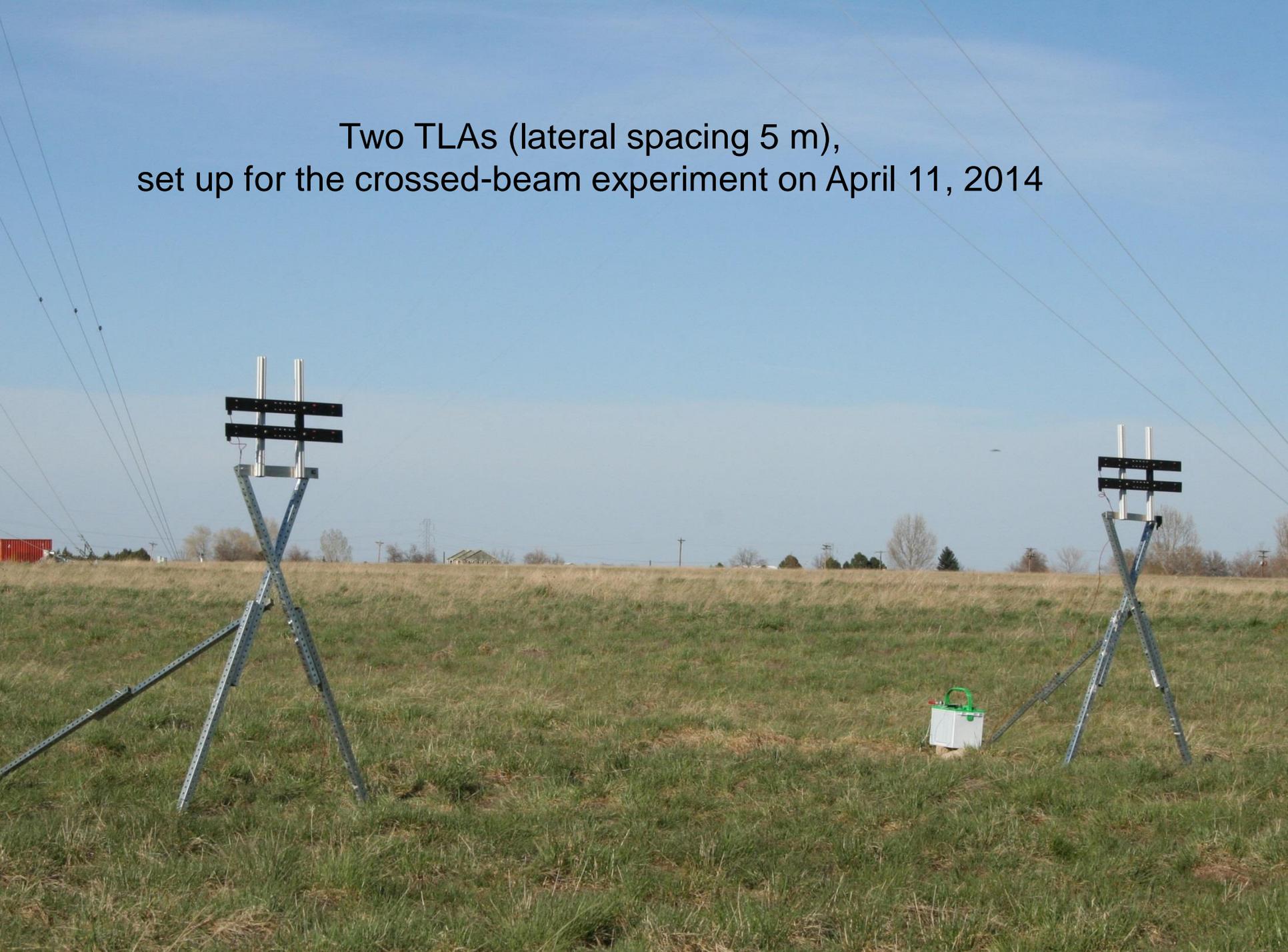
Four “sonics”
mounted on a portable tower,
crossed-beam experiment,
BAO, April 11, 2014



Test-light array (TLA),
crossed-beam experiment,
BAO, April 11, 2014



Two TLAs (lateral spacing 5 m),
set up for the crossed-beam experiment on April 11, 2014



Two telescopes (14 inch aperture),
set up for the crossed-beam experiment on April 11, 2014



Telescopes



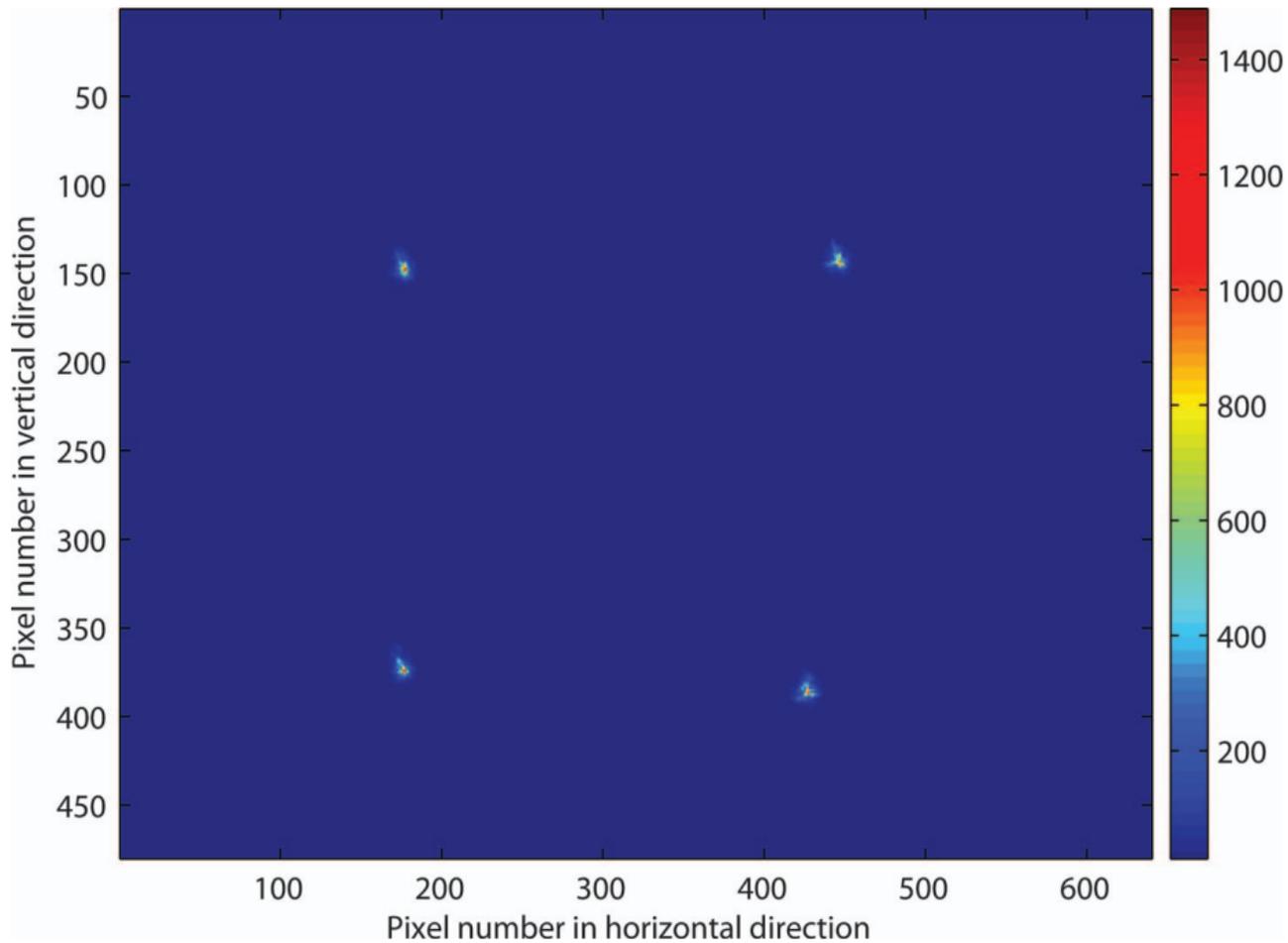


Fig. 3. Image of four lights measured at 21:00:10 LT, September 27, 2006.

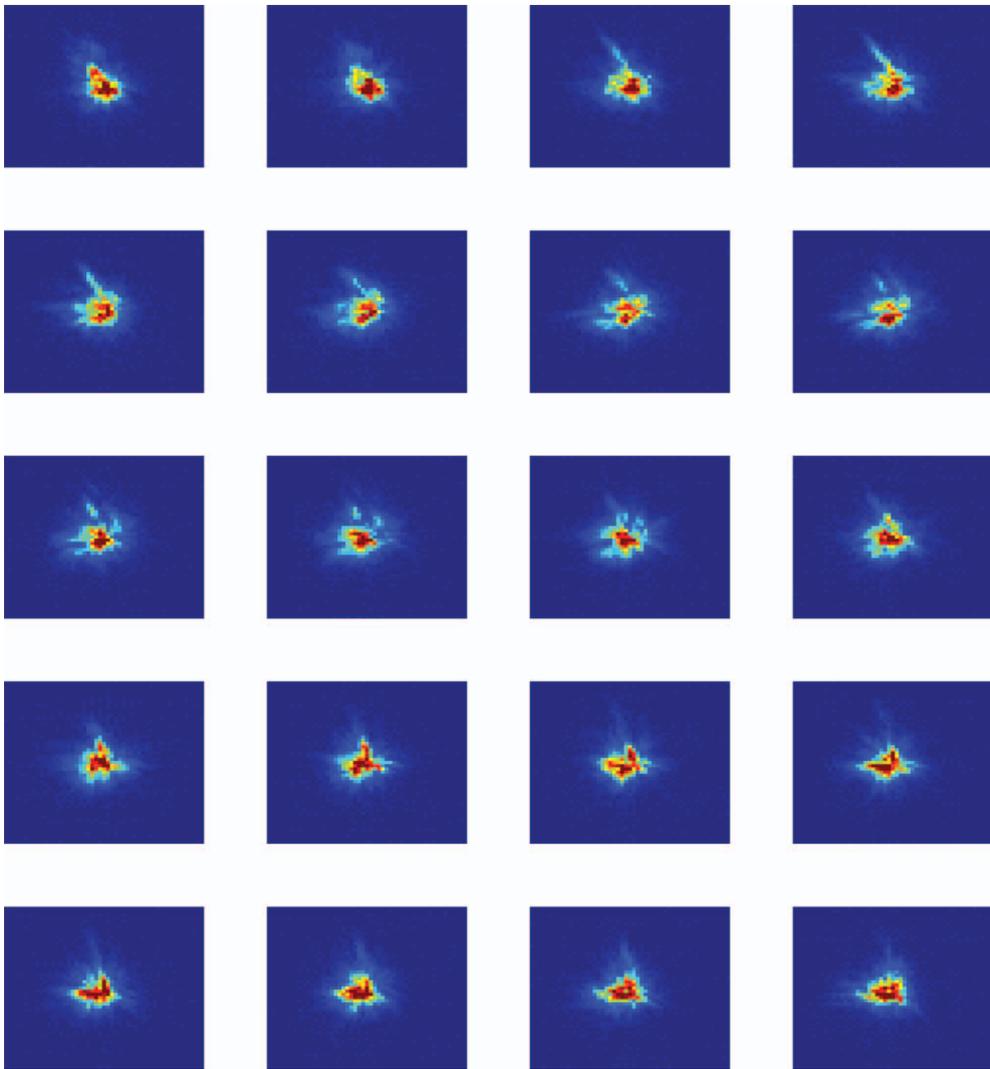


Fig. 4. Sequence of subimages (51×51 pixels) of the bottom left-hand light, measured at 21:00:10.9–21:00:11.6 LT, September 27, 2006. The first row, first column is the first image, the first row, fourth column is the fourth image.

Retrieval of temporal dT/dz fluctuations

Basic theory of refraction

Optical **refractive index** of air:

$$n = 1 + a \frac{p}{T}, \quad (1)$$

where $a = \text{constant} = 7.9 \cdot 10^{-7} \text{K Pa}^{-1}$, $p = \text{air pressure}$, $T = \text{air temperature}$.

Vertical component of the **angle-of-arrival** if source and telescope are at the same height (from eikonal equation):

$$\alpha = -\frac{\Delta\alpha}{2} = \frac{1}{2} \int_0^L \frac{\partial n(x, z)}{\partial z} dx, \quad (2)$$

where $\Delta\alpha = \text{total refraction}$, $x = \text{path coordinate}$, $z = \text{vertical coordinate}$, $L = \text{path length}$.

Hydrostatic equation and ideal **gas equation**:

$$\frac{\partial p}{\partial z} = -\rho g = -\frac{p}{R_a T} g, \quad (3)$$

where $\rho = \text{air density}$, $g = \text{acceleration due to gravity}$, $R_a = 287 \text{ J kg}^{-1} \text{ K}^{-1} = \text{gas constant for (dry) air}$.

Basic theory of refraction (cont'd)

Insert (1) and (3) into (2) to get the **AOA** in terms of the path average of the **vertical temperature gradient** $\gamma = \frac{\partial T}{\partial z}$:

$$\alpha = \frac{aL}{2} \frac{p}{T^2} \left(\frac{g}{R_a} + \gamma \right). \quad (4)$$

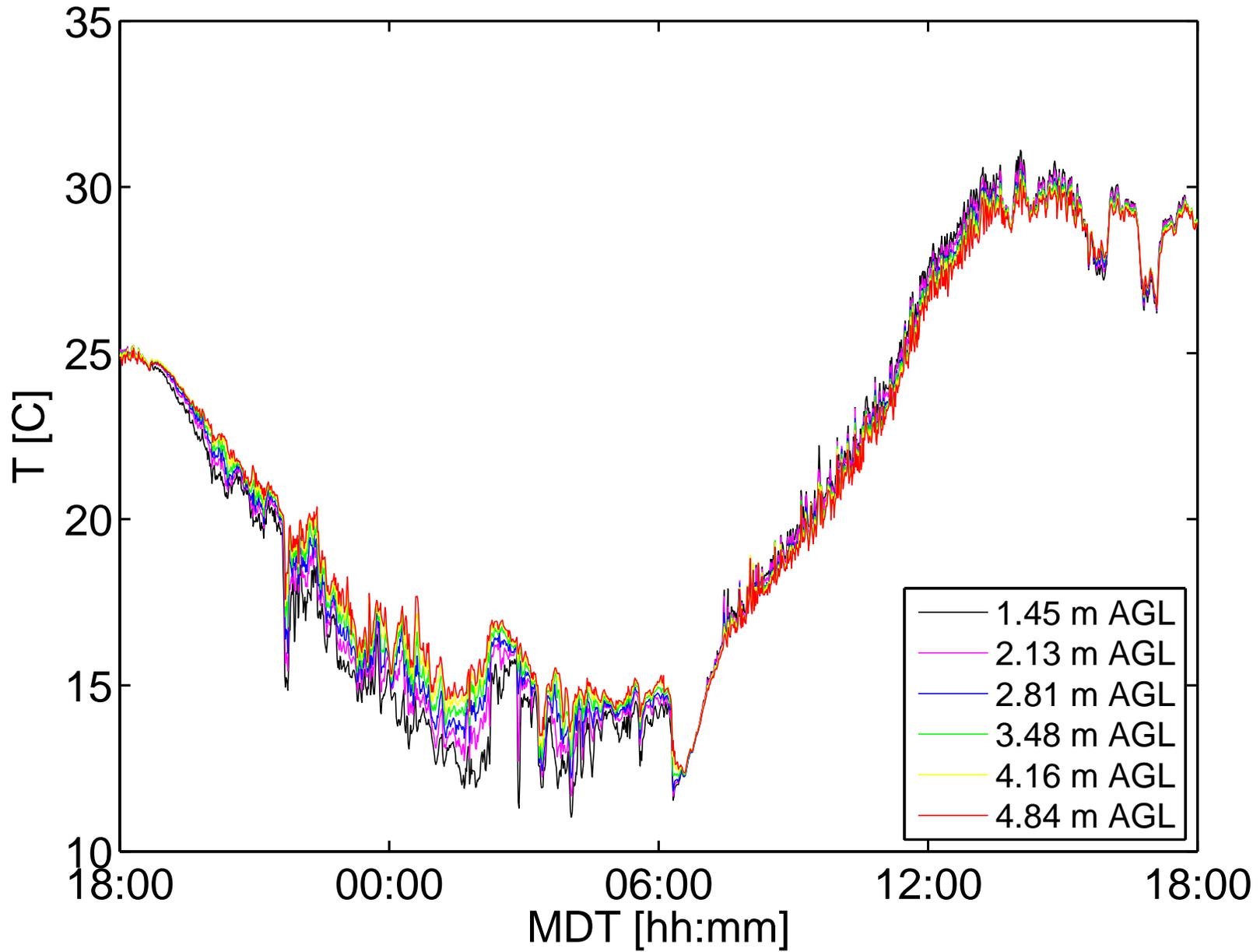
Sensitivity of AOA with respect to γ :

$$s = \frac{\partial \alpha}{\partial \gamma} = \frac{aL}{2} \frac{p}{T^2}. \quad (5)$$

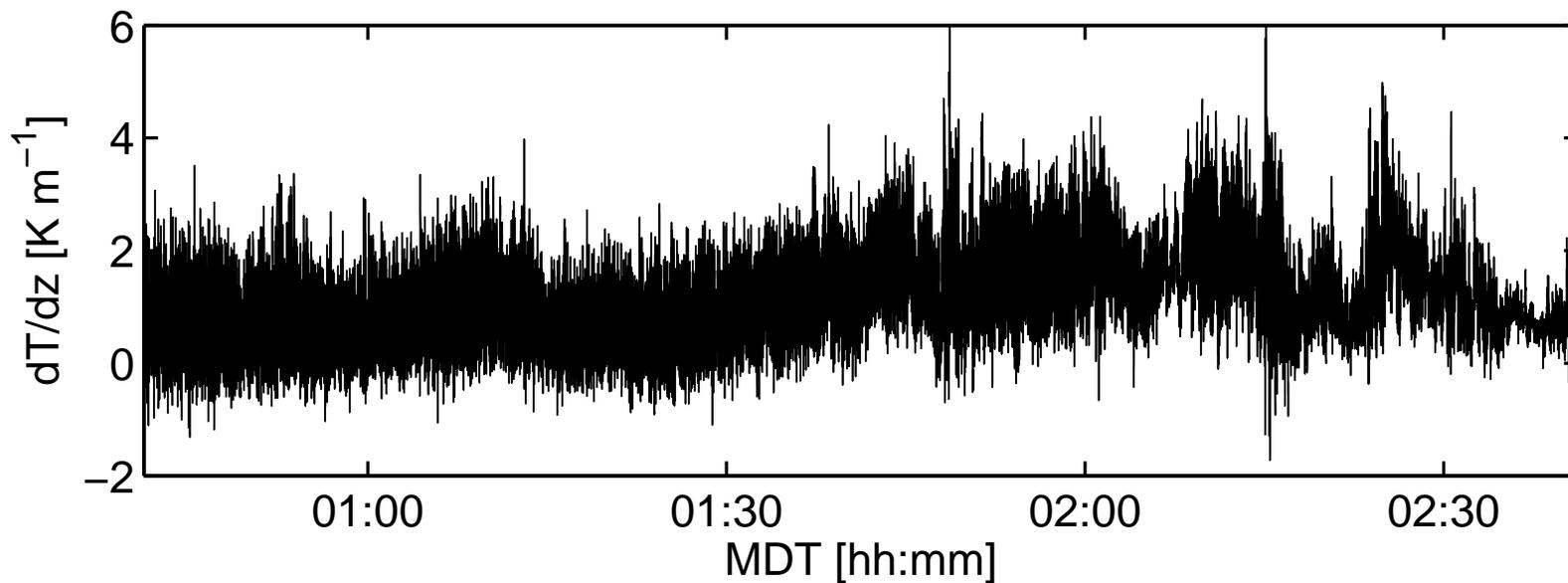
In our case (night of 14/15 August, 2012): $p = 839$ hPa, $T = 288$ K and $L = 174$ m, such that

$$s = 73 \frac{\mu\text{rad}}{\text{K m}^{-1}} = \frac{1 \mu\text{rad}}{0.014 \frac{\text{K}}{\text{m}}}. \quad (6)$$

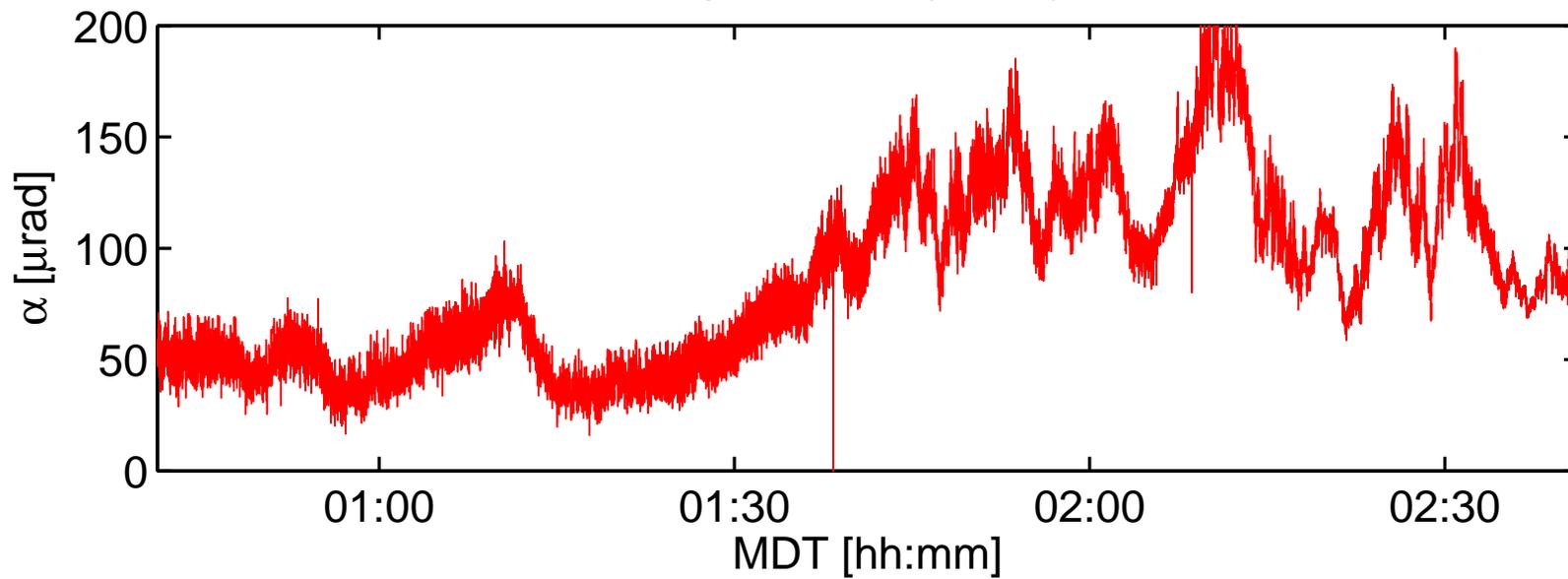
Temperature



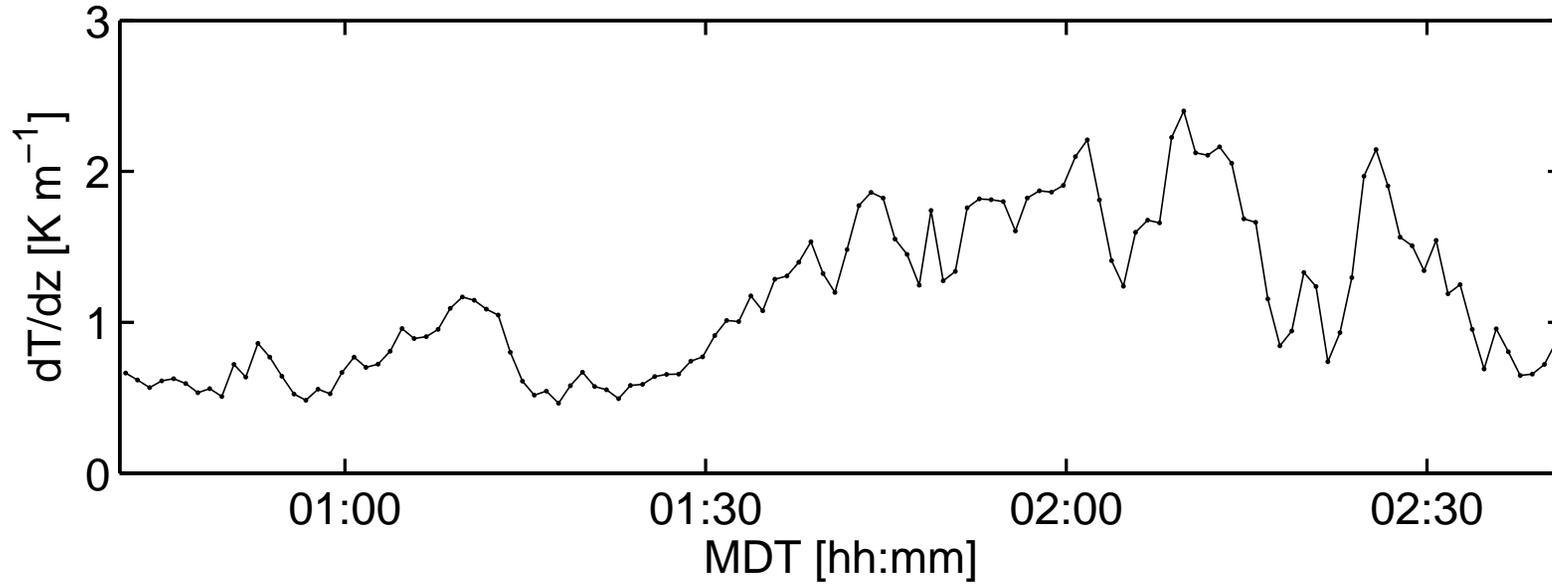
Sonic data (20 Hz)



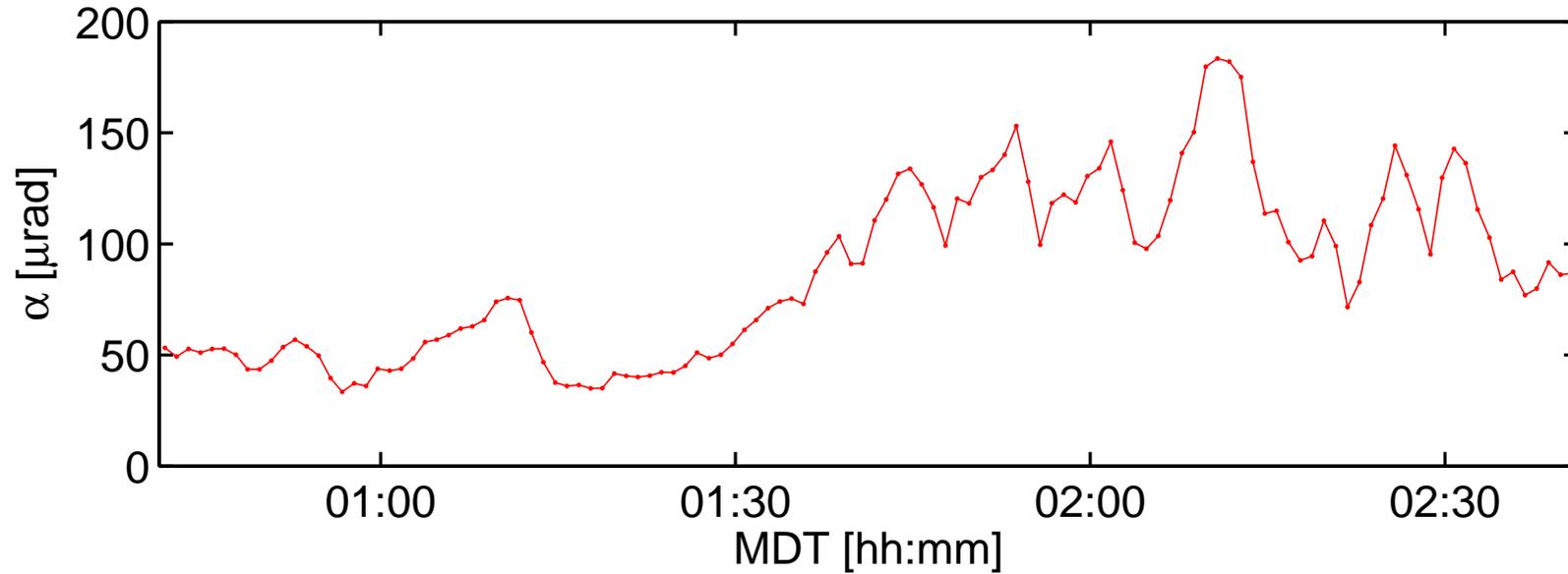
Optical data (20 Hz)



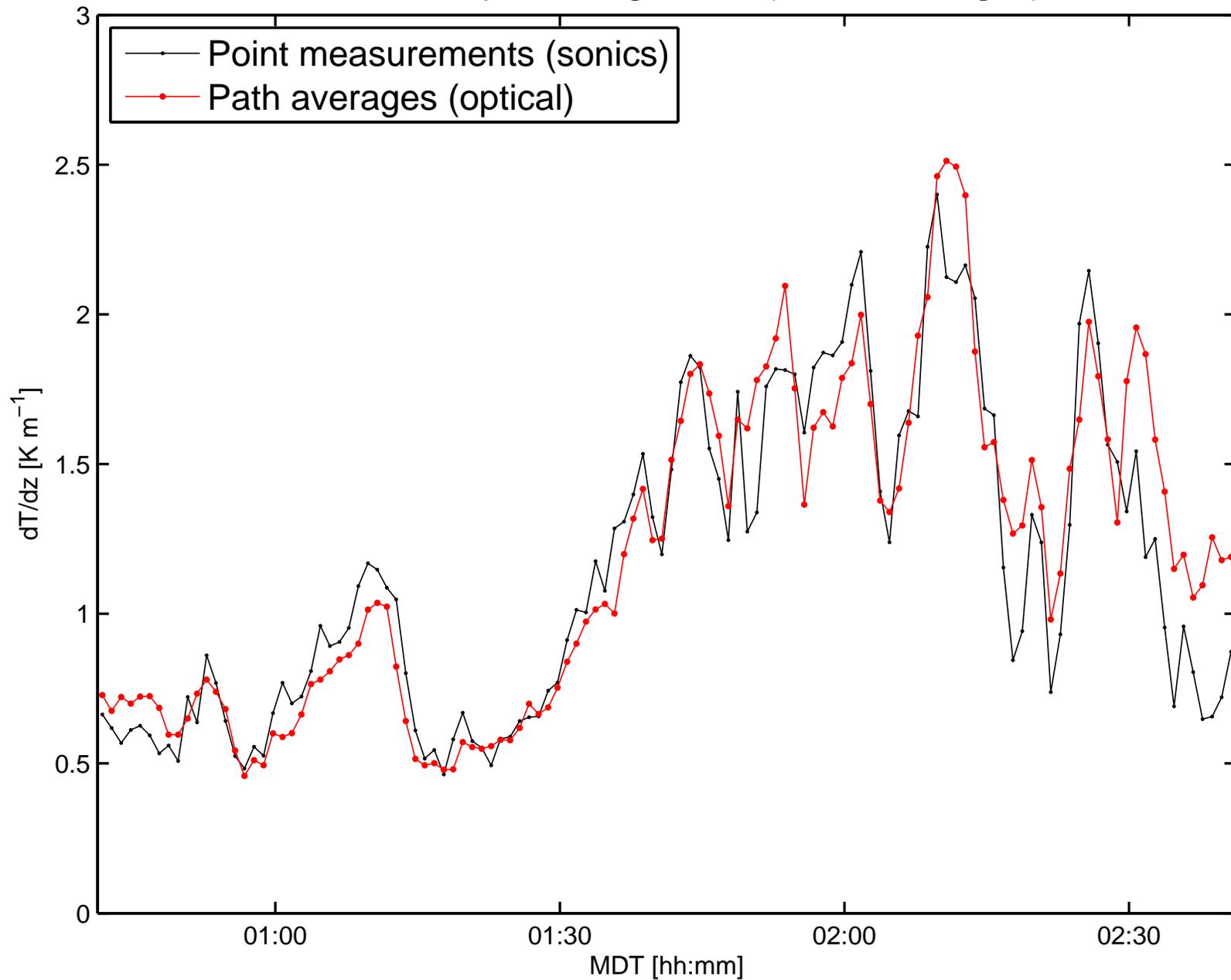
Sonic data (1-min averages)



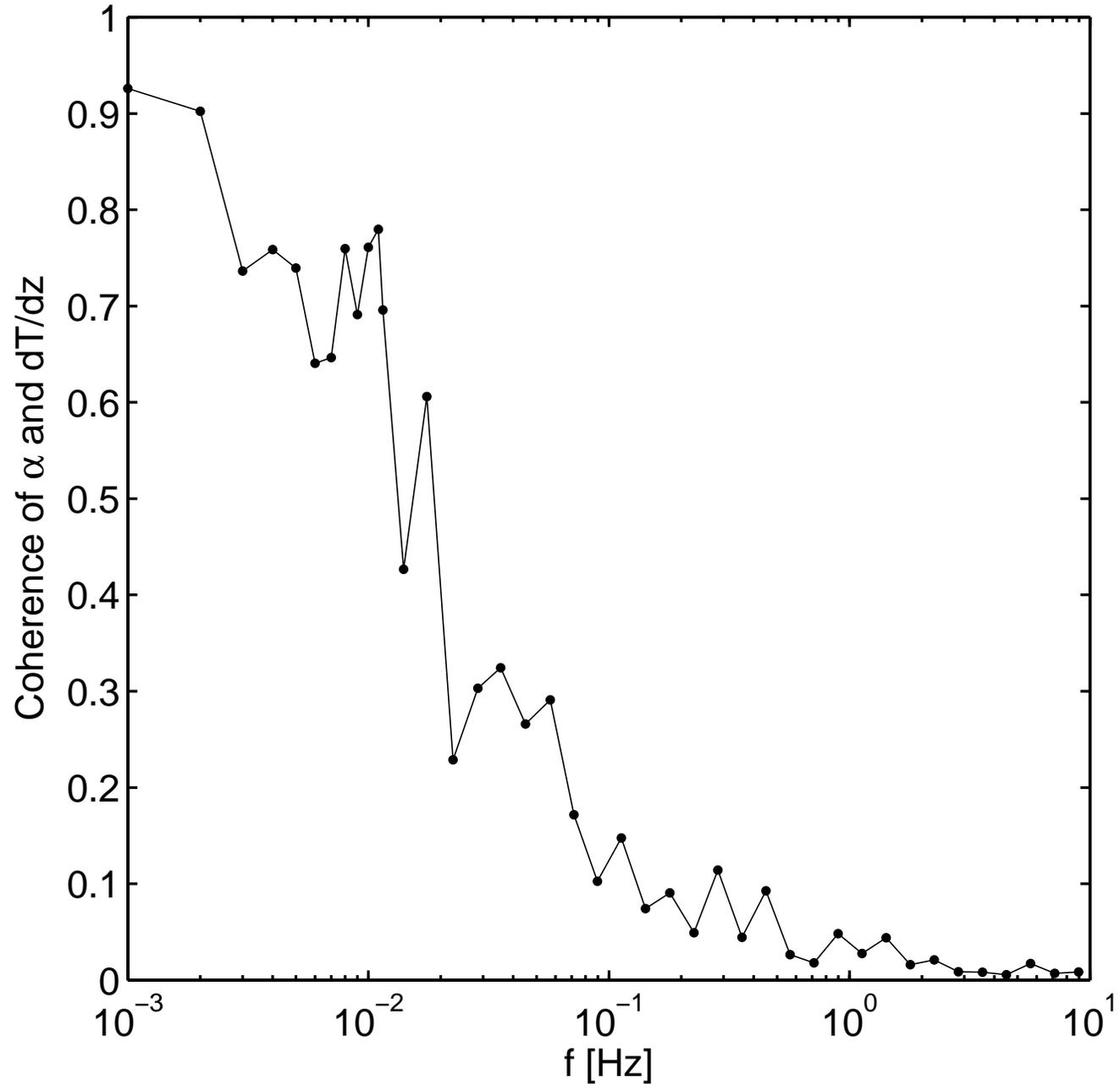
Optical data (1-min averages)



Vertical temperature gradient (1-min averages)



Coherence spectrum of α and dT/dz



Retrieval of optical turbulence intensity
 (C_n^2, C_T^2)

Theory: C_n^2 and C_T^2

Structure function of a turbulent refractive-index field:

$$D_n(\mathbf{x}, \mathbf{r}) = \langle [n(\mathbf{x} + \mathbf{r}) - n(\mathbf{x})]^2 \rangle, \quad (1)$$

where \mathbf{x} = location and \mathbf{r} = spatial separation

Homogeneous turbulence: $D_n(\mathbf{x}, \mathbf{r})$ is independent of \mathbf{x} .

Isotropic turbulence: $D_n(\mathbf{x}, \mathbf{r})$ is independent of \mathbf{r} .

Homogeneous and isotropic turbulence: $D_n(\mathbf{x}, \mathbf{r})$ depends only on $r = |\mathbf{r}|$.

Homogeneous and isotropic turbulence in the **inertial subrange** (Obukhov 1949):

$$D_n(r) = C_n^2 r^{2/3}, \quad (2)$$

where C_n^2 is the refractive-index structure parameter. For dry air, where

$$n = 1 + a \frac{p}{T}, \quad (3)$$

we have

$$D_T(\mathbf{x}, \mathbf{r}) = \langle [T(\mathbf{x} + \mathbf{r}) - T(\mathbf{x})]^2 \rangle = a^2 \frac{p^2}{T^4} \langle [n(\mathbf{x} + \mathbf{r}) - n(\mathbf{x})]^2 \rangle, \quad (4)$$

such that for inertial-range turbulence the temperature structure parameter is

$$C_T^2 = a^2 \frac{p^2}{T^4} C_n^2. \quad (5)$$

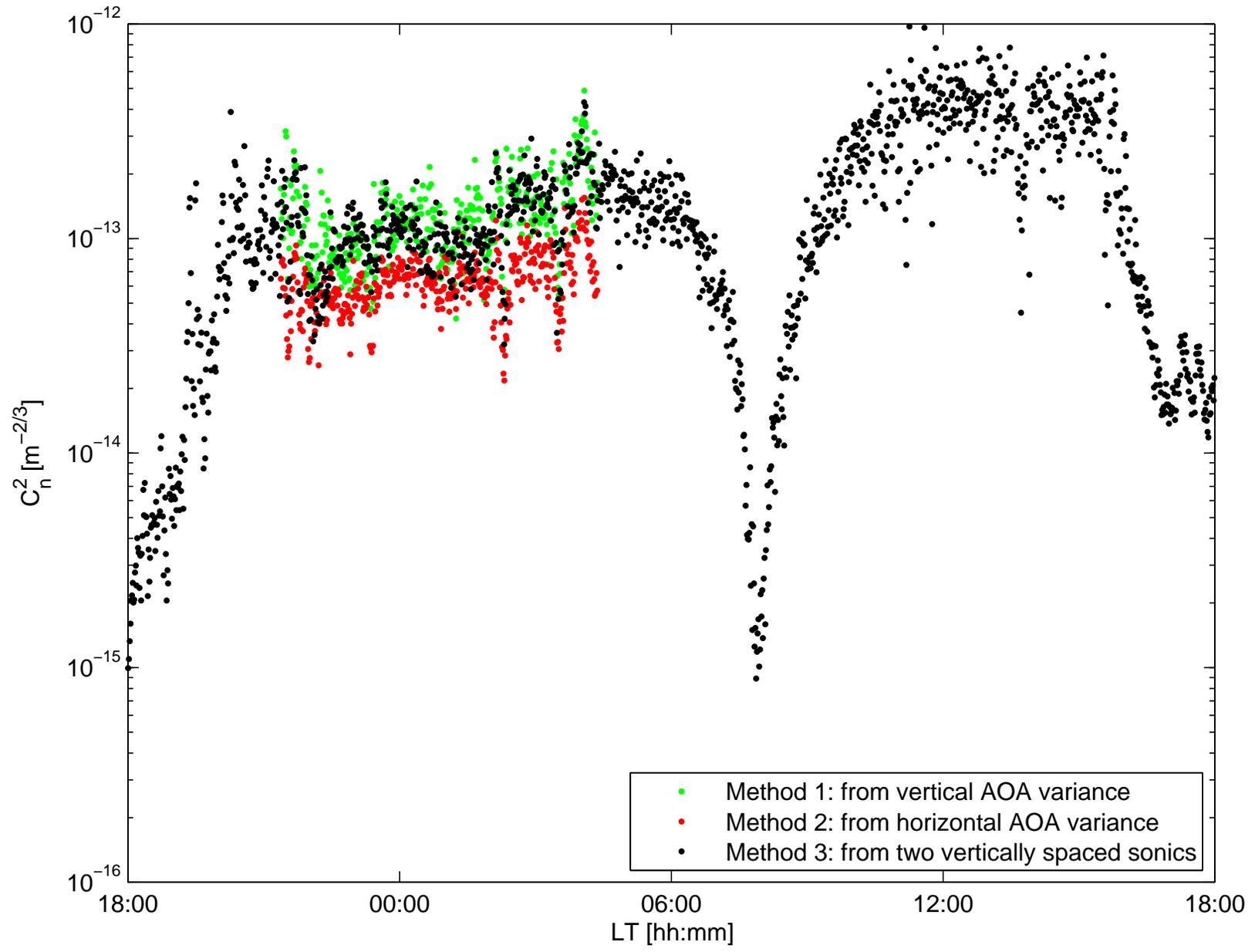
For **optical turbulence** in the inertial subrange, the variance of the vertical AOA fluctuations, σ_α^2 , is equal to the variance of the horizontal AOA fluctuations, σ_β^2 :

$$\sigma_\alpha^2 = \sigma_\beta^2. \quad (6)$$

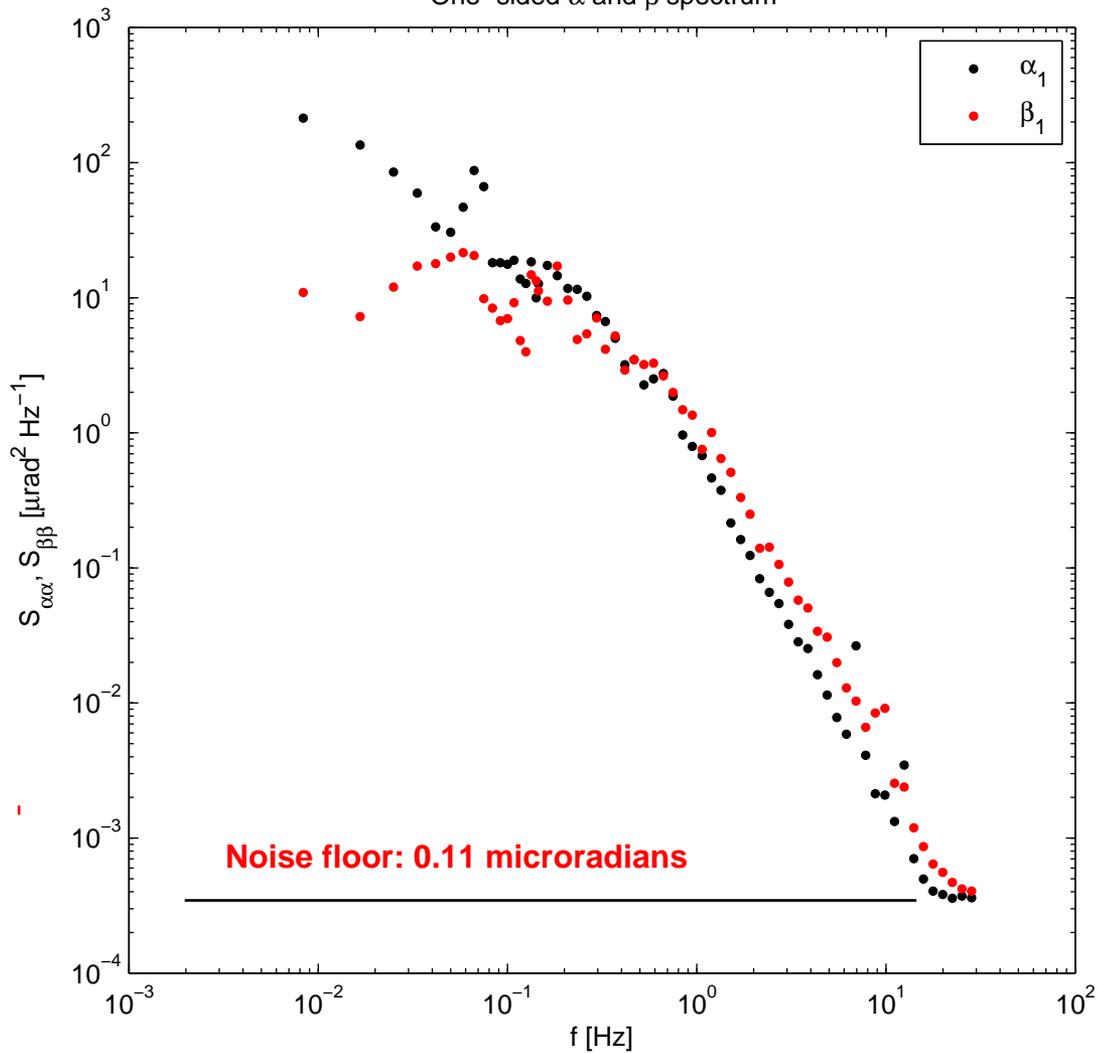
By means of geometrical optics, we obtain for a spherical wave propagating through inertial-range turbulence and received with a circular aperture:

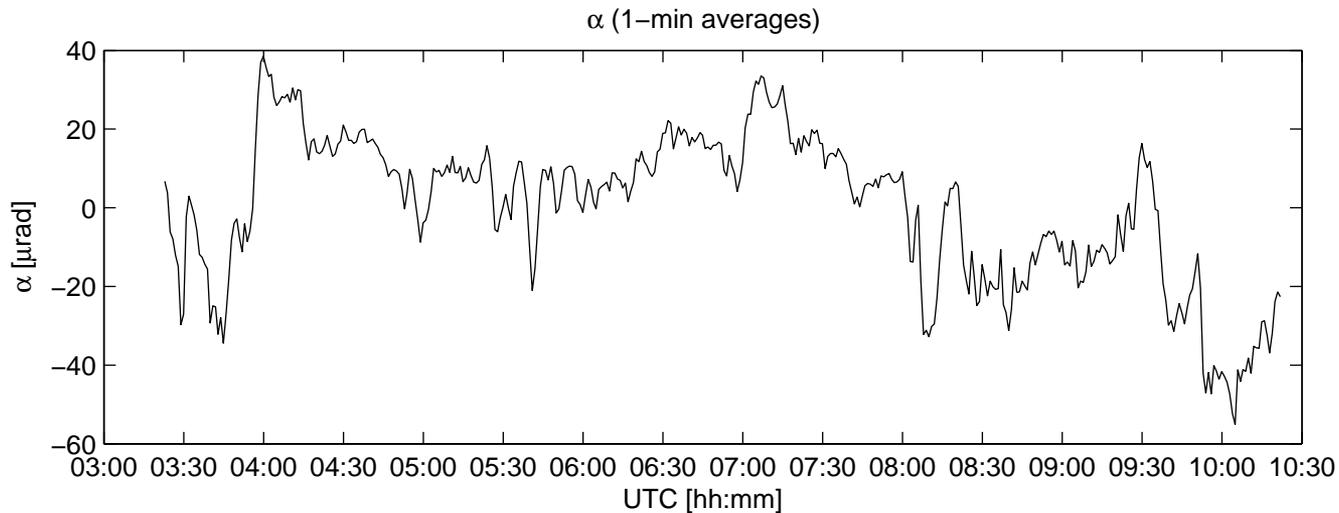
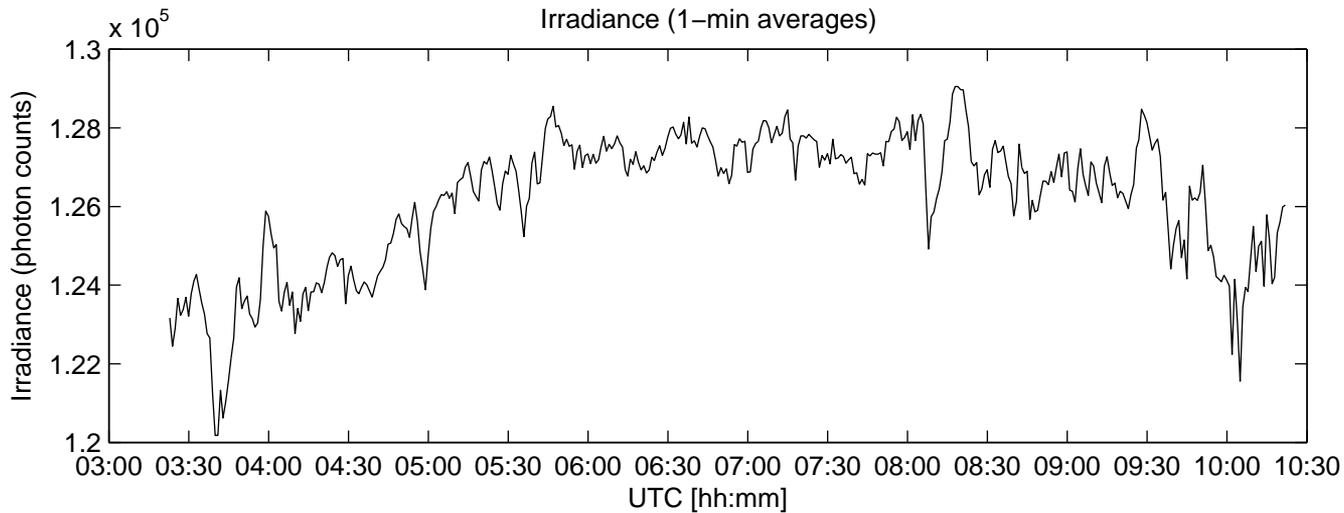
$$\sigma_\alpha^2 = \sigma_\beta^2 = 1.064 LD^{-1/3} C_n^2, \quad (7)$$

where L = propagation path length and D = telescope's aperture diameter.

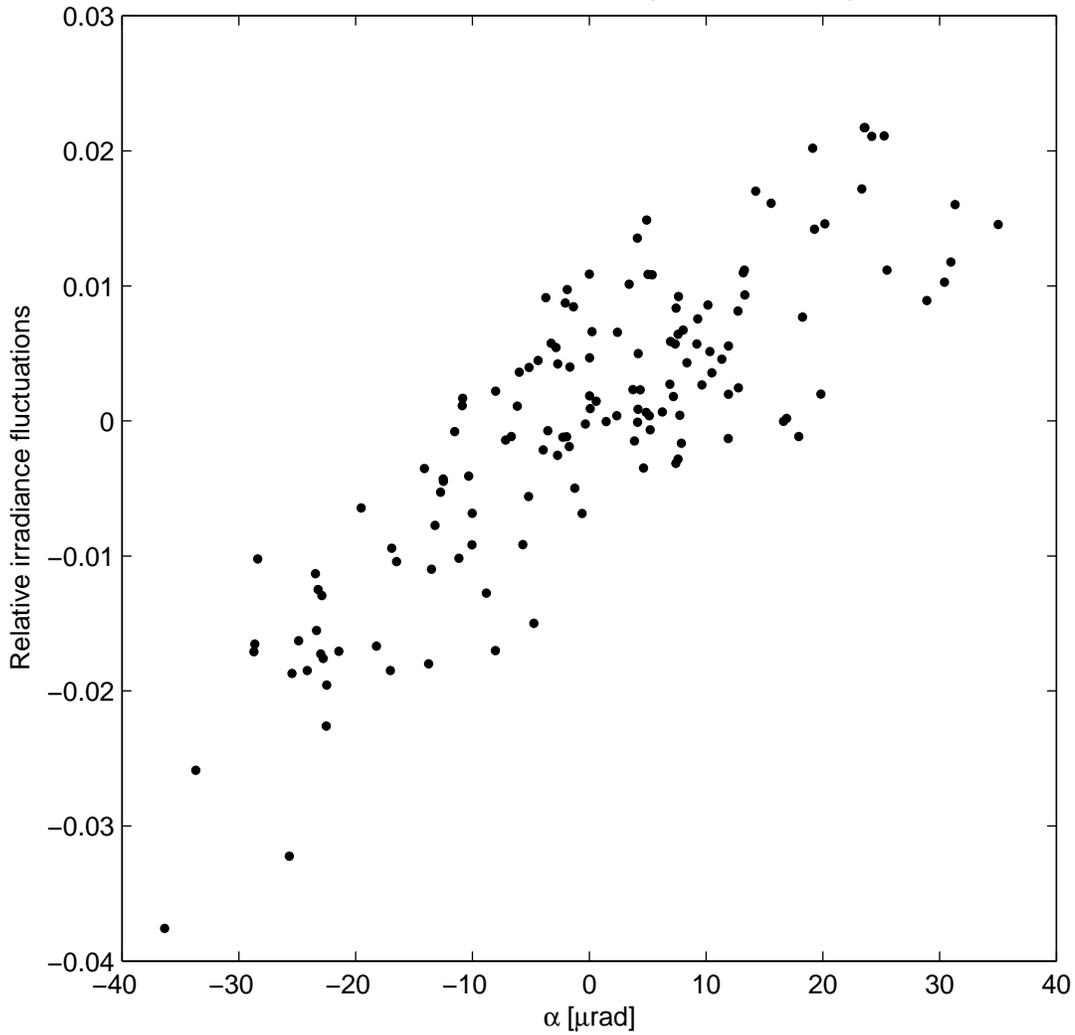


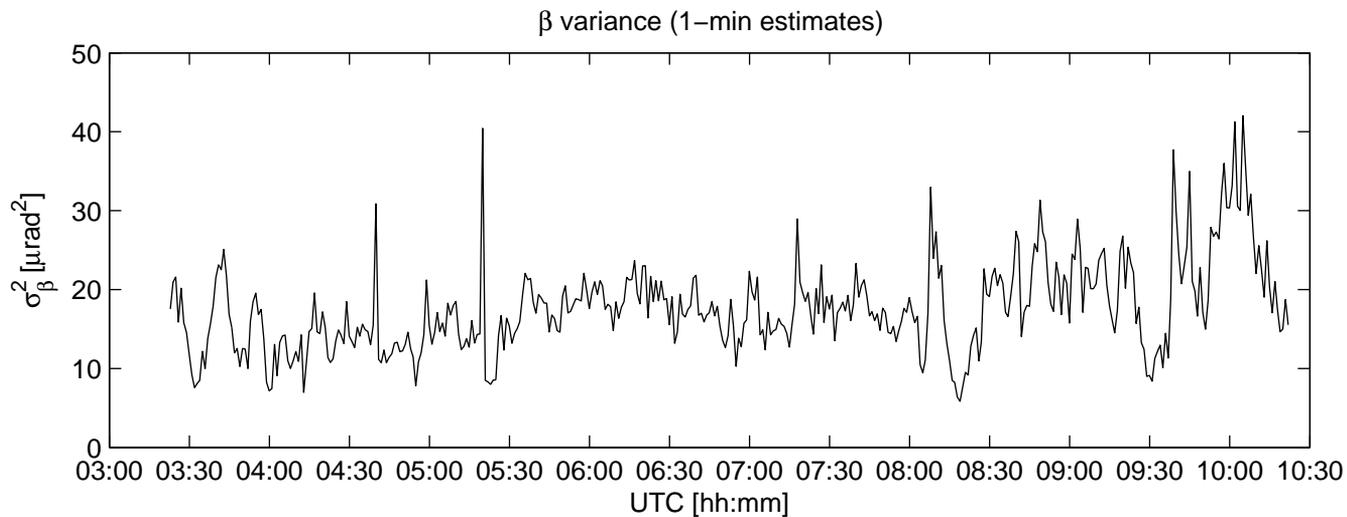
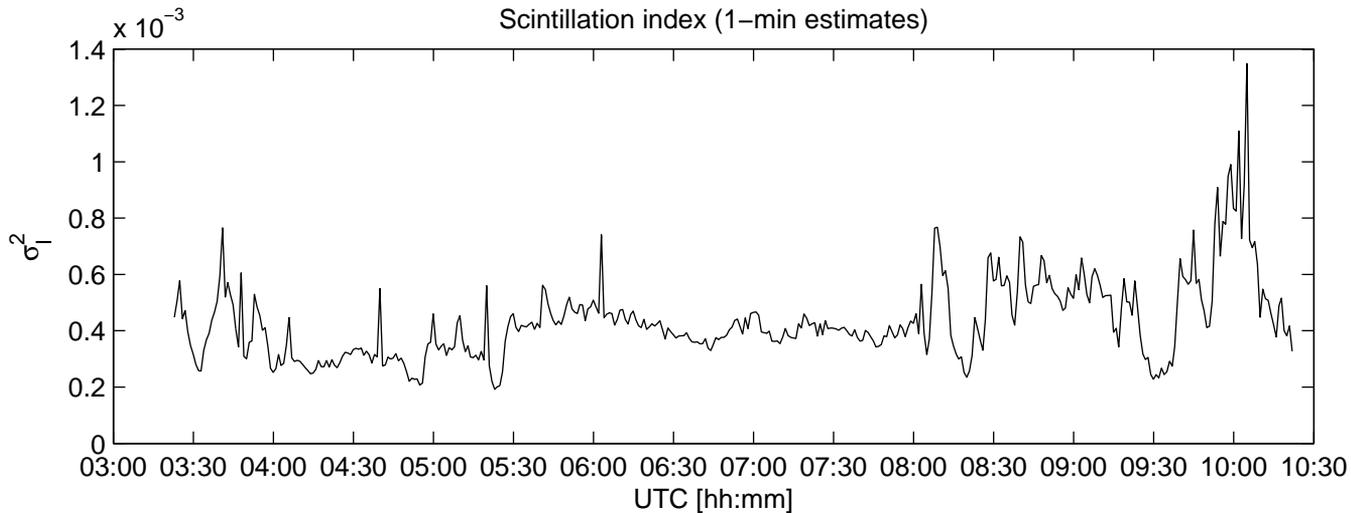
One-sided α and β spectrum

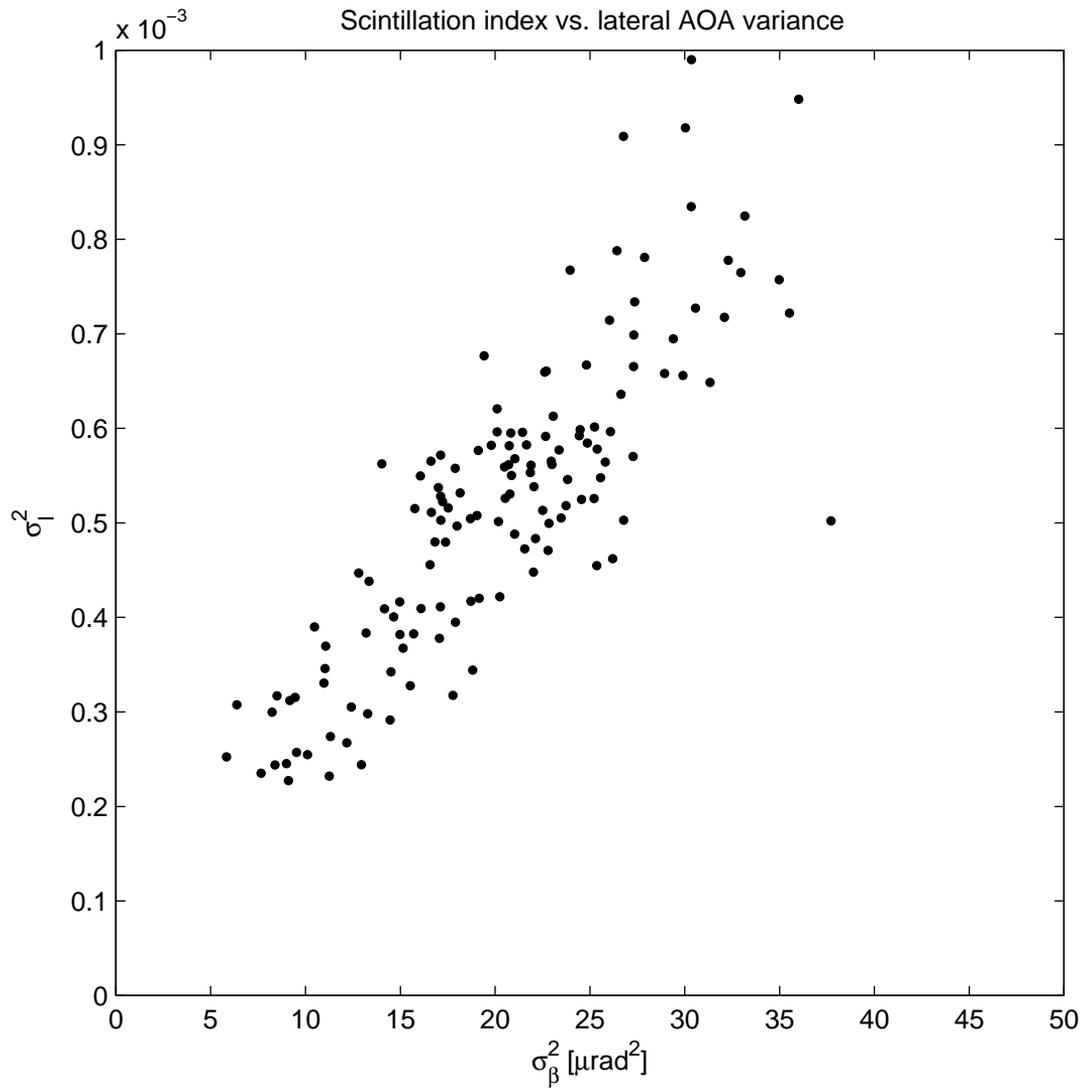




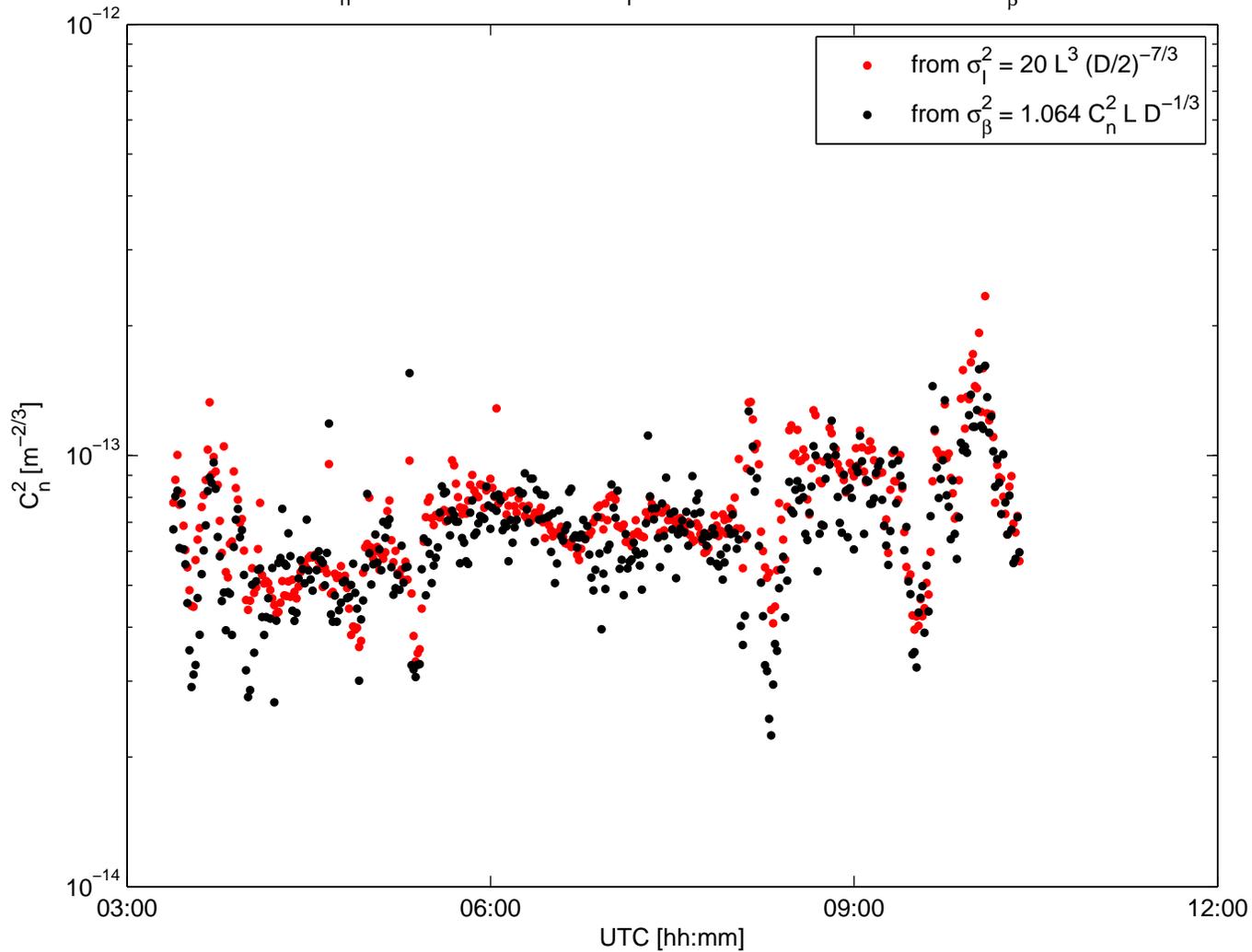
Irradiance vs. vertical AOA (1-min estimates)







C_n^2 from scintillation index (σ_1^2) and from lateral AOA variance (σ_β^2)



Retrieval of transverse wind velocities

Basic theory of the AOA spectrum

Assumptions:

- Homogeneous and isotropic inertial-range turbulence in the temperature field
- Taylor's frozen-turbulence hypothesis valid, constant baseline wind speed v_b
- Spherical wave emitted from point source
- Two-point interferometer with baseline length b
- Geometrical optics (frequency f much lower than Fresnel frequency $v_b/\sqrt{L\lambda}$)

Clifford (*J. Opt. Am. Soc. A*, 1971)

and Cheon, Hohreiter, Behn, and Muschinski (*J. Opt. Am. Soc. A*, 2007) predict:

$$S_\beta(f) = \frac{2^{4/3}}{9\pi^{7/6}\Gamma(5/6)} C_n^2 v_b^{5/3} L b^{-2} f^{-8/3} \left[1 - \frac{\sin(2\pi b f / v_b)}{2\pi b f / v_b} \right], \quad (1)$$

where C_n^2 = refractive-index structure parameter,

L = path length,

f = frequency,

β = horizontal AOA (fluctuation).

Two special cases

High frequencies, $f \gg v_b/b$ (aperture filtering):

$$S_\beta(f) = \frac{2^{4/3}}{9\pi^{7/6}\Gamma(5/6)} C_n^2 v_b^{5/3} L b^{-2} f^{-8/3}. \quad (2)$$

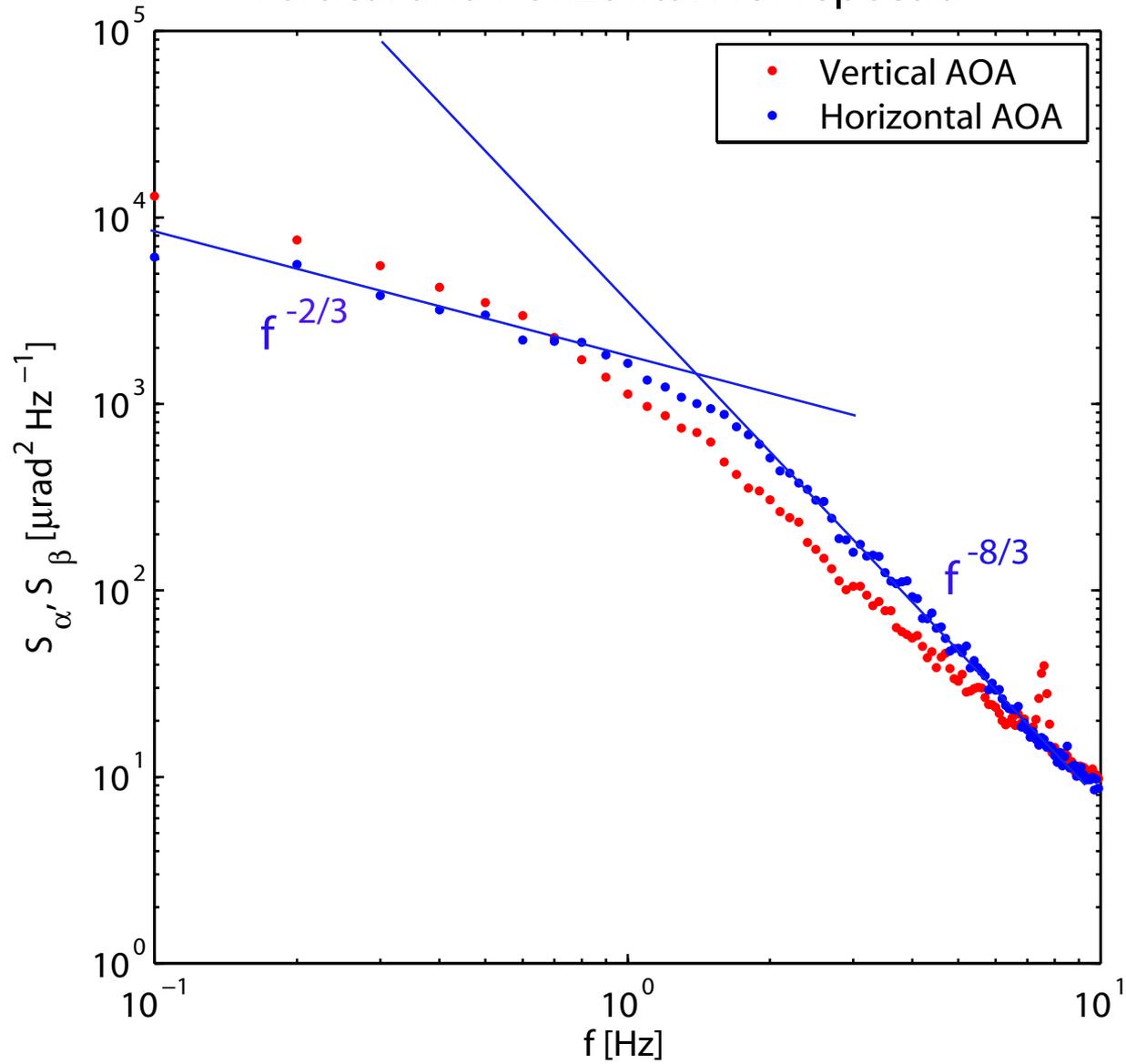
Low frequencies, $f \ll v_b/b$ (no aperture filtering):

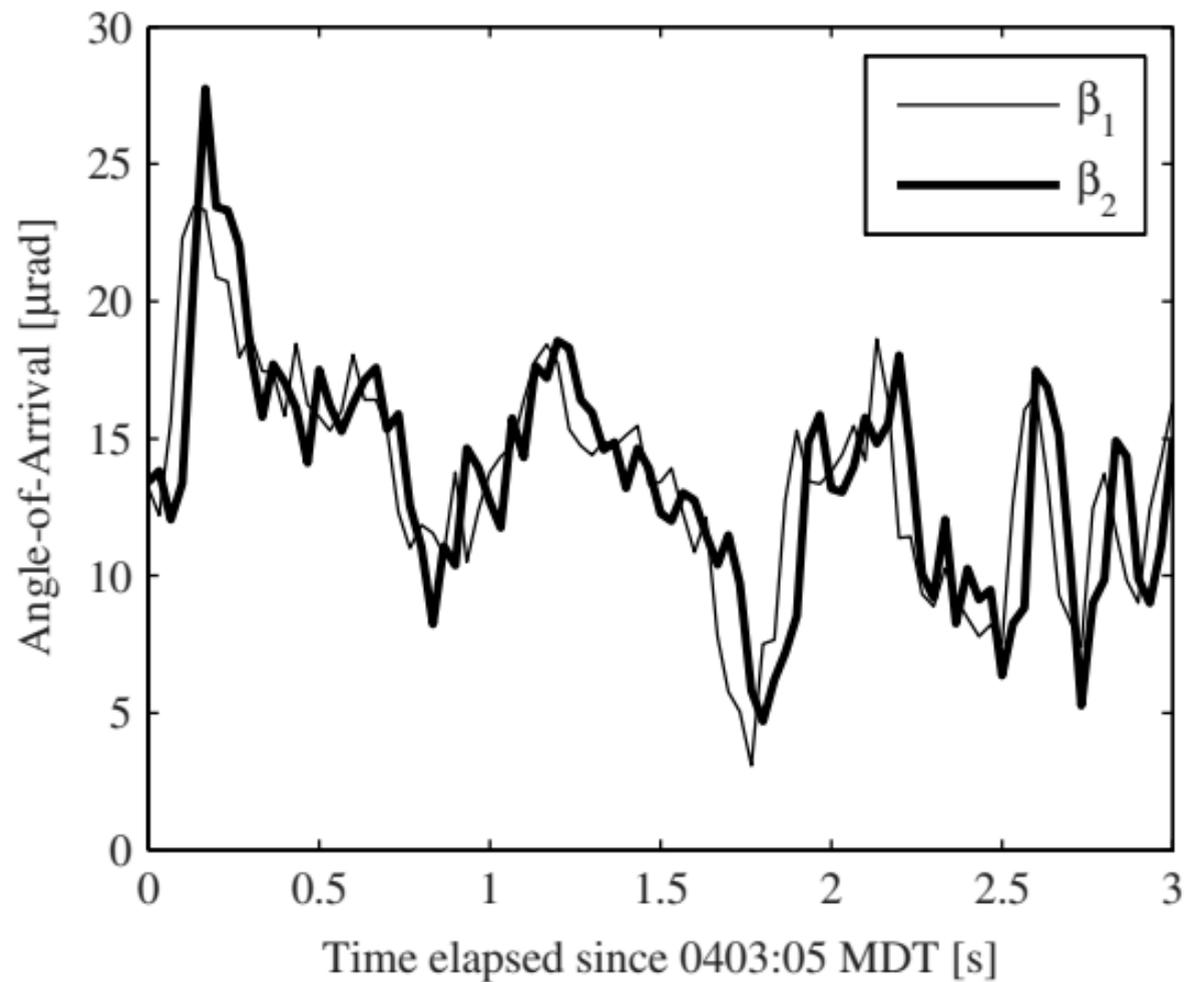
$$S_\beta(f) = \frac{2^{7/3}}{27\pi^{1/6}\Gamma(11/6)} C_n^2 v_b^{-1/3} L f^{-2/3}. \quad (3)$$

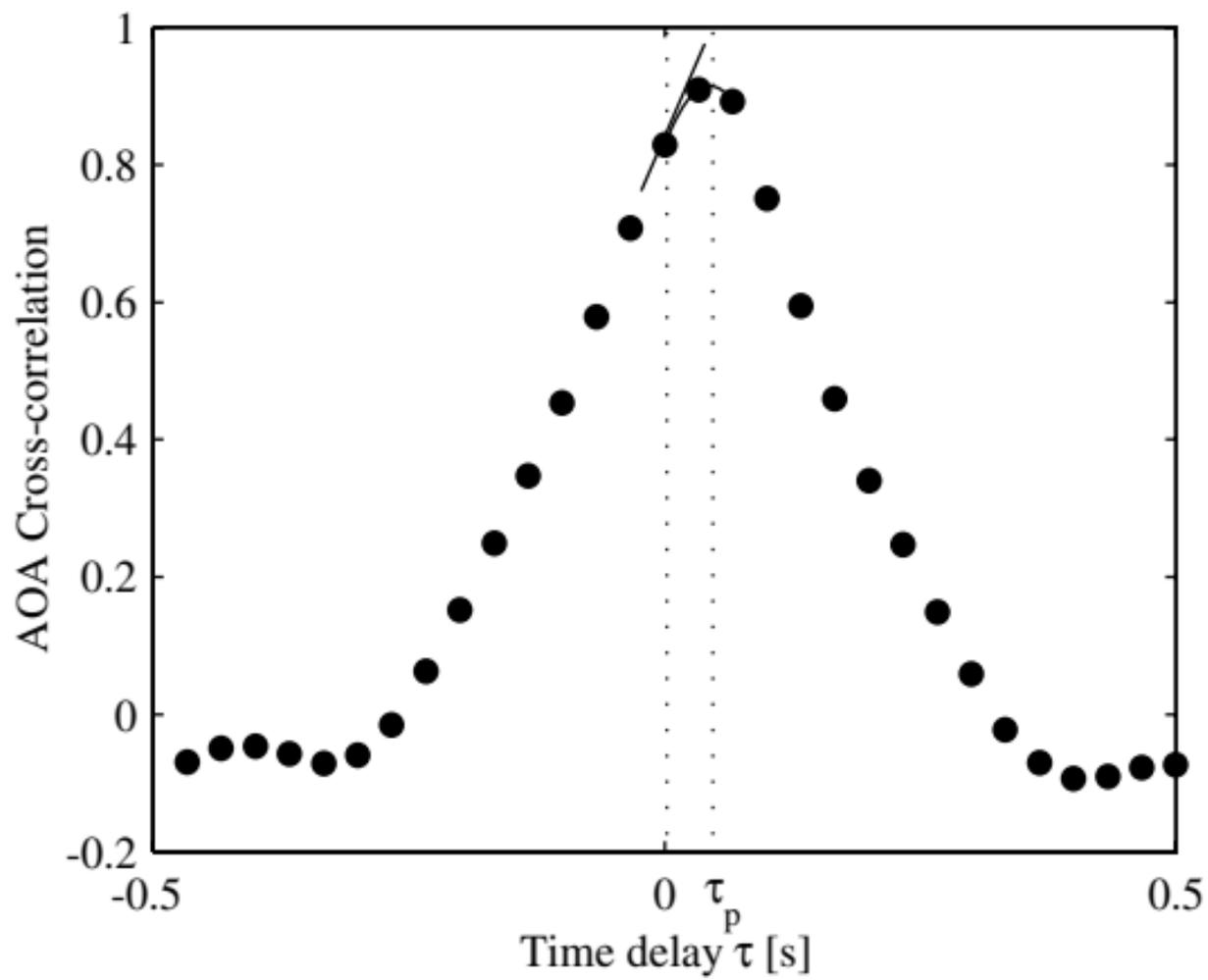
“Knee frequency” (frequency of intersection of the asymptotes):

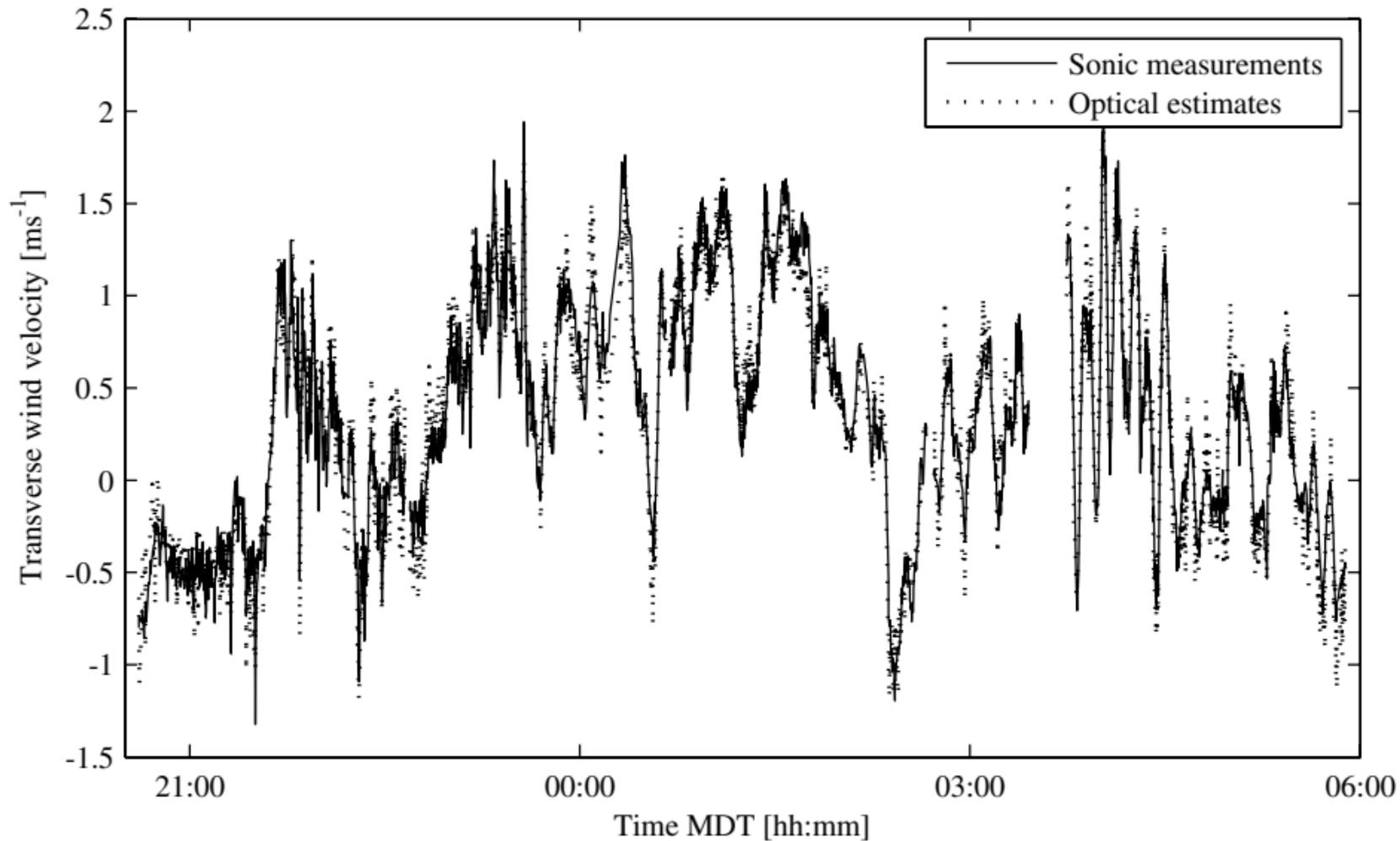
$$v_k = \frac{\sqrt{6} v_b}{2\pi b} = 0.39 \frac{v_b}{b}. \quad (4)$$

Vertical and horizontal AOA spectra









Summary and conclusions

- Our propagation testbed enables us to quantitatively test hypotheses and approximations about optical propagation through the turbulent atmosphere
- Aperture-averaged AOAs (vertical and lateral) and irradiances contain valuable information about turbulent and non-turbulent characteristics in the atmospheric wind and refractive-index fields.

Future work

- Investigate propagation scenarios involving longer ranges and multiple paths (multiple source/receiver combinations)
- Investigate effects of turbulent fluctuations of aerosol concentrations on long-range AOA and irradiance measurements.