

# THE VIRTUAL FIELDS METHOD

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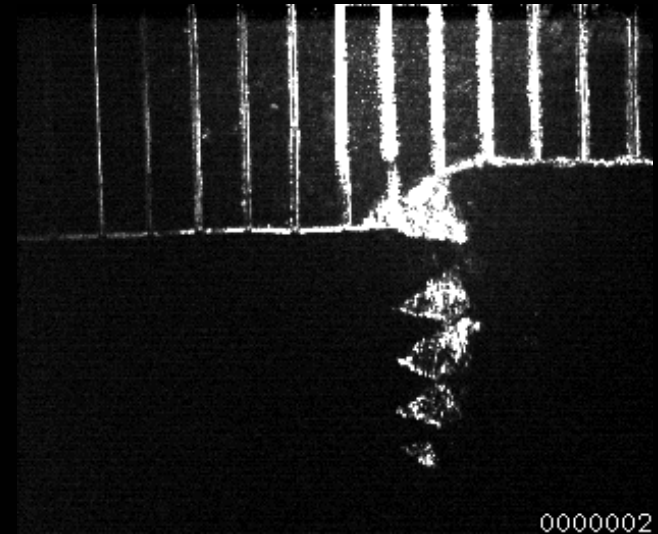
# Course overview

- Introduction and basic idea
- Principle of virtual work
- The Virtual Fields Method (VFM): principle in elasticity
- Complements on the VFM
- Conclusion

# Introduction and basic idea

# Introduction

- Huge progress in computational mechanics
  - Simulation of machining
    - Large strains elasto-plasticity
    - Large strain rates
    - Localization
    - Friction/thermal behaviour
- Problem
  - Many material parameters required
  - How to obtain them?



# Introduction

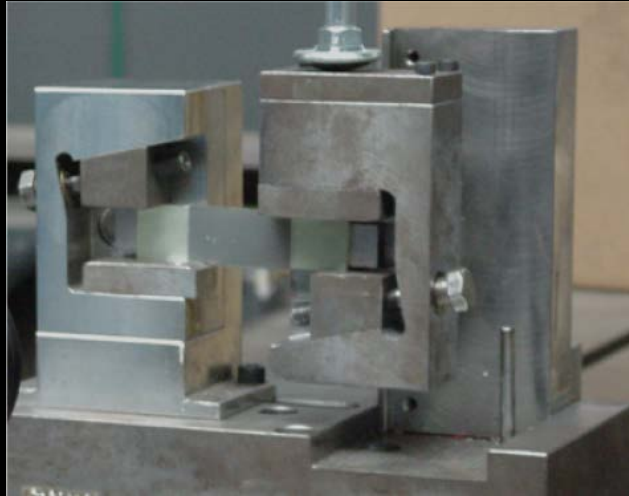
- Standard tests: Tensile test on rectangular specimen
  - Uniform stress state
  - Uniaxial stress strain curve



- Very poor information
- Very restrictive assumptions (constraints)

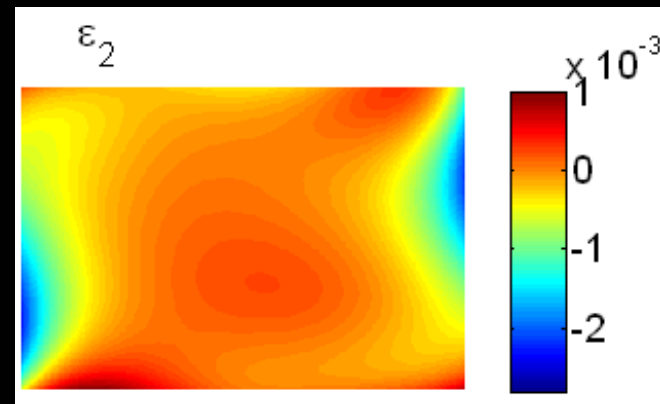
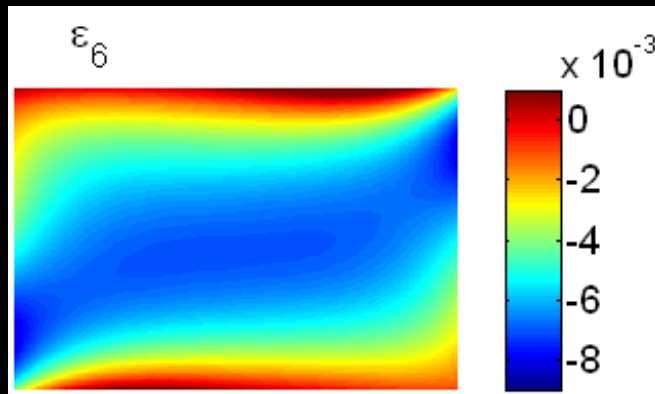
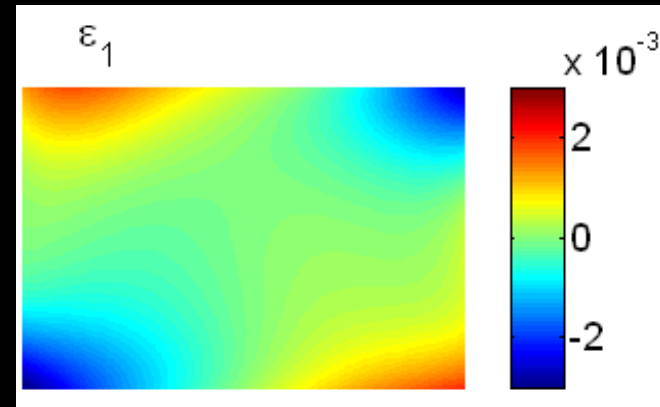
Develop the experimental identification procedures of the future !

# Alternative



Composite specimen

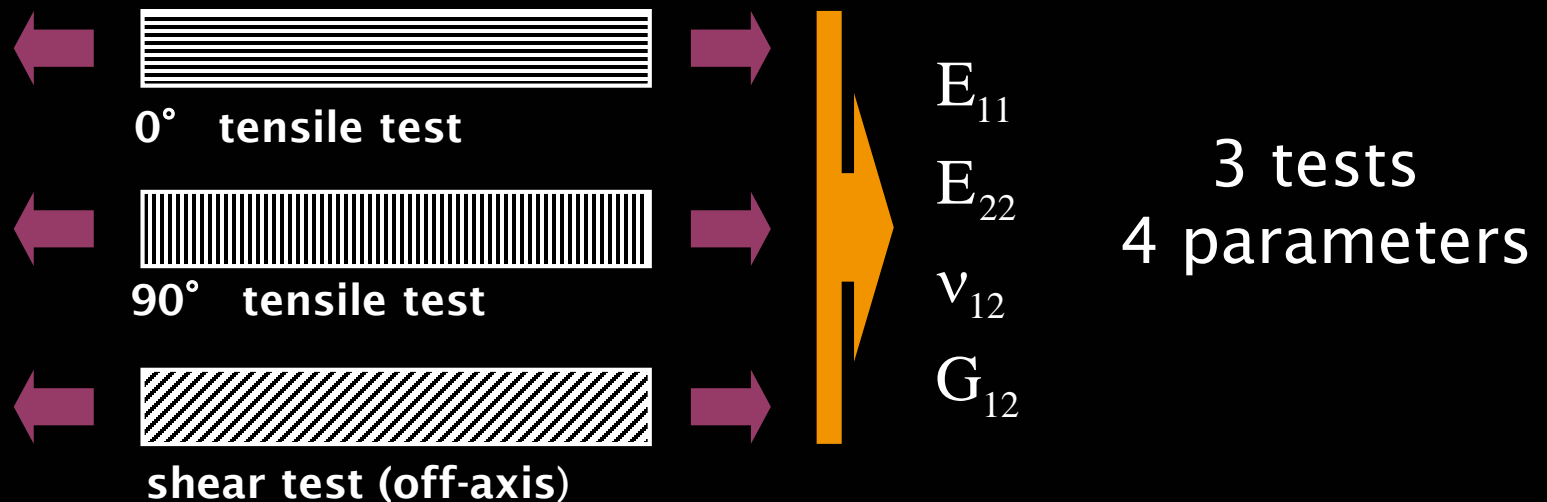
## Grid method



Extract material properties ( $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ ,  $Q_{66}$ ) from strain fields and load

# Motivation

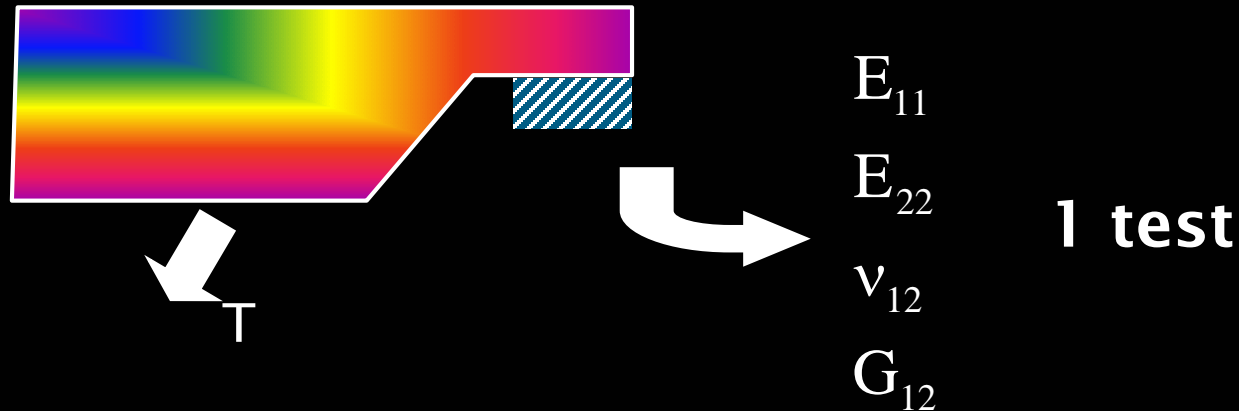
- Extract more information from 1 test



Local strain measurements  
Uniform stress fields (closed-form solution)  
Spatial stress distribution

# Motivation

- Novel strategy for orthotropic elastic moduli measurements

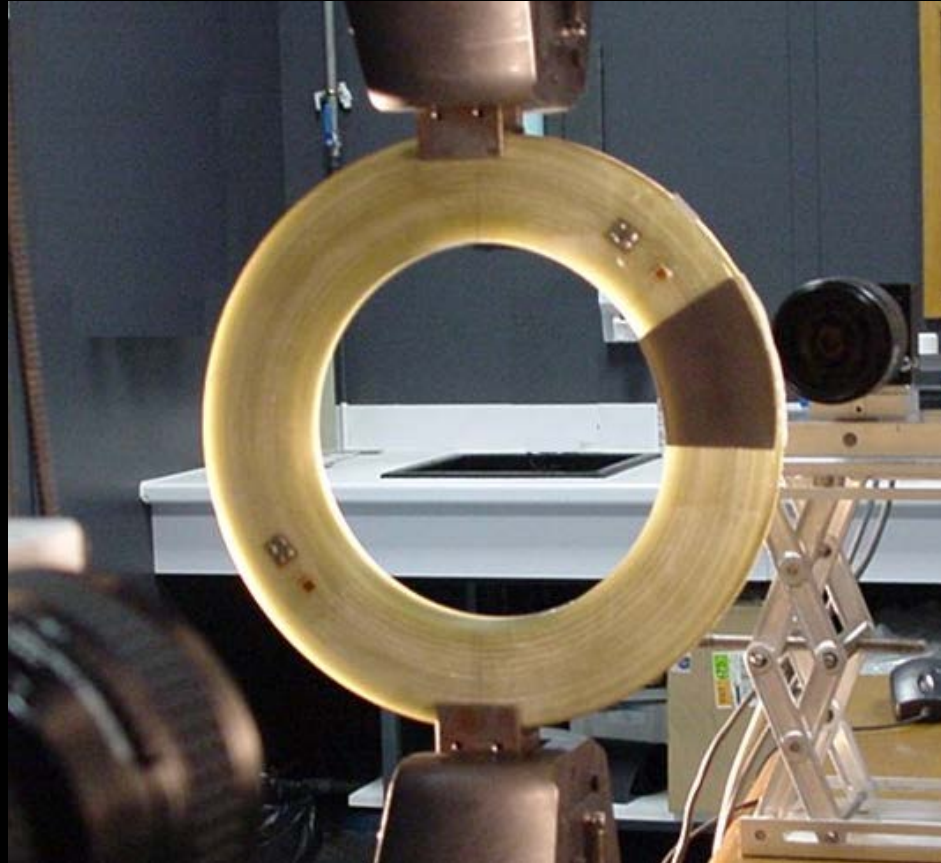


Heterogeneous stress fields  
(no closed-form solution)



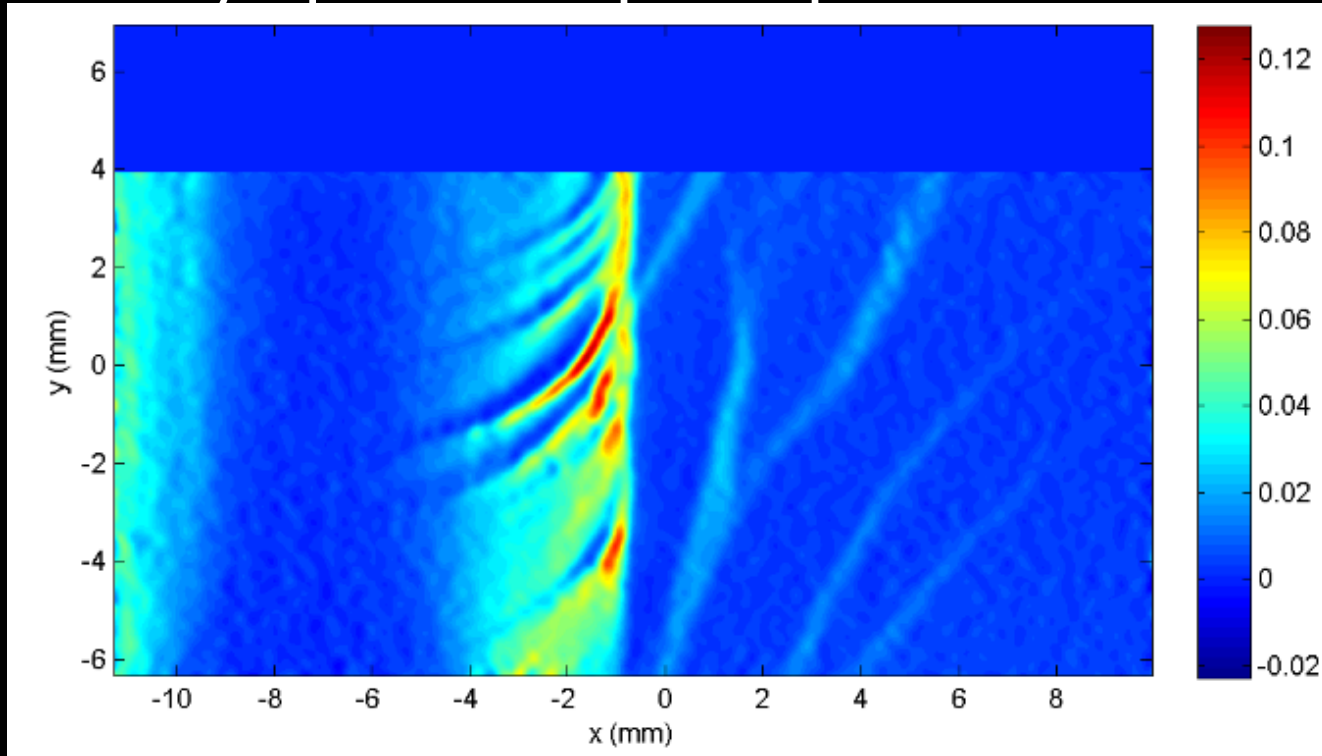
# Motivation

- Complex test geometry



# Motivation

- Heterogeneous materials
  - Identify spatial maps of parameters



Longitudinal strain in a magnesium FSW weld  
(speckle interferometry)

# Key issues: full-field measurements

- Why?
  - Complex strain state
  - Need for more experimental information
- Interferometry (moiré, speckle, etc...)
  - High sensitivity, high spatial resolution
  - Small displacements (rigid body), low strains and high strain gradients
- White light techniques
  - Image correlation : large displacements, 3D
  - Grid method : intermediate

# Key issues: full-field measurements

- Extraordinary potential
  - From a few measurements points to a few 100.000's!!
- Qualitative / quantitative ?
  - Metrological properties not fully assessed yet
  - Often reduced to visual information
  - Very few quantitative use of the data!
- International effort
  - DIC Challenge (Society for Experimental Mechanics) <https://www.sem.org/dic-challenge/>

# Key issues: inverse resolution

- Extensive literature on inverse problems (mathematical)
- Very few tools targeted at full-field processing
- Very few groups with expertise in measurements and inverse problems at the same time
- Inversion closely related to measurement performance!
- Huge investigation area, critically underexplored

# Inverse resolution

## ■ Basic equations

I Equilibrium equations (static)

$$\sigma_{ij,j} + f_i = 0 \quad + \text{boundary conditions} \quad \text{strong (local)}$$

or

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS + \int_V f_i u_i^* dV = 0 \quad \text{weak (global)}$$

II Constitutive equations (elasticity)  $\sigma_{ij} = C_{ijkl} \varepsilon_{kl}$

III Kinematic equations (small strains/displacements)

$$\varepsilon_{ij} = \frac{1}{2}(u_{i,j} + u_{j,i})$$

# Inverse resolution

Known

Unknown

Direct problem

$C_{ijkl}$

Geometry

Boundary conditions

$\sigma_{ij}, \varepsilon_{ij}, u_i$

- Tools for solving this problem
  - Direct integration (closed-form solution)
  - Approximate solutions
  - Galerkin, Ritz
  - Finite elements, boundary elements...
  - etc...

# Inverse resolution

Known

Unknown

Inverse  
problem

$\varepsilon_{ij}, u_i$  (measured)

Geometry

Some information on the  
boundary conditions (load cell)

$C_{ijkl}$

$\sigma_{ij}$

- Tools for solving this problem
  - Statically determined tests:  
Closed form solution of Eq. I (uncoupled system)  
Force BC, simple geometry  
Ex.: tensile test, bending tests (on rect. beams)  
etc...



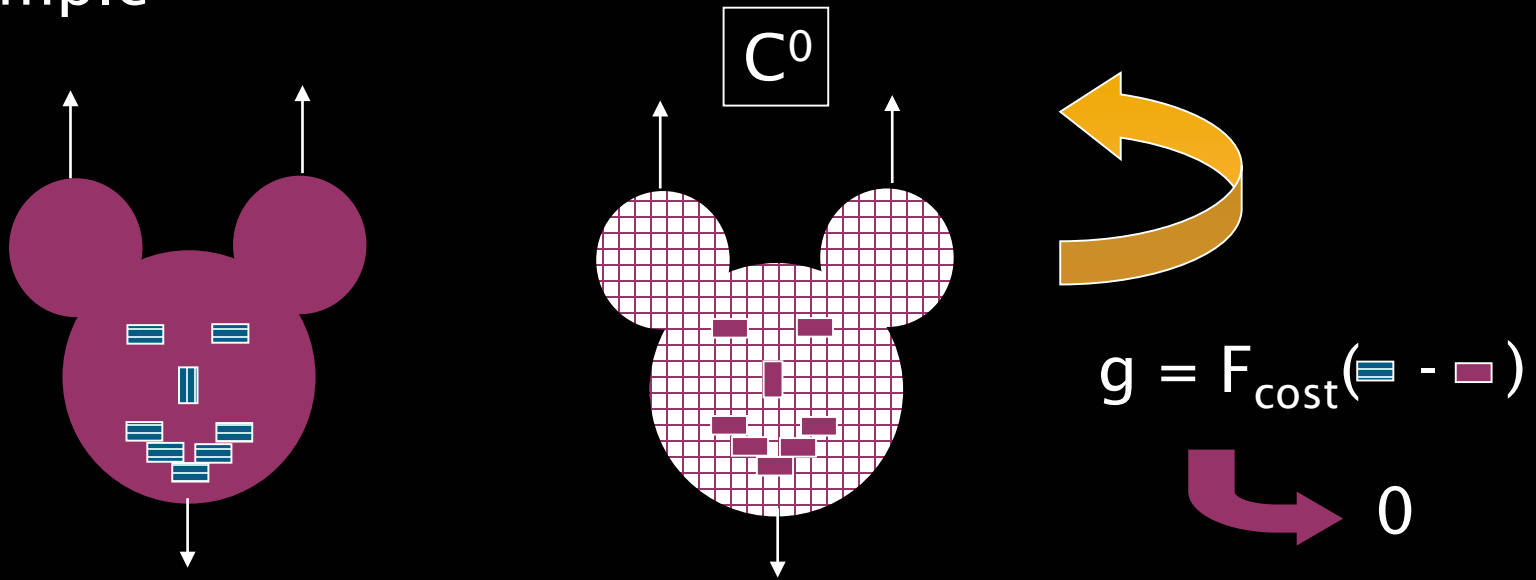
# Inverse resolution

- Tools for solving this problem

- Model updating

Idea: iterative use of tool for direct problem  
(analytical or approximate)

Example



# Inverse resolution

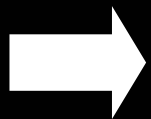
- Model updating

- Advantages

- General method (full-field measurements not compulsory)
    - Tools already developed

- Shortcomings

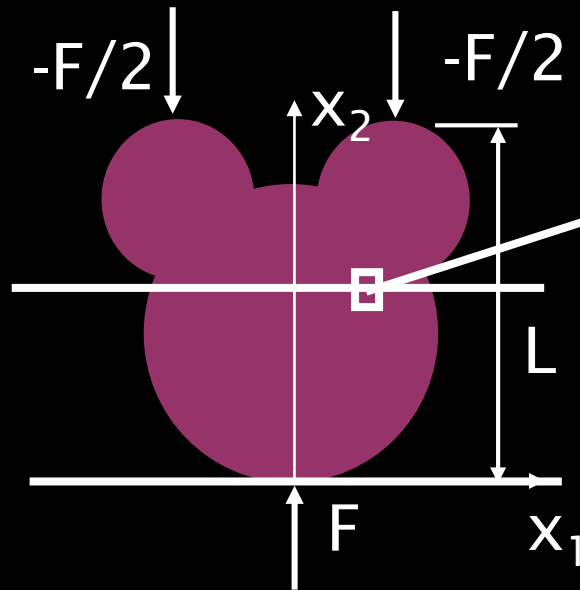
- Sensitive to boundary conditions (generally badly known)
    - CPU intensive (for numerical approximations and non-linear equations...)
    - Not fully dedicated to full-field measurements



Alternative tool: the Virtual Fields Method

# Basic idea

- Use global equations (and not local)



$$\varepsilon_{22}(x_1^0, x_2^0) \quad \sigma_{22}(x_1^0, x_2^0)?$$

$$\int_S \sigma_{22} dx_1 dx_3 = -F$$

Integrate over  $x_2$

$$\int_V \sigma_{22} dx_1 dx_2 dx_3 = -FL$$

# Basic idea

- Constitutive behaviour

$$\begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{pmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & 0 \\ Q_{12} & Q_{11} & 0 \\ 0 & 0 & \frac{Q_{11} - Q_{12}}{2} \end{bmatrix} \begin{pmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ 2\varepsilon_{12} \end{pmatrix} \quad \text{In-plane linear elastic isotropy}$$

$$\int_V \sigma_{22} dV = -FL \quad \longrightarrow \quad \int_V (Q_{12}\varepsilon_{11} + Q_{11}\varepsilon_{22}) dV = -FL$$

Material is homogeneous

$$Q_{11} \int_V \varepsilon_{22} dV + Q_{12} \int_V \varepsilon_{11} dV = -FL$$

# Basic idea

## ■ Surface measurements only

Constant strains through the thickness

$$Q_{11} \int_S \epsilon_{22} dx_1 dx_2 + Q_{12} \int_S \epsilon_{11} dx_1 dx_2 = -\frac{FL}{e}$$

$$\int_S \epsilon_{22} dx_1 dx_2 \approx \sum_{i=1}^n \epsilon_{22}^i s^i$$

$s^i$  is the surface of each pixel  
 $n$  is the number of strain data points

If all pixels have the same size  $s$  (usually the case for CCD/CMOS based measurements)

$$\sum_{i=1}^n \epsilon_{22}^i s^i = s \sum_{i=1}^n \epsilon_{22}^i = \frac{S_d}{n} \sum_{i=1}^n \epsilon_{22}^i = S_d \overline{\epsilon_{22}} \quad \overline{\epsilon_{22}} = \frac{1}{n} \sum_{i=1}^n \epsilon_{22}^i$$

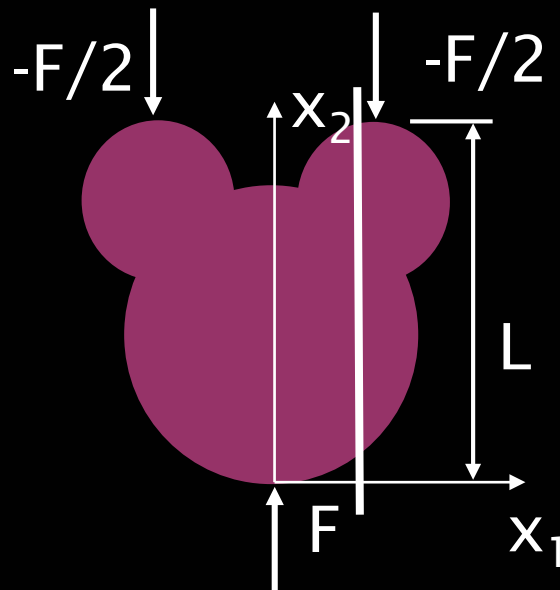
$S_d$  is the surface of the disc

# Basic idea

## ■ Finally

$$Q_{22} \bar{\varepsilon}_{22} + Q_{12} \bar{\varepsilon}_{11} = \frac{-FL}{eS_d}$$

$$Q_{11} \bar{\varepsilon}_{11} + Q_{12} \bar{\varepsilon}_{22} = 0$$



$$\int_S \sigma_{11} dx_2 dx_3 = 0$$

Integrate over  $x_1$

$$\int_V \sigma_{11} dV = 0$$

$$\begin{bmatrix} \bar{\varepsilon}_{11} & \bar{\varepsilon}_{22} \\ \bar{\varepsilon}_{22} & \bar{\varepsilon}_{11} \end{bmatrix} \begin{pmatrix} Q_{11} \\ Q_{12} \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{-FL}{eS_d} \end{pmatrix}$$

$$Q_{11} = \frac{-FL \bar{\varepsilon}_{22}}{eS_d (\bar{\varepsilon}_{22}^2 - \bar{\varepsilon}_{11}^2)}$$

$$Q_{12} = \frac{FL \bar{\varepsilon}_{11}}{eS_d (\bar{\varepsilon}_{22}^2 - \bar{\varepsilon}_{11}^2)}$$

# Basic idea

- What if the disc material is anisotropic?
  - Need for more equations
- Tool to generate integral equilibrium equations: principle of virtual work

# The principle of virtual work



# The Principle of Virtual Work

Equilibrium  
equation

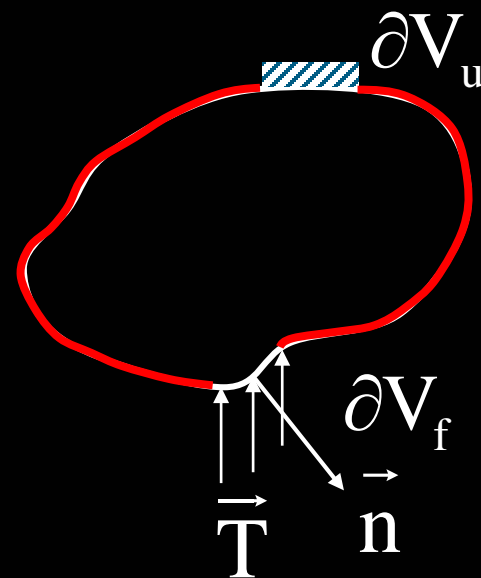
$$\overrightarrow{\text{div}} \overline{\overline{\sigma}} + \vec{f} = \vec{0}$$

$$\sigma_{ij,j} + f_i = 0$$

Boundary  
conditions

$$T_i = \sigma_{ij} n_j \text{ on } \partial V_f$$

$$\mathbf{u} = \overline{\mathbf{u}} \text{ on } \partial V_u$$



Let  $u_i^*$  be a vectorial function

$$\sigma_{ij,j} u_i^* + f_i u_i^* = 0$$

Integration over the whole domain

$$\int_V \sigma_{ij,j} u_i^* dV + \int_V f_i u_i^* dV = 0$$

# The Principle of Virtual Work

- Integration by part and use of divergence theorem

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS + \int_V f_i u_i^* dV = 0$$

$u_i^*$  Virtual displacement field (cont. and diff)

$\varepsilon_{ij}^* = \frac{1}{2}(u_{i,j}^* + u_{j,i}^*)$  virtual strain tensor

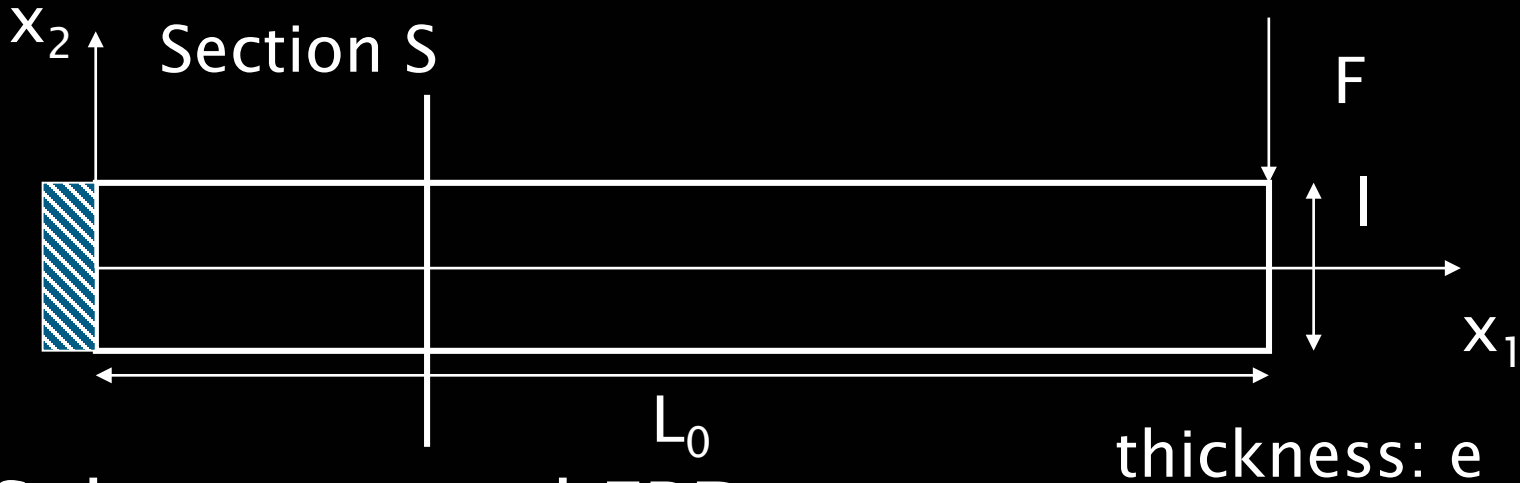
$T_i$  Applied stress vector on solid boundary  $\partial V_f$

$f_i$  External volume forces

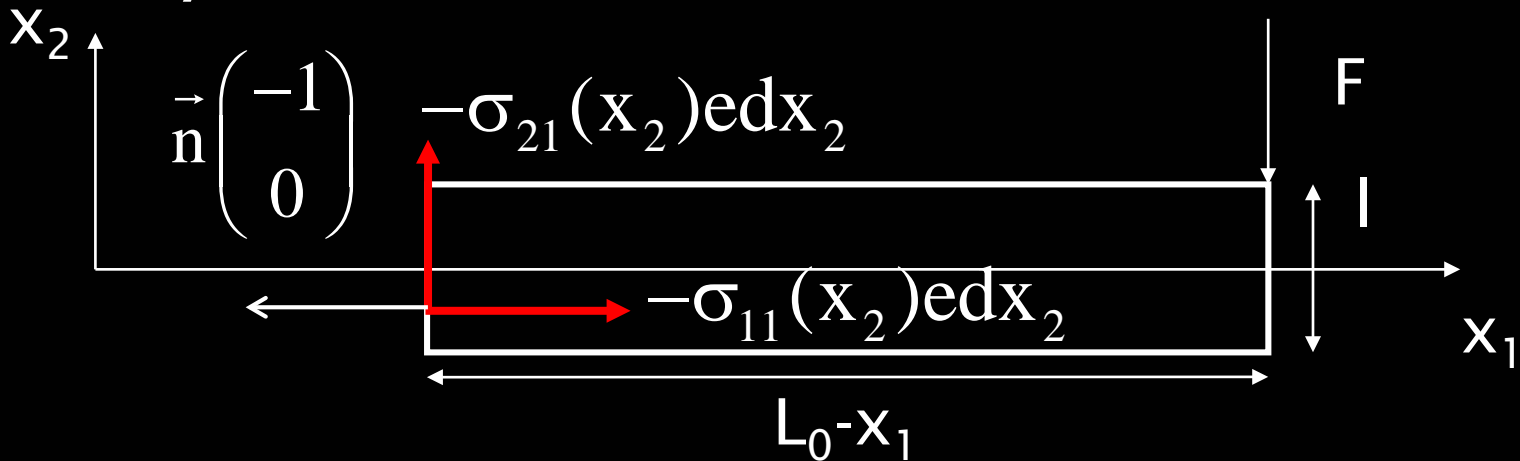
$\sigma_{ij}$  Stress tensor

# Illustration of the PVW

## ■ Cantilever beam

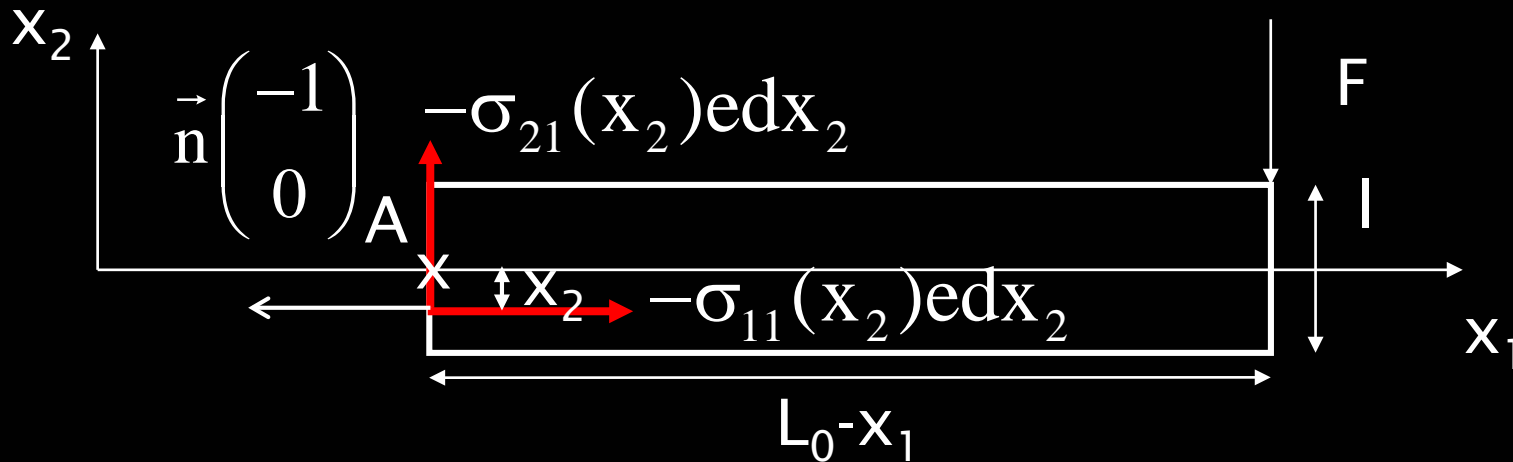


## ■ Subsystem and FBD



# Illustration of the PVW

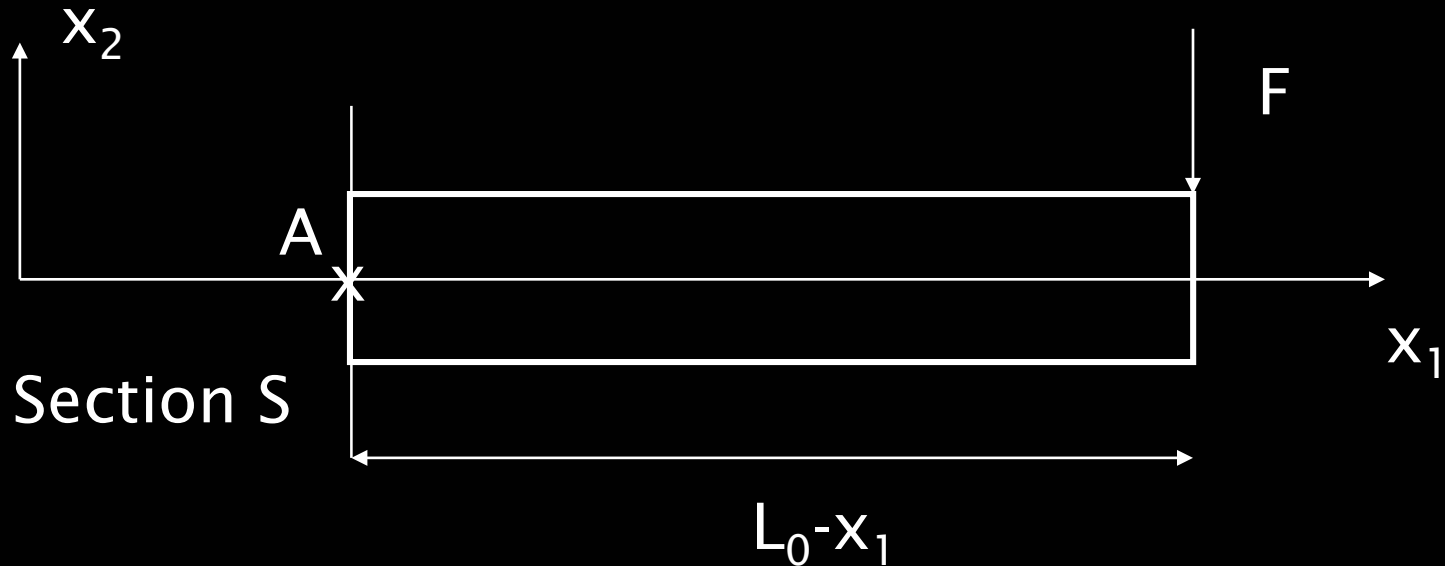
## ■ Resultant of internal forces



$$\vec{F}_i = \begin{pmatrix} -e \int_{-1/2}^{1/2} \sigma_{11} dx_2 \\ -e \int_{-1/2}^{1/2} \sigma_{21} dx_2 \end{pmatrix} \quad M_{i/A}^{x_3} = e \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_2$$

# Illustration of the PVW

- Resultant of external force



$$\vec{F}_e = \begin{pmatrix} 0 \\ -F \end{pmatrix} ; M_{e/A}^{x_3} = -F(L_0 - x_1)$$

# Illustration of the PVW

## ■ Equilibrium

$$\sum \vec{F} = \vec{0} \quad \Rightarrow \quad \begin{pmatrix} -e \int_{-1/2}^{1/2} \sigma_{11} dx_2 \\ -e \int_{-1/2}^{1/2} \sigma_{21} dx_2 \end{pmatrix} + \begin{pmatrix} 0 \\ -F \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$
$$\sum M^{x_3} = 0 \quad \Rightarrow \quad F(L_0 - x_1) - e \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_2 = 0$$

$$\int_{-1/2}^{1/2} \sigma_{11} dx_2 = 0$$

$$e \int_{-1/2}^{1/2} \sigma_{21} dx_2 = -F$$

$$e \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_2 = F(L_0 - x_1)$$

# Illustration of the PVW

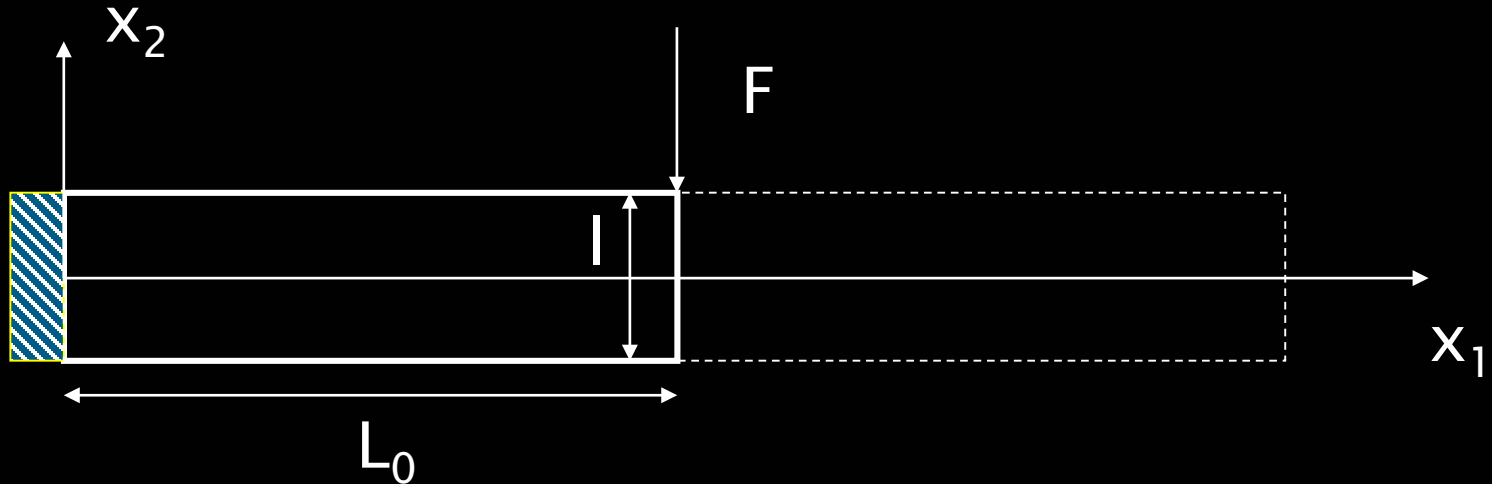
- Valid over any section S of the beam:
  - integration over  $x_1$

$$\int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} dx_1 dx_2 = 0 \quad \text{Eq. 1}$$

$$e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{21} dx_1 dx_2 = -FL_0 \quad \text{Eq. 2}$$

$$e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_1 dx_2 = \frac{FL_0^2}{2} \quad \text{Eq. 3}$$

# Illustration of the PVW



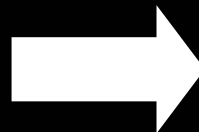
Principle of virtual work (static, no volume forces)

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS = 0$$

Let us write a virtual field:

$$u_1^* = x_1$$

$$u_2^* = 0$$



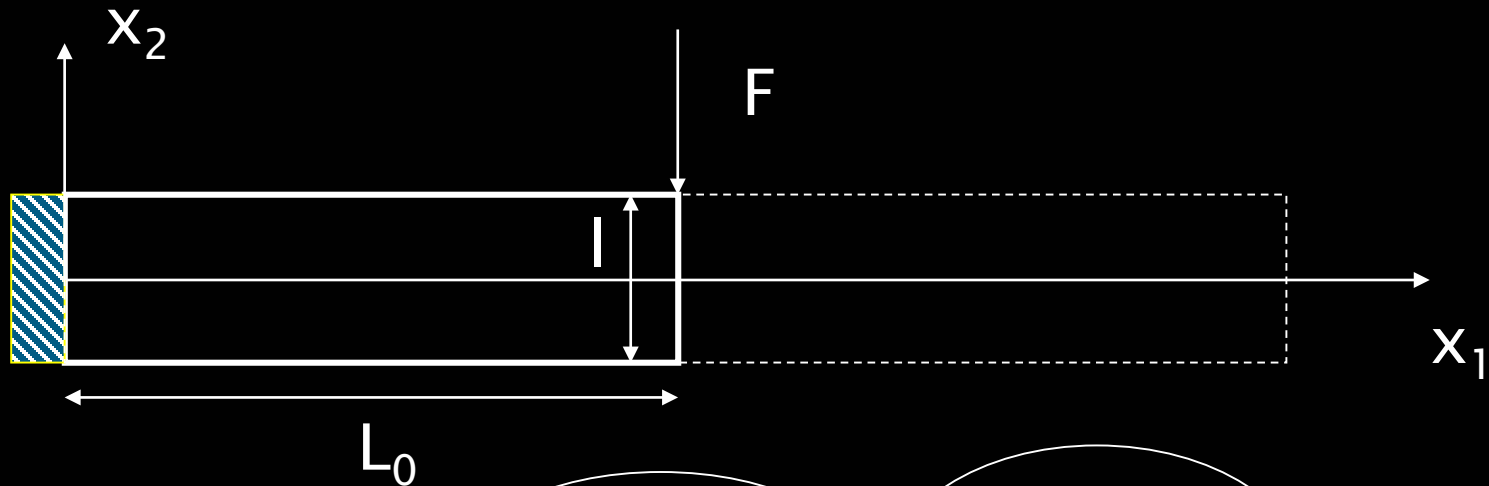
$$\varepsilon_{11}^* = 1$$

$$\varepsilon_{22}^* = 0$$

$$\varepsilon_{12}^* = 0$$



# Illustration of the PVW

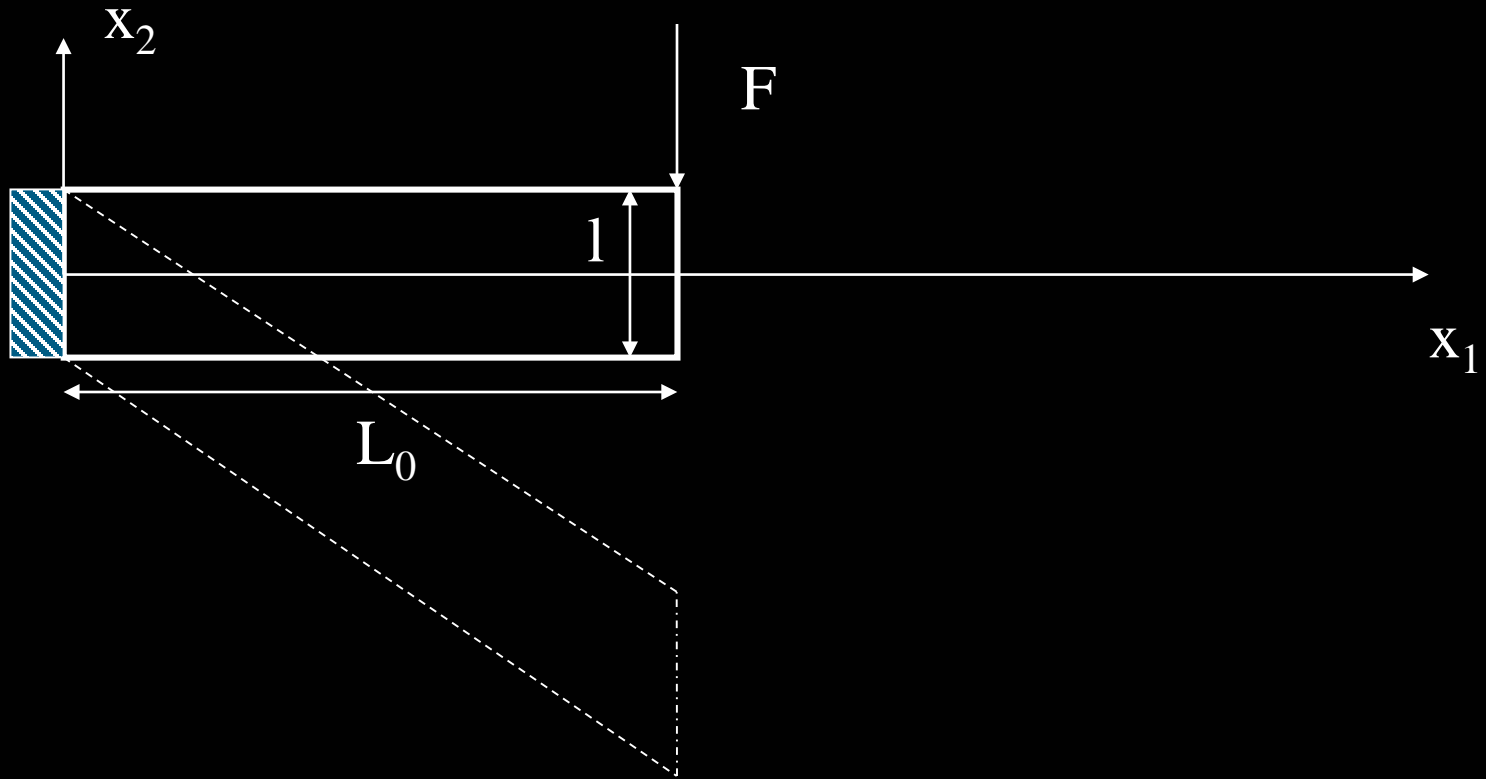


$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS = 0$$

$$-\int_V \sigma_{11} \varepsilon_{11}^* dV = -e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} dx_1 dx_2 = 0$$

$$\int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} dx_1 dx_2 = 0 \quad \text{Eq. 1}$$

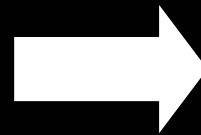
# Illustration of the PVW



Let us write another virtual field:

$$u_1^* = 0$$

$$u_2^* = -x_1$$

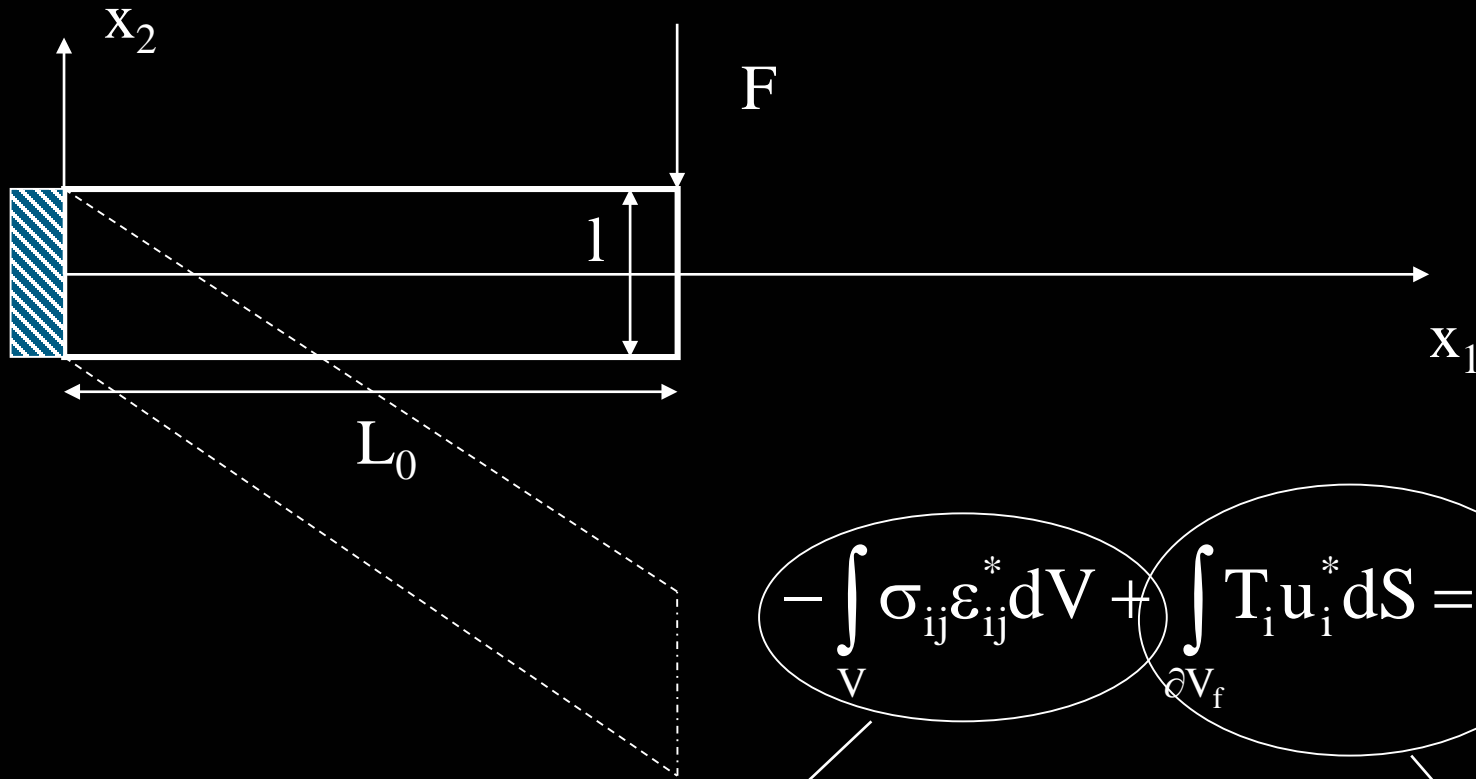


$$\varepsilon_{11}^* = 0$$

$$\varepsilon_{22}^* = 0$$

$$\varepsilon_{12}^* = -1/2$$

# Illustration of the PVW

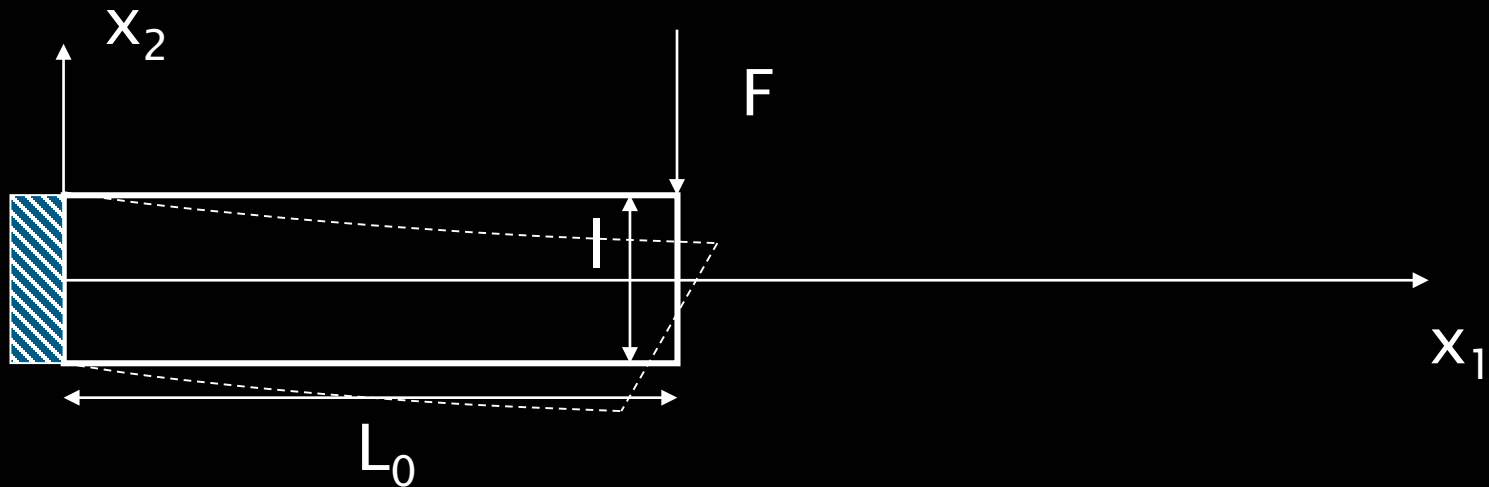


$$-\int_V 2\sigma_{12} \epsilon_{12}^* dV = e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{12} dx_1 dx_2$$

$$e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{12} dx_1 dx_2 = -FL_0$$

Eq. 2

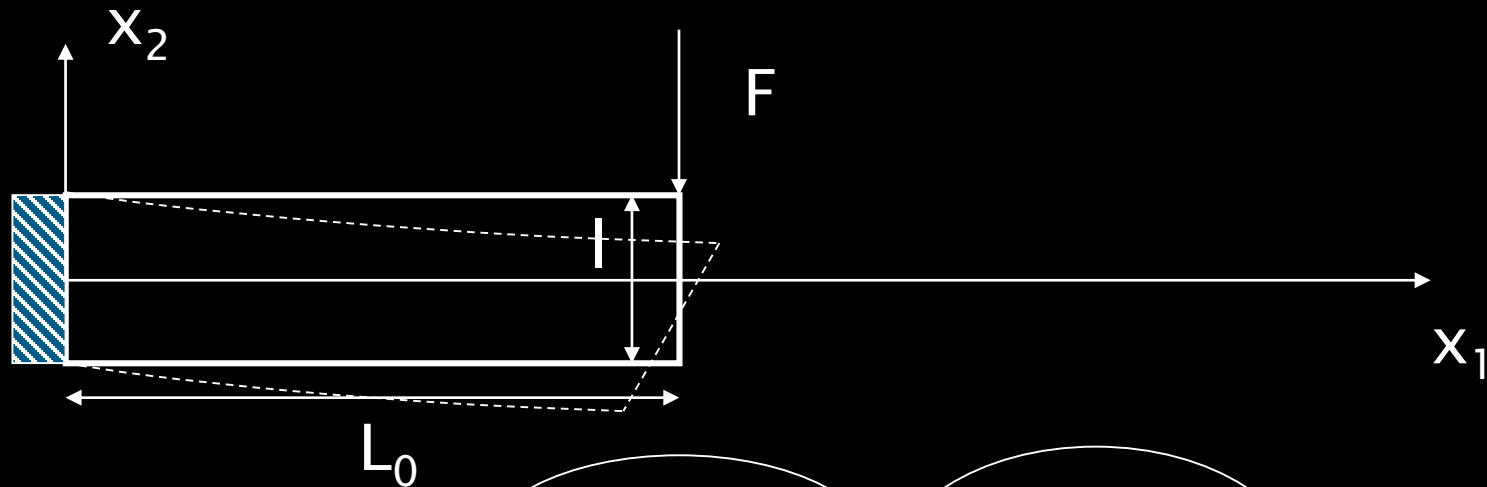
# Illustration of the PVW



Let us write a 3rd field: virtual bending

$$\begin{aligned}
 u_1^* &= x_1 x_2 \\
 u_2^* &= -\frac{x_1^2}{2}
 \end{aligned}
 \quad \longrightarrow \quad
 \begin{aligned}
 \varepsilon_{11}^* &= x_2 \\
 \varepsilon_{22}^* &= 0 \\
 \varepsilon_{12}^* &= 0
 \end{aligned}$$

# Illustration of the PVW



$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS = 0$$

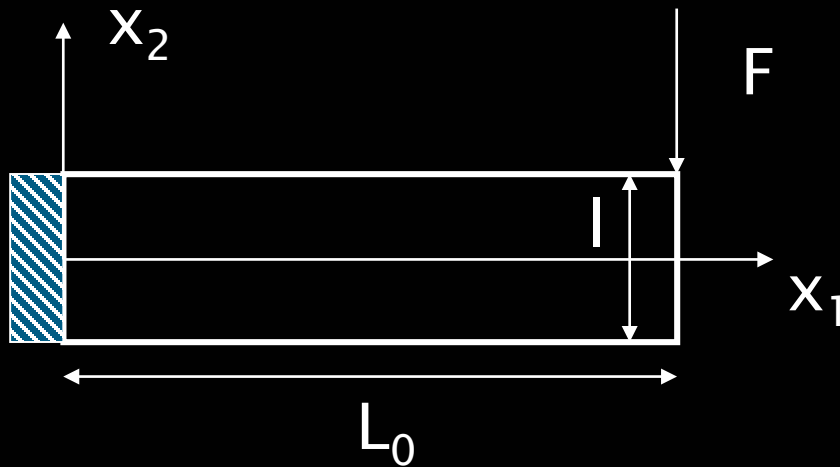
$$-\int_V \sigma_{11} \varepsilon_{11}^* dV = -e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_1 dx_2$$

$$= \frac{F L_0^2}{2}$$

$$e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} x_2 dx_1 dx_2 = \frac{F L_0^2}{2}$$

Eq. 3

# Illustration of the PVW



Any virtual field can be selected

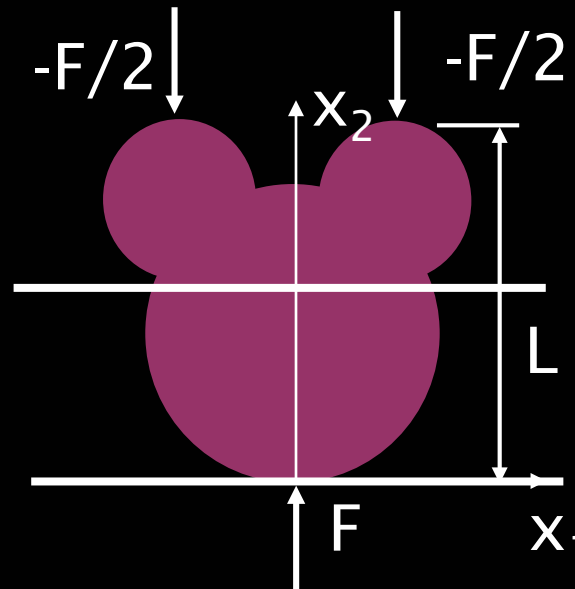
$$\begin{cases} u_1^* = 1 - e^{x_1} \\ u_2^* = 0 \end{cases} \quad \longrightarrow \quad \begin{aligned} \varepsilon_{11}^* &= -e^{x_1} \\ \varepsilon_{22}^* &= 0 \\ \varepsilon_{12}^* &= 0 \end{aligned}$$

$$e \int_0^{L_0} \int_{-1/2}^{1/2} \sigma_{11} e^{x_1} dx_1 dx_2 = 0$$

# The Virtual Fields Method

# The Virtual Fields Method

- Back to the funny shaped disc



- Principle of virtual work

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V} T_i u_i^* dS = 0 \quad \text{No volume forces, static}$$



# The Virtual Fields Method

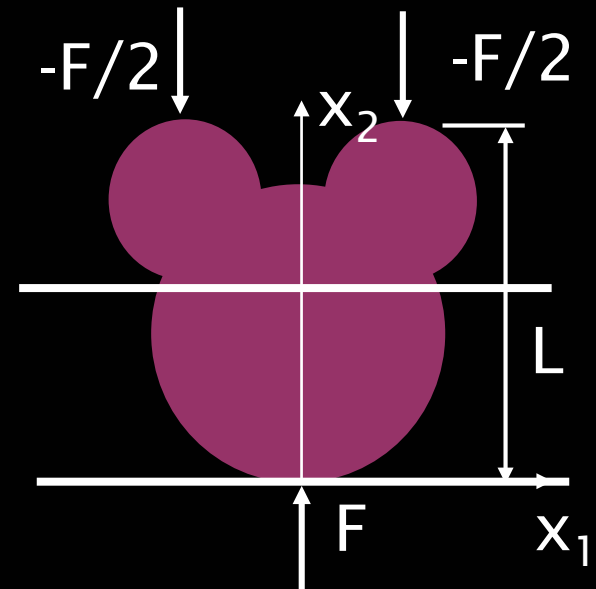
- Let's write two virtual fields
  - 1<sup>st</sup> virtual field: virtual compression field

$$u_1^* = 0 ; u_2^* = -x_2$$

$$\varepsilon_{11}^* = 0 ; \varepsilon_{22}^* = -1 ; \varepsilon_{12}^* = 0$$

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV = \int_V \sigma_{22} dV$$

$$\int_{\partial V_f} T_i u_i^* dS = FL \quad \int_V \sigma_{22} dV = -FL$$



# The Virtual Fields Method

- 2<sup>nd</sup> virtual field: transverse shrinking

$$u_1^* = -x_1 ; u_2^* = 0 \quad \varepsilon_{11}^* = -1 ; \varepsilon_{22}^* = 0 ; \varepsilon_{12}^* = 0$$

$$-\int_V \sigma : \varepsilon^* dV = \int_V \sigma_{11} dV \quad \int_{\partial V_f} \vec{T} \cdot \vec{u}^* dS = 0 \quad \Rightarrow \quad \int_V \sigma_{11} dV = 0$$

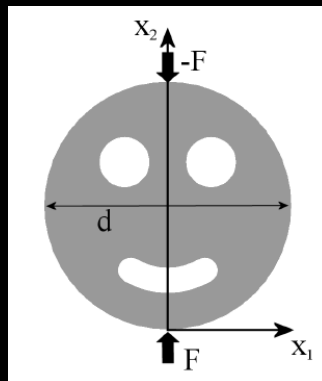
- Other virtual fields (an infinity!)

$$u_1^* = 0 ; u_2^* = -\sin\left(\frac{\pi}{2L} x_2\right)$$

$$Q_{11} \left[ \varepsilon_{22} \cos\left(\frac{\pi}{2L} x_2\right) \right] + Q_{12} \left[ \varepsilon_{11} \cos\left(\frac{\pi}{2L} x_2\right) \right] = \frac{-2FL}{\pi e S_d}$$

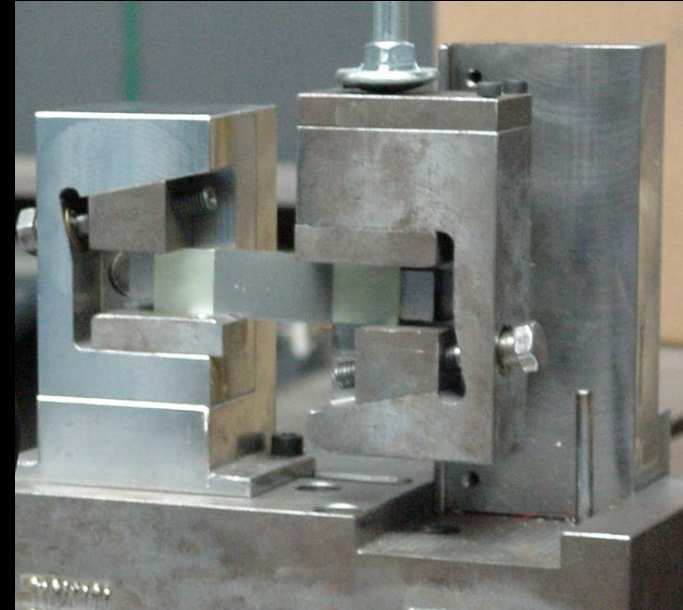
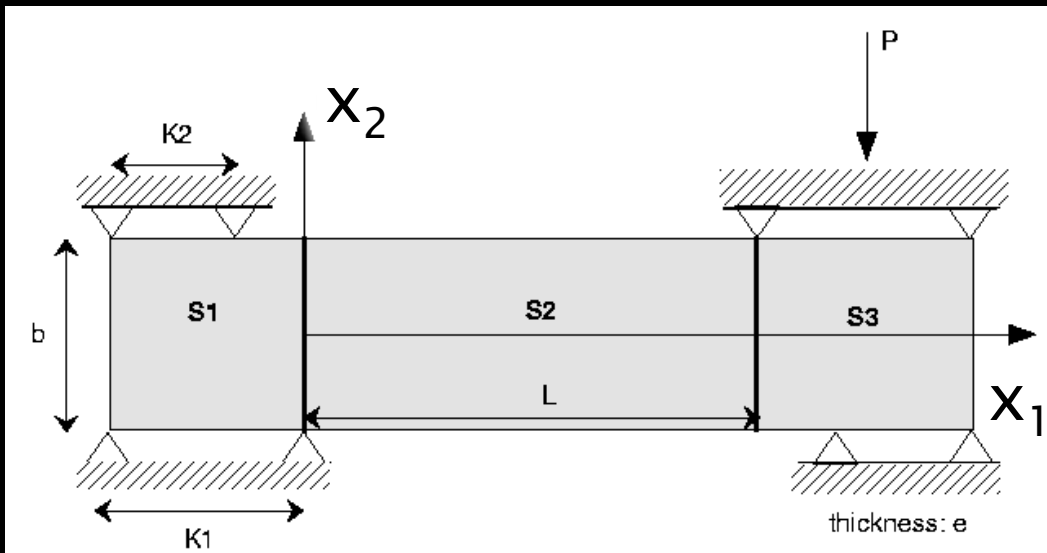
# The Virtual Fields Method

- Principal advantages
  - Independent from stress distribution; etc. geometry
  - Direct identification in elasticity (no updating)
- Limitations
  - Kinematic assumption through the thickness (plane stress, plane strain, bending...)



# VFM: more complex example

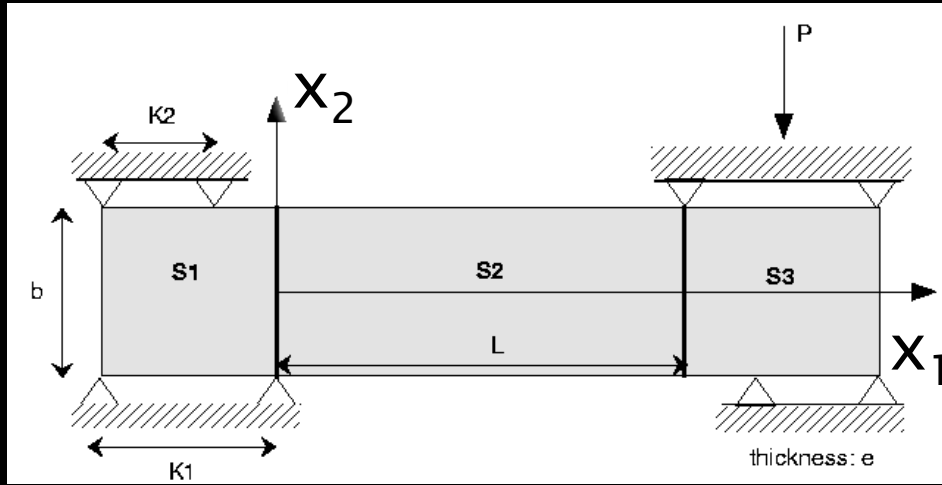
- Orthotropic elasticity



Orthotropic material

$$\begin{pmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{pmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & 0 \\ Q_{12} & Q_{22} & 0 \\ 0 & 0 & Q_{66} \end{bmatrix} \begin{pmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ 2\varepsilon_{12} \end{pmatrix}$$

# VFM: more complex example



Choice of the virtual fields

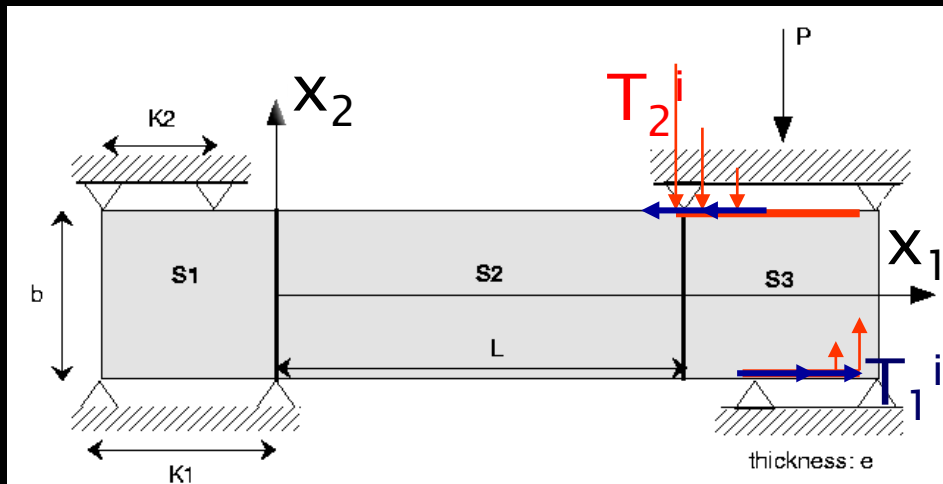
1. Measurement on  $S_2$  only (optical system)

Over  $S_1$  and  $S_3$ :  $\varepsilon_{11}^* = 0$  ;  $\varepsilon_{22}^* = 0$  ;  $\varepsilon_{12}^* = 0$  (rigid body)

2. A priori choice:

$$\text{over } S_1: u_1^* = 0 ; u_2^* = 0 \quad \Rightarrow \quad \int_{S_1} T_i u_i^* dS = 0$$

# VFM: more complex example



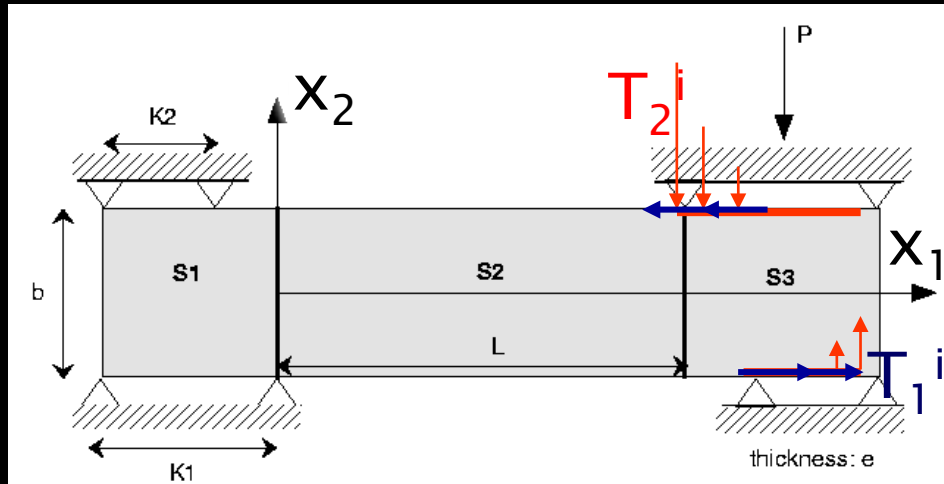
Unknown force distribution over  $S_1$  and  $S_3$ .  
Resultant  $P$  measured

3. Over  $S_3$  (rigid body) : 2 possibilities

$$3.1 \quad u_1^* = 0; u_2^* = 0 \quad \Rightarrow \quad \int_{S_3} T_i u_i^* dS = 0$$

$$3.2 \quad \int_{S_3} T_i u_i^* dx_1 dx_3 = e \left[ \int_{\text{blue}} T_1 u_1^* dx_1 + \int_{\text{red}} T_2 u_2^* dx_1 \right]$$

# VFM: more complex example



No information on  $t_1$

$$\Rightarrow u_1^* = 0$$

$$\int_{S_3} T_i u_i^* dx_1 dx_2 = e \left[ \int_{\text{red}} T_2 u_2^* dx_1 \right]$$

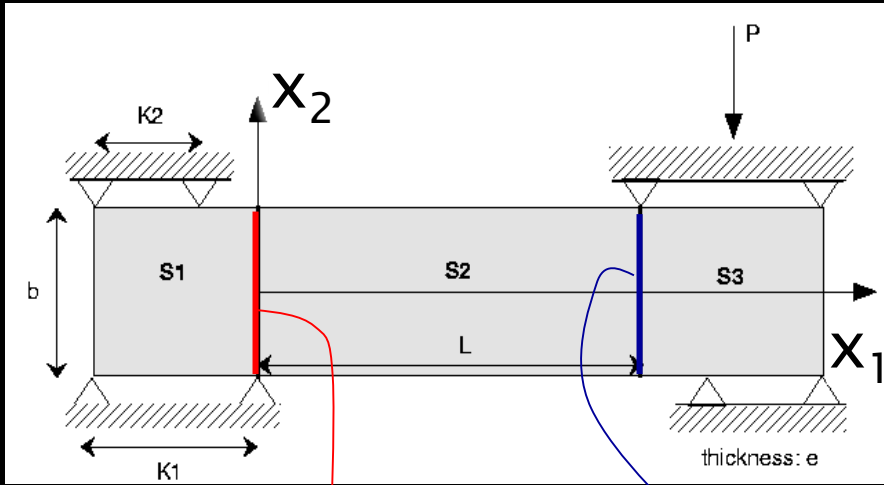
Distribution  $T_2$   
unknown

$$\Rightarrow u_2^* = k$$

Filtering capacity  
of the VF

$$\int_{\text{red}} T_2 u_2^* dx_1 = \int_{\text{red}} T_2 k dx_1 = k \int_{\text{red}} T_2 dx_1 = -kP$$

# VFM: more complex example



4. Continuity of the virtual displacement field  
Conditions over  $S_2$

$$u_1^* = u_2^* = 0$$

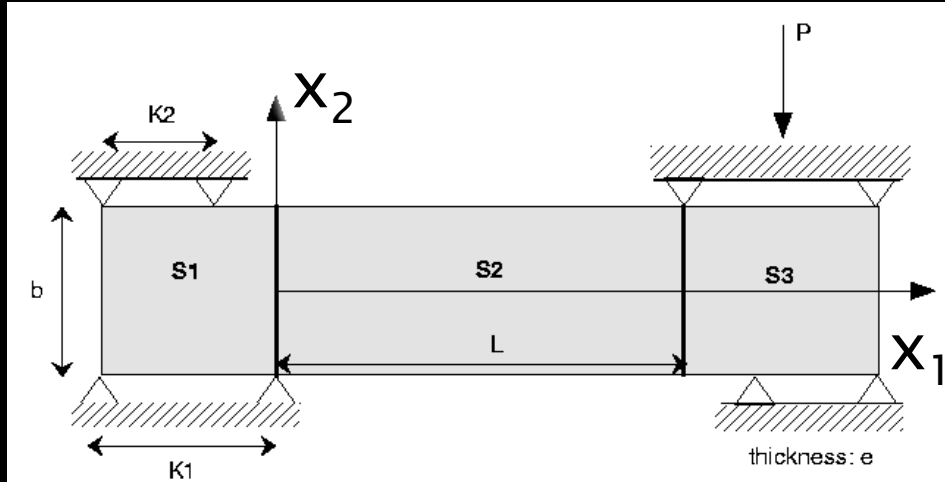
$$u_1^* = 0; u_2^* = k$$

Virtual strain field discontinuous

Choice of 4 virtual fields at least: example



# VFM: more complex example

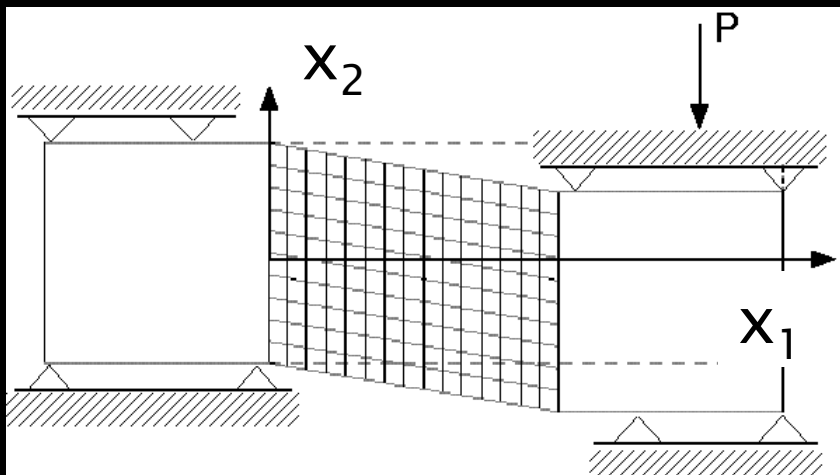


Over  $S_2$

$$u_1^* = 0; u_2^* = -X_1$$

$$\varepsilon_{11}^* = 0; \varepsilon_{22}^* = 0; \varepsilon_{12}^* = -\frac{1}{2}$$

Over  $S_3$        $k = -L$



Uniform virtual shear

# VFM: more complex example

Plane stress orthotropic elasticity

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS = 0$$

↓ Plane stress

$$\int_V \sigma_{12} dV \rightarrow e \int_S \sigma_{12} dS$$

↓ Homogeneous material

$$2e \int_S Q_{66} \varepsilon_{12} dS \rightarrow 2e Q_{66} \int_S \varepsilon_{12} dS$$

$$Q_{66} \int_{S_2} \varepsilon_{12} dx_1 dx_2 = \frac{-PL}{2e}$$

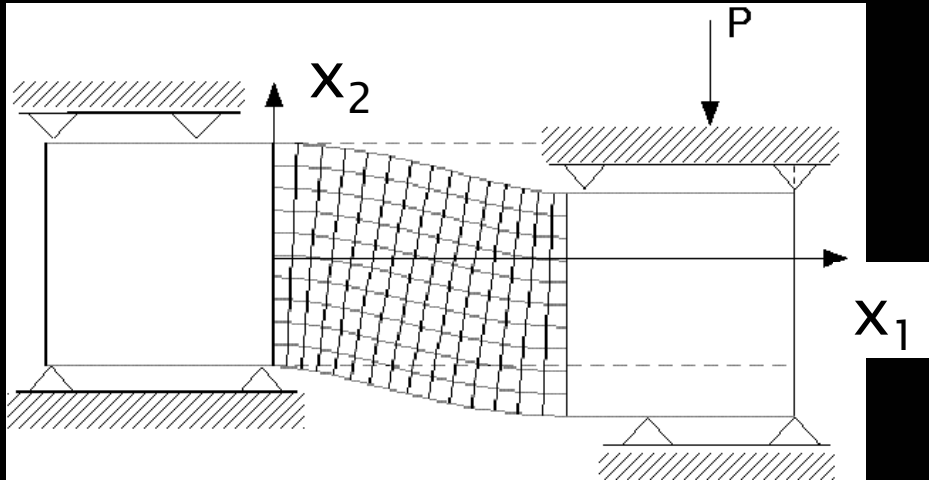
$$\boxed{Q_{66} \bar{\varepsilon}_{12} = \frac{-P}{2eb}}$$

$$\int_{\partial V_f} T_i u_i^* dS$$

↓ PL

# VFM: more complex example

Field n° 2: Bernoulli bending



Over  $S_2$

$$\begin{cases} u_1^* = 6x_1x_2(L - x_1) \\ u_2^* = x_1^2(2x_1 - 3L) \end{cases}$$

$$\begin{cases} \varepsilon_{11}^* = 6x_2(L - 2x_1) \\ \varepsilon_{22}^* = 0 ; \varepsilon_{12}^* = 0 \end{cases}$$

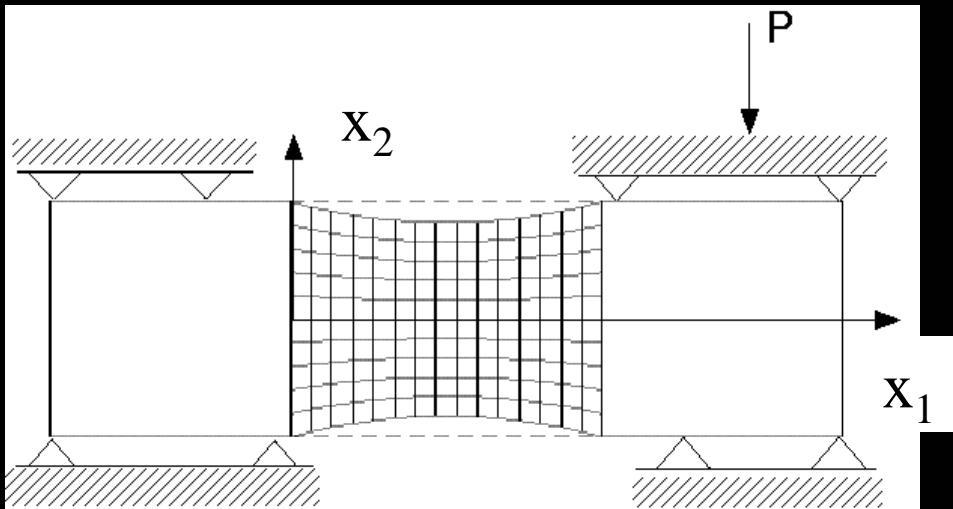
Over  $S_3$

$$k = -L^3$$

$$\overline{Q_{11} 6x_2(L - 2x_1) \varepsilon_{11}} + \overline{Q_{12} 6x_2(L - 2x_1) \varepsilon_{22}} = \frac{PL^2}{eb}$$

# VFM: more complex example

Field n° 3: Global compression



Over  $S_2$

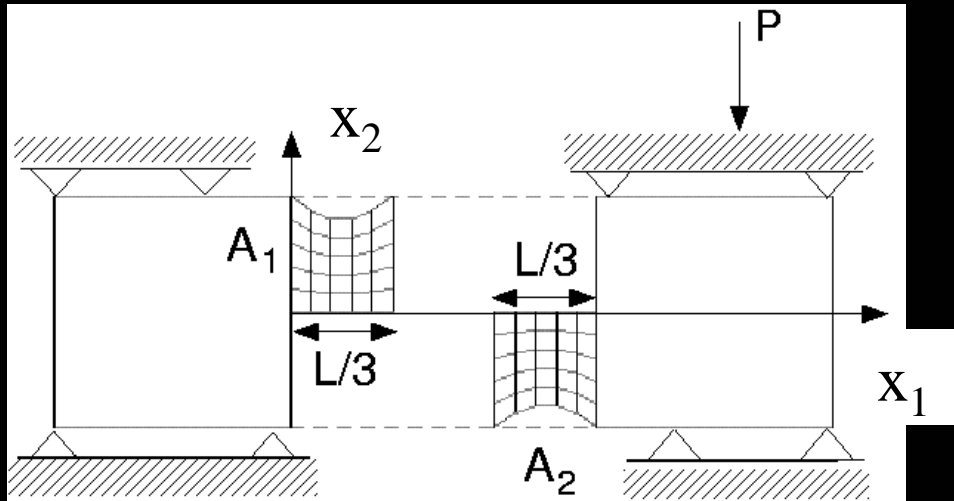
$$\begin{cases} u_1^* = 0 \\ u_2^* = x_1 x_2 (x_1 - L) \\ \varepsilon_{11}^* = 0 \\ \varepsilon_{22}^* = x_1 (x_1 - L) \\ \varepsilon_{12}^* = \frac{1}{2} x_2 (2x_1 - L) \end{cases}$$

$$\overline{Q_{22} x_1 (x_1 - L) \varepsilon_{22}} + \overline{Q_{12} x_1 (x_1 - L) \varepsilon_{11}} + \overline{Q_{66} 2x_2 (2x_1 - L) \varepsilon_{12}} = 0$$

Over  $S_3$   
 $k = 0$

# VFM: more complex example

Field n° 4: Local compression



Over  $A_1$

$$\begin{cases} u_1^* = 0 \\ u_2^* = x_1 x_2 (x_1 - L/3) \\ \varepsilon_{11}^* = 0 \\ \varepsilon_{22}^* = x_1 (x_1 - L/3) \\ \varepsilon_{12}^* = \frac{1}{2} x_2 (2x_1 - L/3) \end{cases}$$

Over  $S_3$   $k = 0$

Over  $A_2$

$$\begin{cases} u_1^* = 0 \\ u_2^* = x_2 (x_1 - L)(x_1 - 2L/3) \end{cases}$$

$$\begin{cases} \varepsilon_{11}^* = 0 \\ \varepsilon_{22}^* = (x_1 - L)(x_1 - L/3) \\ \varepsilon_{12}^* = \frac{1}{2} x_2 (2x_1 - 5L/3) \end{cases}$$

# VFM: more complex example

Field n° 4: Local compression

$$\begin{aligned} & Q_{22} \left[ \overline{x_1(x_1 - L)\varepsilon_{22}}^{A_1} + \overline{(x_1 - L)(x_1 - L/3)\varepsilon_{22}}^{A_2} \right] \\ & + Q_{12} \left[ \overline{x_1(x_1 - L)\varepsilon_{11}}^{A_1} + \overline{(x_1 - L)(x_1 - L/3)\varepsilon_{11}}^{A_2} \right] \\ & + Q_{66} \left[ \overline{2x_2(2x_1 - L/3)\varepsilon_{12}}^{A_1} + \overline{2x_2(2x_1 - 5L/3)\varepsilon_{12}}^{A_2} \right] = 0 \end{aligned}$$

# VFM: more complex example

Final system

$$\begin{bmatrix} 0 & 0 & 0 & \overline{\varepsilon_{12}} \\ \overline{6x_2(L-2x_1)\varepsilon_{11}} & 0 & \overline{6x_2(L-2x_1)\varepsilon_{22}} & 0 \\ 0 & \overline{x_1(x_1-L)\varepsilon_{22}} & \overline{x_1(x_1-L)\varepsilon_{11}} & \overline{x_2(2x_1-L)\varepsilon_{12}} \\ 0 & \left\{ \overline{x_1(x_1-L)\varepsilon_{22}}^{A_1} + \overline{(x_1-L)(x_1-L/3)\varepsilon_{22}}^{A_2} \right\} & \left\{ \overline{x_1(x_1-L)\varepsilon_{11}}^{A_1} + \overline{(x_1-L)(x_1-L/3)\varepsilon_{11}}^{A_2} \right\} & \left\{ \overline{x_2(2x_1-L/3)\varepsilon_{12}}^{A_1} + \overline{x_2(2x_1-5L/3)\varepsilon_{12}}^{A_2} \right\} \end{bmatrix} \begin{pmatrix} Q_{11} \\ Q_{22} \\ Q_{12} \\ Q_{66} \end{pmatrix} = \begin{pmatrix} -\frac{P}{2eb} \\ \frac{PL^2}{eb} \\ 0 \\ 0 \end{pmatrix}$$

$$AQ = B \quad \Rightarrow \quad Q = A^{-1}B \quad \text{If VF independent !!}$$

Pierron F. et Grédiac M., *Composites Part A*, vol. 31, pp. 309-318, 2000.

# Complements on the VFM



# Summary

- Virtual fields selection in elasticity
- Heterogeneous materials
- Non-linear constitutive models
- Dynamics
- Force identification

# VF selection

# VF selection

- Exact data
  - All sets of VF produce the SAME results
- Noisy data
  - Different sets provide different results
  - Optimal selection?

# VF selection

- Strategy for virtual fields selection in elasticity
  - Projection of VF on a certain basis of functions
  - Polynomial

$$\begin{cases} u_1^* = \sum_{i=0}^m \sum_{j=0}^n a_{ij} \left( \frac{x_1}{L} \right)^i \left( \frac{x_2}{b} \right)^j \\ u_2^* = \sum_{i=0}^m \sum_{j=0}^n b_{ij} \left( \frac{x_1}{L} \right)^i \left( \frac{x_2}{b} \right)^j \end{cases}$$

- Choice of VF: choice of a set of  $m, n, a_{ij}, b_{ij}$

# VF selection

- Strategy to find  $a_{ij}, b_{ij}$  for given  $m$  and  $n$ 
  - Special virtual fields
  - For instance, special virtual field to identify  $Q_{11}$

$$Q_{11} \underbrace{\overline{\varepsilon_{11} \varepsilon_{11}^*}}_{=1} + Q_{22} \underbrace{\overline{\varepsilon_{22} \varepsilon_{22}^*}}_{=0} + Q_{12} \underbrace{\left( \overline{\varepsilon_{11} \varepsilon_{22}^*} + \overline{\varepsilon_{22} \varepsilon_{11}^*} \right)}_{=0} + 4Q_{66} \underbrace{\overline{\varepsilon_{12} \varepsilon_{12}^*}}_{=0} = -\frac{Pu_2^*(x_1 = L)}{Lbe}$$

$$Q_{11} = -\frac{Pu_2^{*(1)}(x_1 = L)}{Lbe}$$

- 4 linear equations linking the  $a_{ij}$  and  $b_{ij}$  coefficients
- Virtual BCs: a few more linear equations
- Still a very large number of VF in general!

# VF selection

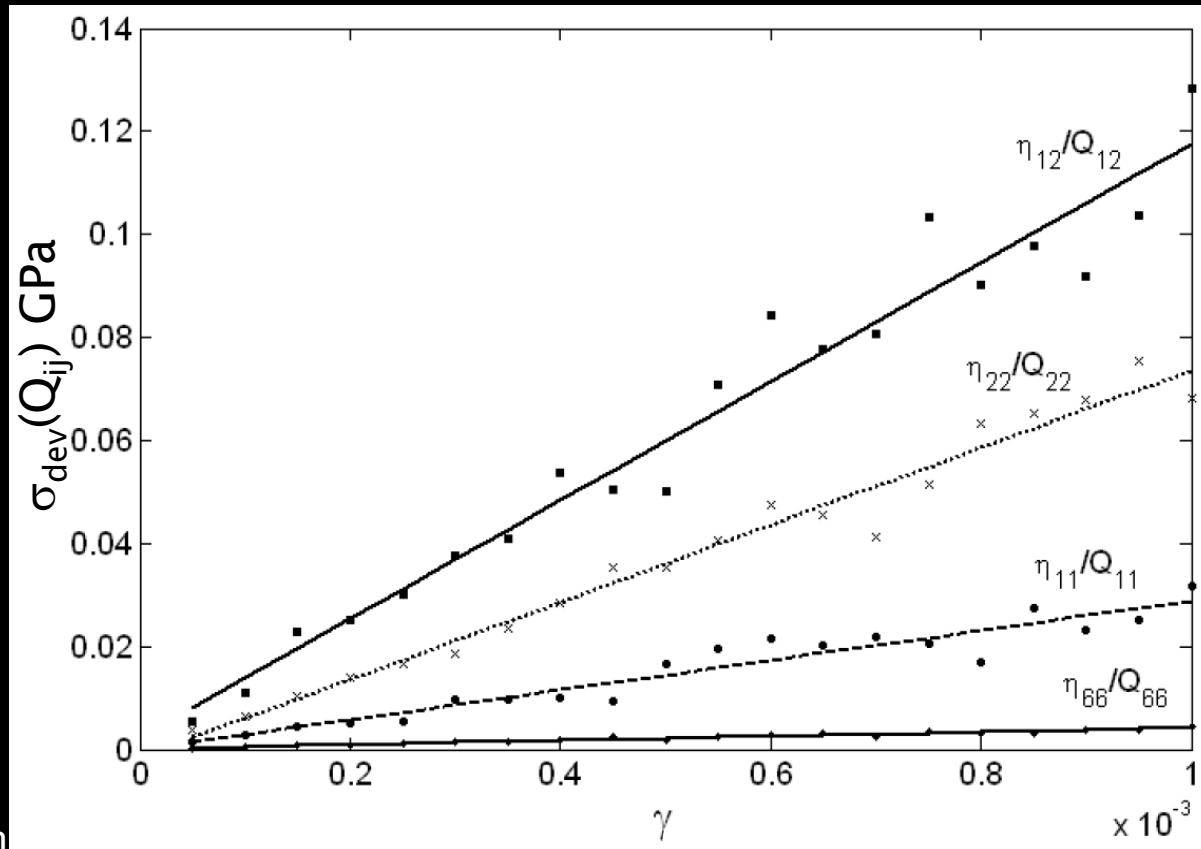
- Strategy to find  $a_{ij}, b_{ij}$  for given  $m$  and  $n$ 
  - Minimization of the sensitivity to noise
  - Constrained minimization problem where  $a_{ij}$  and  $b_{ij}$  coefficients are the unknown
  - Provides the maximum likelihood solution, with:

$$\left\{ \begin{array}{l} \sigma_{\text{dev}}(Q_{11}) = \eta_{11}\gamma \\ \sigma_{\text{dev}}(Q_{22}) = \eta_{22}\gamma \\ \sigma_{\text{dev}}(Q_{12}) = \eta_{12}\gamma \\ \sigma_{\text{dev}}(Q_{66}) = \eta_{66}\gamma \end{array} \right. \quad \begin{array}{l} \gamma \text{ is the standard} \\ \text{deviation of strain noise} \\ \eta_{ij} \text{ are the sensitivity to} \\ \text{noise coefficients} \end{array}$$

# VF selection

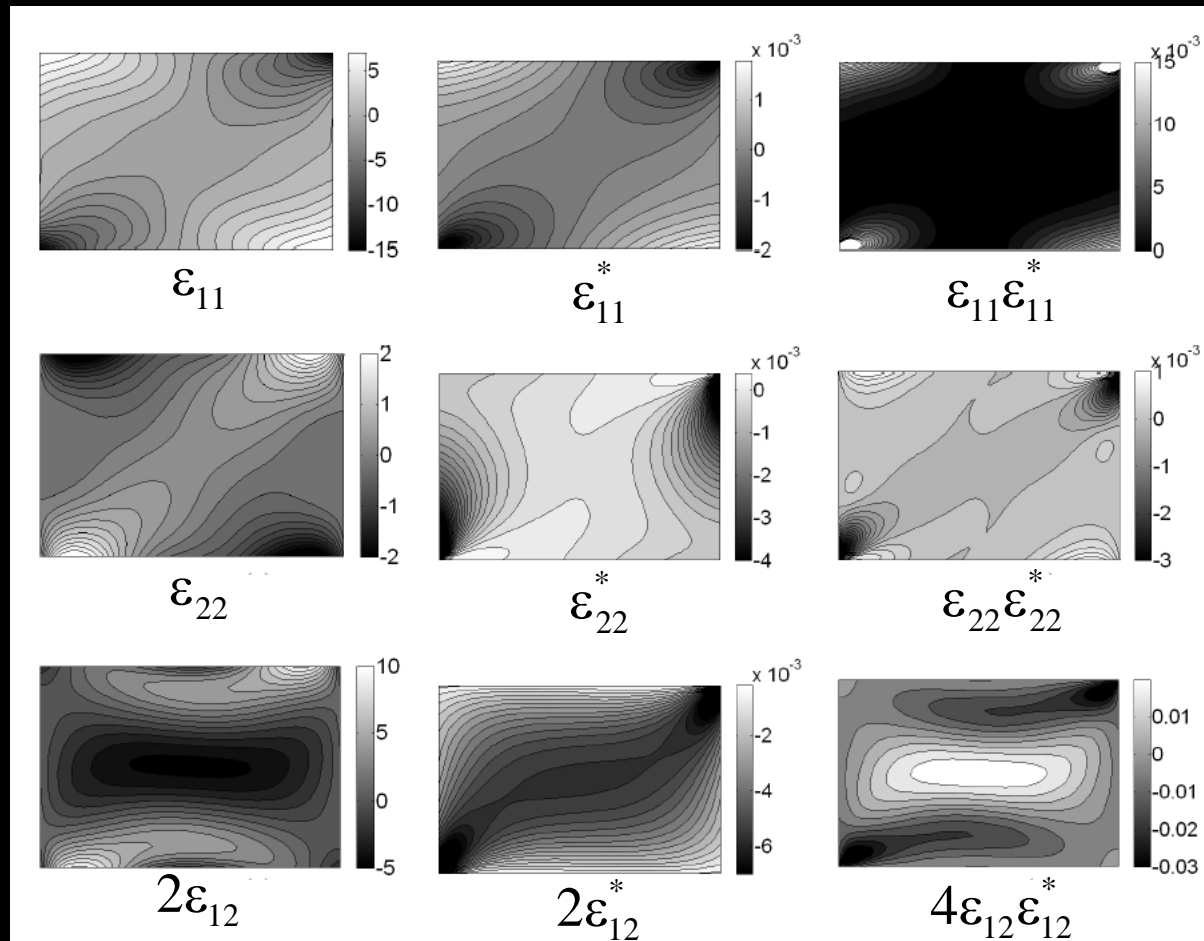
## ■ Validation

- 30 copies of noise: distribution of stiffness
- Increasing levels of noise



# VF selection

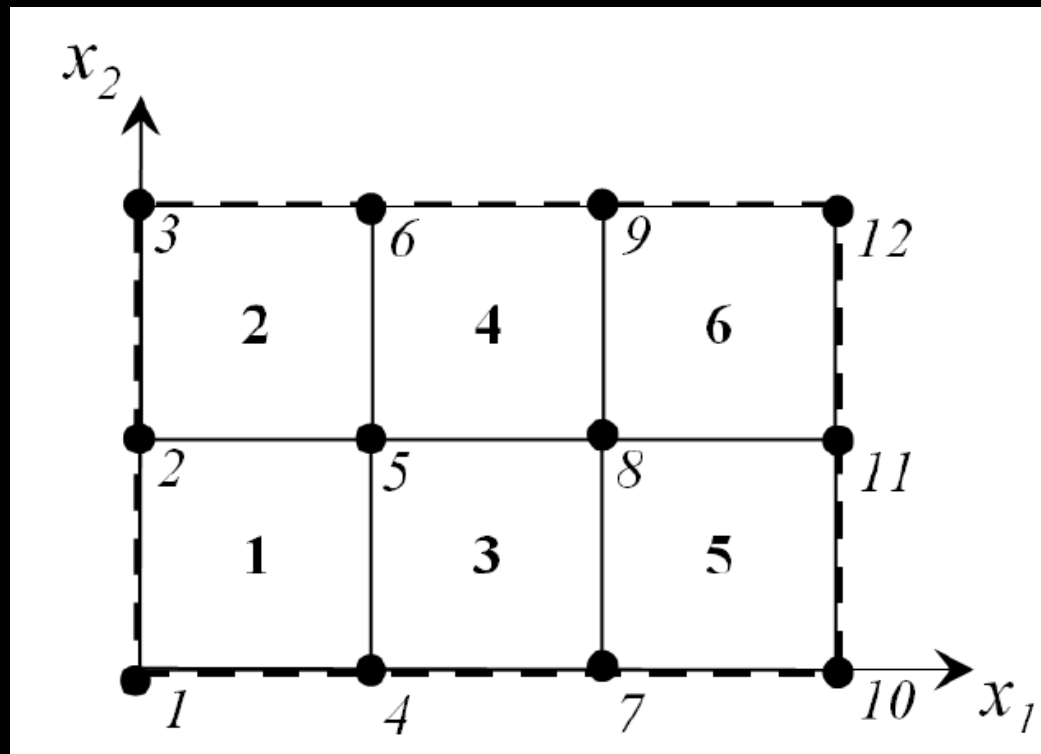
- Actual and virtual optimized virtual fields for  $Q_{11}$  ( $m=4, n=5$ )





# VF selection

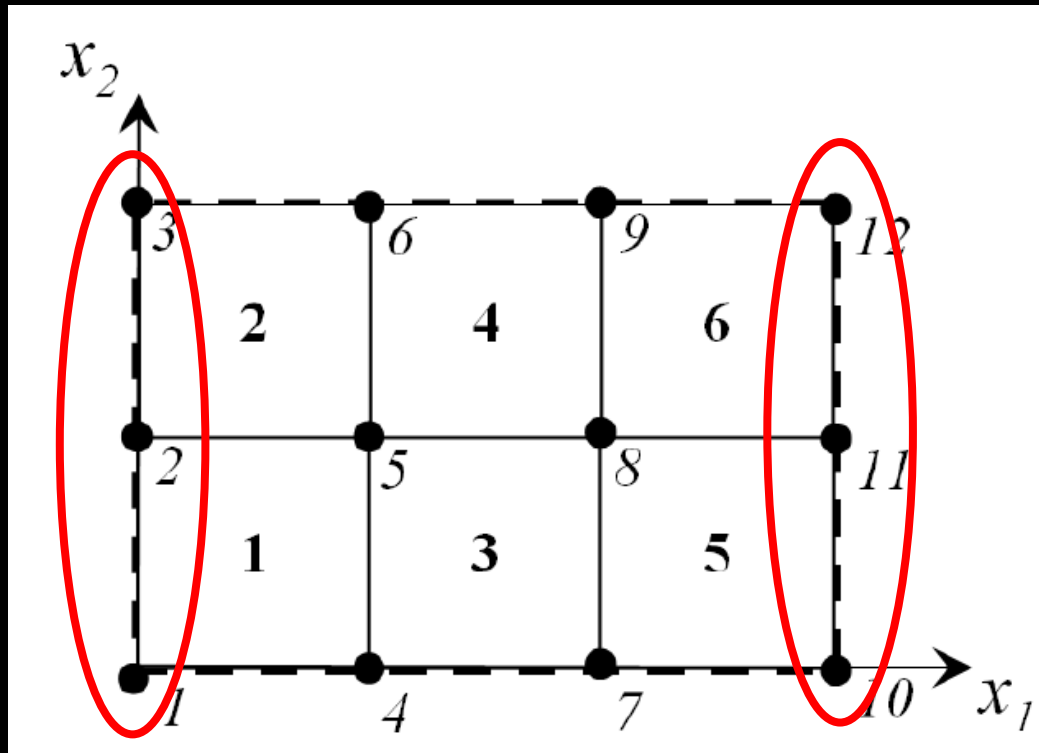
- Piecewise virtual fields
  - Finite element mesh
  - Virtual Dofs unknown of minimization problem



# VF selection

- Piecewise virtual fields
  - Advantage: more flexible
  - Easy for virtual BCs

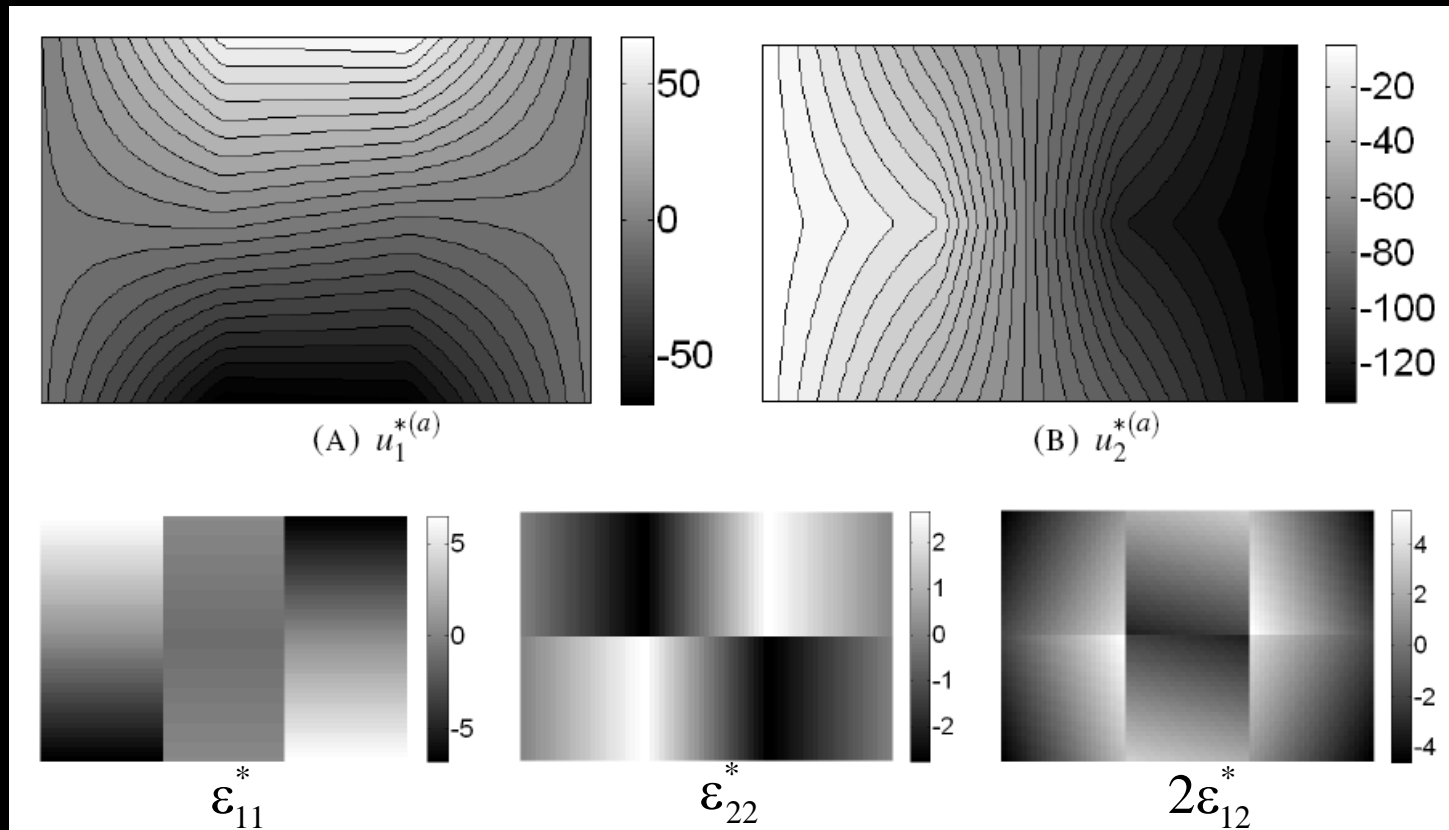
All dofs  
To zero



All  $u_1^* = 0$   
All  $u_2^* = \text{cst}$

# VF selection

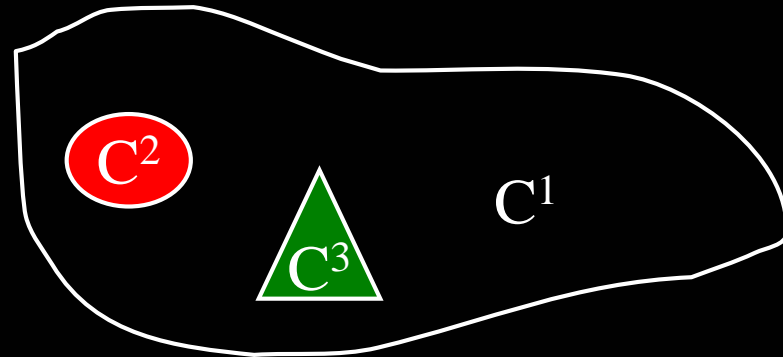
- Optimized piecewise virtual fields for  $Q_{11}$ 
  - Virtual strains discontinuous



# Heterogeneous materials

# Heterogeneous materials

- Discrete parameterization
  - Multimaterials
  - Sharp transitions



$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV = -\int_{C^1} \sigma_{ij} \varepsilon_{ij}^* dV - \int_{C^2} \sigma_{ij} \varepsilon_{ij}^* dV - \int_{C^3} \sigma_{ij} \varepsilon_{ij}^* dV$$

# Heterogeneous materials

- Discrete parameterization
  - 6 unknowns in isotropic linear elasticity
  - Contrast of stiffness: strain derivative!!
  - Subject to noise

$$\begin{aligned} &= -Q_{11}^1 \int_{C1} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV - Q_{12}^1 \int_{C1} (\varepsilon_1 \varepsilon_2^* + \varepsilon_2 \varepsilon_1^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV \\ &\quad - Q_{xx}^2 \int_{C2} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV - Q_{12}^1 \int_{C2} (\varepsilon_1 \varepsilon_2^* + \varepsilon_2 \varepsilon_1^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV \\ &\quad - Q_{xx}^3 \int_{C3} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV - Q_{12}^1 \int_{C3} (\varepsilon_1 \varepsilon_2^* + \varepsilon_2 \varepsilon_1^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV \end{aligned}$$

# Heterogeneous materials

- Continuous parameterization
  - Functionally graded material
  - Smooth transitions



- Polymer/composite with moisture ingress from one face
- Ceramic with varying pore density
- Etc...

# Heterogeneous materials

## ■ Continuous parameterization



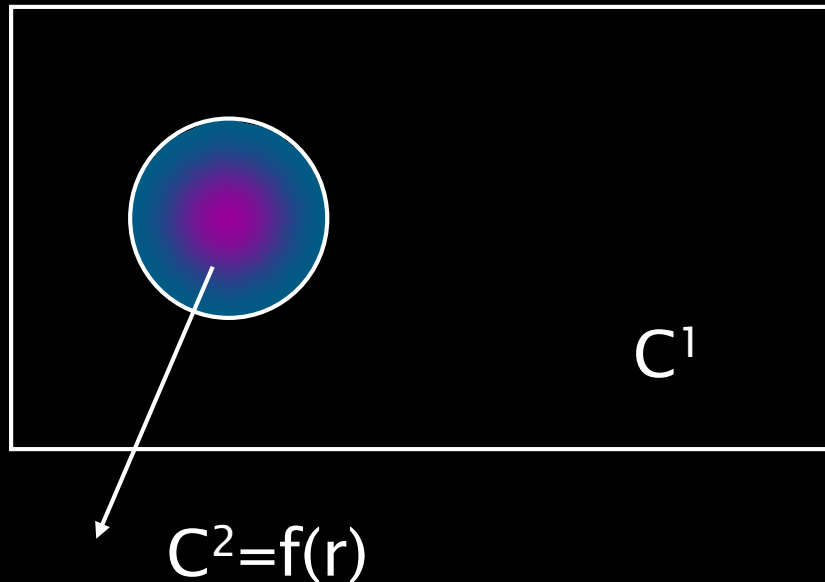
$$\begin{cases} Q_{11} = Q_{11}^0 + Q_{11}^1 \frac{x_2}{w} + Q_{11}^2 \left( \frac{x_2}{w} \right)^2 \\ Q_{12} = Q_{12}^0 + Q_{12}^1 \frac{x_2}{w} + Q_{12}^2 \left( \frac{x_2}{w} \right)^2 \end{cases}$$

$$\begin{aligned} & -Q_{11}^0 \int_{C1} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV - Q_{12}^0 \int_{C1} (\varepsilon_1 \varepsilon_2^* + \varepsilon_1 \varepsilon_2^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) dV \\ & -Q_{11}^1 \int_{C2} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) \frac{x_2}{w} dV - Q_{12}^1 \int_{C2} (\varepsilon_1 \varepsilon_2^* + \varepsilon_2 \varepsilon_1^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) \frac{x_2}{w} dV \\ & -Q_{11}^2 \int_{C3} (\varepsilon_1 \varepsilon_1^* + \varepsilon_2 \varepsilon_2^* + \frac{1}{2} \varepsilon_6 \varepsilon_6^*) \left( \frac{x_2}{w} \right)^2 dV - Q_{12}^2 \int_{C3} (\varepsilon_1 \varepsilon_2^* + \varepsilon_2 \varepsilon_1^* - \frac{1}{2} \varepsilon_6 \varepsilon_6^*) \left( \frac{x_2}{w} \right)^2 dV \end{aligned}$$



# Heterogeneous materials

- Mixed parameterization
  - Plate with local damage or manufacturing defect

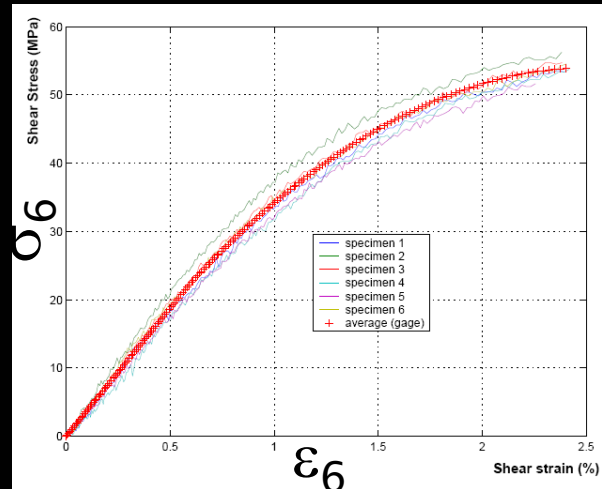


# Non-linear constitutive model

# Non-linear constitutive model

- Non-linear model with linear VFM formulation
  - Composite with shear strain softening

$$\sigma_6 = Q_{66}^0 \varepsilon_6 - K \varepsilon_6^3$$



$$Q_{11} \overline{\varepsilon_1 \varepsilon_1^*} + Q_{22} \overline{\varepsilon_2 \varepsilon_2^*} + Q_{12} \left( \overline{\varepsilon_1 \varepsilon_2^*} + \overline{\varepsilon_2 \varepsilon_1^*} \right) + Q_{66}^0 \overline{\varepsilon_6 \varepsilon_6^*} - K \overline{\varepsilon_6^3 \varepsilon_6^*} = - \frac{P u_2^*(x_1 = L)}{L b e}$$

# ... in elasto-plasticity

- Virtual fields method for non-linear constitutive equations

$$\sigma(t) = f(\epsilon, \boxed{Y_1^0, \dots, Y_m^0})$$

Measured:  
 $\epsilon(t), F(t)$

$$\min \left( - \int_V \underline{\underline{\sigma}}(\underline{\underline{\epsilon}}, Y_1 \dots Y_m) : \underline{\underline{\epsilon}}^* dV + \int_{\partial V} \underline{\underline{T}} \cdot \underline{\underline{u}}^* dS \right)^2 = 0$$

Iterative but no direct  
problem solved  
Fast (minutes)

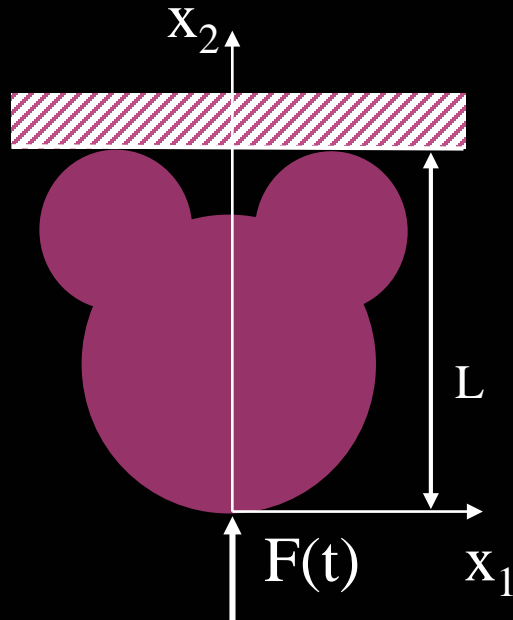
$u_1^*, \dots, u_n^*$

$Y_1, \dots, Y_m$

# Dynamics

# VFM in dynamics

## ■ Principle of virtual work in dynamics



$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \cancel{\int_{\partial V_f} T_i u_i^* dS} = \int_V \rho a_i u_i^* dV$$

If the load is unknown

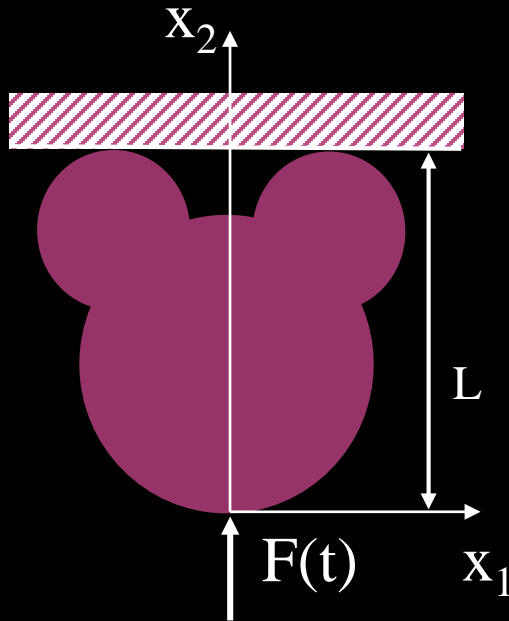
$$u_1^* = 0; u_2^* = x_2(x_2 - L)$$

In statics: with other virtual fields

$$[A]\{Q\} = \{0\}$$

Only stiffness ratios can be identified  
(Poisson's ratio)

# VFM in dynamics



$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV = \int_V \rho a_i u_i^* dV$$

$$u_1^* = 0; u_2^* = x_2(x_2 - L)$$

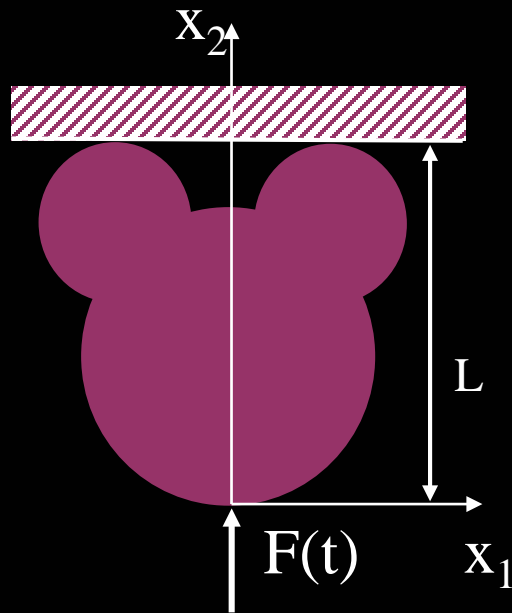
$$\varepsilon_1^* = 0; \varepsilon_2^* = 2x_2 - L; \varepsilon_6^* = 0$$

$$\overline{Q_{11}(2x_2 - L)\varepsilon_2} + \overline{Q_{12}(2x_2 - L)\varepsilon_1} = -\overline{\rho x_2(x_2 - L)a_2}$$

Measurements  $\begin{cases} u_1(x_1, x_2, t) \\ u_2(x_1, x_2, t) \end{cases}$

- Single spatial differentiation
- Double temporal differentiation

# VFM in dynamics



$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV = \int_V \rho a_i u_i^* dV$$

$$u_1^* = x_1 ; u_2^* = 0$$

$$\varepsilon_1^* = 1 ; \varepsilon_2^* = 0 ; \varepsilon_6^* = 0$$

$$Q_{11} \overline{\varepsilon_1} + Q_{12} \overline{\varepsilon_2} = -\rho \overline{x_1 a_1}$$

$$[A]\{Q\} = \{B(\rho, \vec{a})\}$$

Volume distributed load cell  
Embedded in the full-field data!!

More on Tuesday!



# Force identification

# Force identification

## ■ Standard VFM

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} T_i u_i^* dS = \int_V \rho a_i u_i^* dV$$

Measured  $\downarrow$

$$\sigma_{ij} = f(\varepsilon_{ij}, \boxed{E, \nu \dots})$$

Measured      Measured      **Unknown**

## ■ Force VFM

$$-\int_V \sigma_{ij} \varepsilon_{ij}^* dV + \int_{\partial V_f} \boxed{T_i} u_i^* dS = \int_V \rho a_i u_i^* dV$$

$\downarrow$

$$\sigma_{ij} = f(\varepsilon_{ij}, (E, \nu \dots))$$

Berry, A., Robin, O., & Pierron, F. (2014). *Journal of Sound and Vibration*, 333(26), 7151-7164.

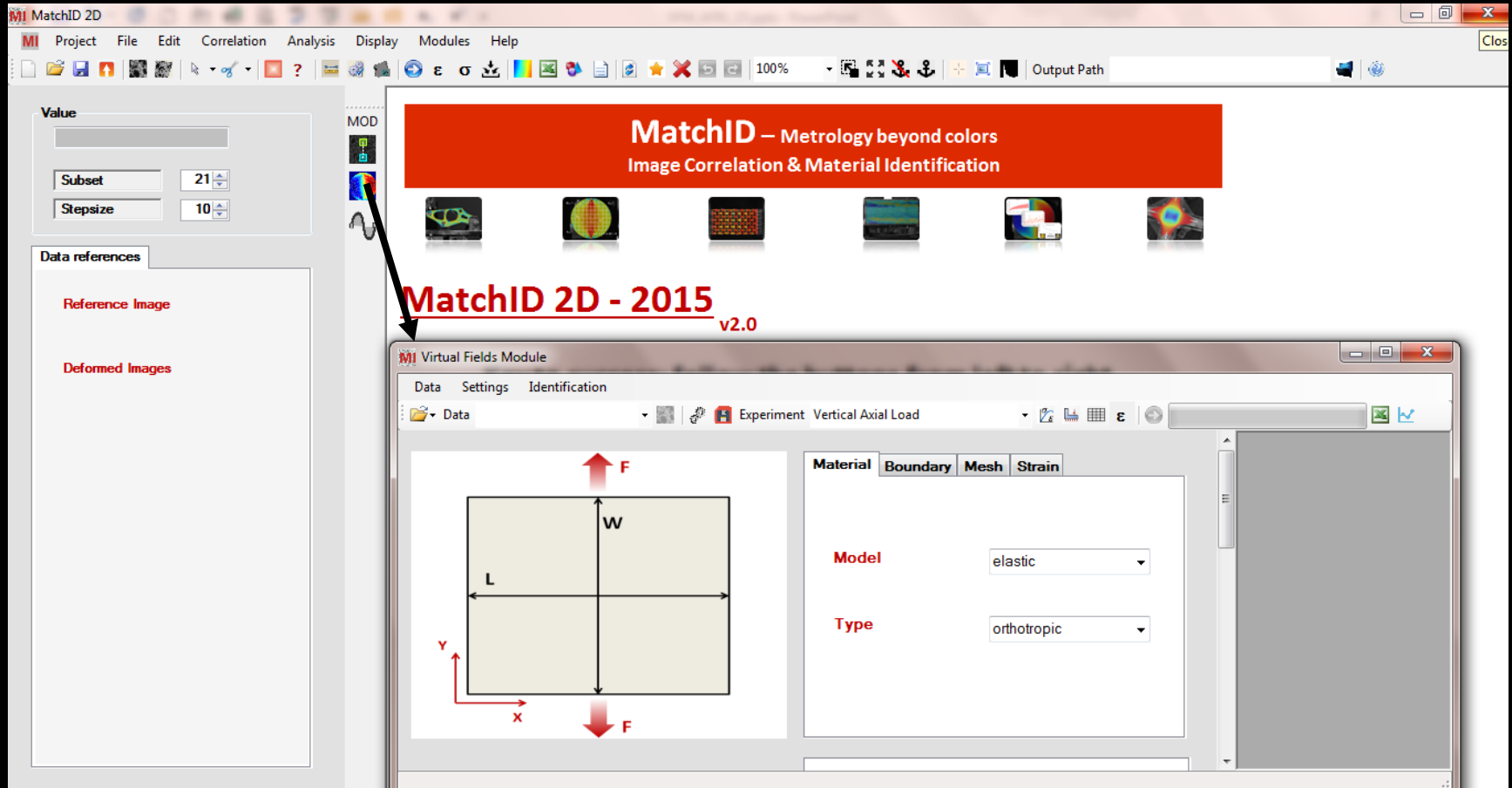
# Conclusion

# Conclusion

- VFM: operational tool, computationally simple and efficient
- Future challenges
  - Design and optimize test configurations, predict error bars (new operational standards)
  - Apply to 3D (X-ray CT + DVC)
  - High rate dynamics (see talk on Tuesday)
  - Optimization of VFs for non-linear VFM
  - Integrate with measurements: MatchID platform

# MatchID

- [www.matchIDmbc.com](http://www.matchIDmbc.com)



# More on the VFM

- Theory
- Applications
- Self-training

